

5D Meyers-Perry Black Holes

Can not be Over-Spun

by Cgedanken Experiments

Hongbao Zhang

Beijing Normal University

Based mainly on the work

arxiv:1711.04310 w/

Jincheng An, Jieru Shan, Suting Zhao

ICTS @ USTC

March 29, 2018

Outline of this talk

1 Cosmic Censorship Conjecture

① Weak

② Strong

2 High Dimensional case

3 Gedanken Experiments

① Old

② New

4 Iyer-Wald Formalism

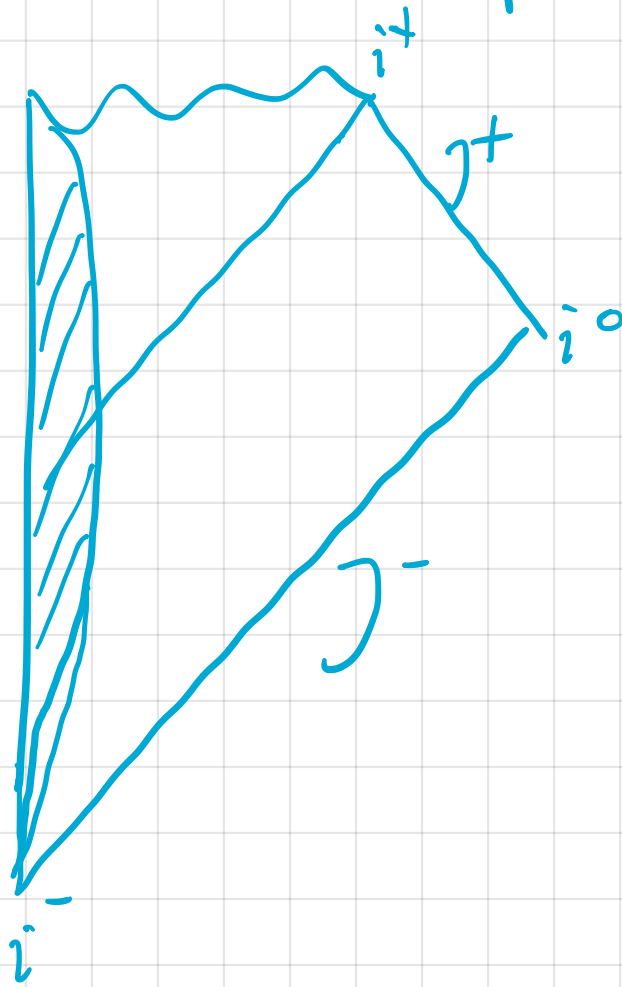
① First

② Second

5 5D Myers-Perry BHs

6 Conclusion

Cosmic Censorship Conjecture



4D
AF

① WCC:

The gravitational collapse from generic initial condition of suitable matters, gives rise to the singularity hidden inside the BHs.

② SCC:

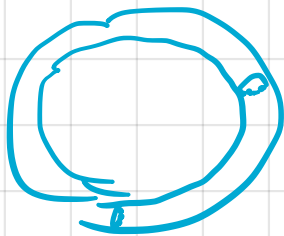
The singularity is spacelike or null

High Dimensional Case

① 5D Black strings

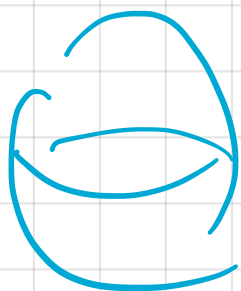
Lehner and Pretorius
2010 PRL

② 5D Black Rings



Figueras et al.
2016 PRL

③ 6D MP Black Holes

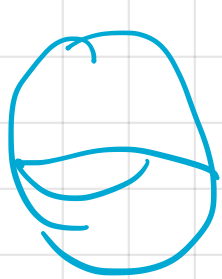


Figueras et al.
2017 PRL

What happens to 5D MP BHs?

Gedanken Experiments

① Old Version



$$\begin{aligned} E &= -\left(m\dot{u}_a + qA_a\right)^2 \\ &\geq -qA_a^2 \\ &= q\Phi_H \\ &= q \end{aligned}$$

Wald

1974 Ann. Phys.

($\Phi_H = 1$ for extremal BHs)

Nearly extremal BHs

WCC can be violated!

Kubeny

1999 PRD

② New Version

point particles
trajectories
first order

Generic Matters
NEC
Second order

Sorce and Wald

2017 PRD

Iyer-Wald formalism

Lagrangian n-form Iyer and Wald
1994, 1995 PRD

$$\underline{L} = \mathcal{L}(\phi) \underline{\varepsilon} \quad \phi = (\mathcal{G}_{ab}, \psi)$$

Variation of Lagrangian

$$\delta L = \underline{\varepsilon} \delta \phi + d \underbrace{\mathbb{W}(\phi, \delta \phi)}_{\substack{\uparrow \\ \text{Symplectic potential}}}$$

Symplectic Current

$$\omega(\phi, \delta \phi_1, \delta \phi_2) = \delta_1 \mathbb{W}(\phi, \delta \phi_2) - \delta_2 \mathbb{W}(\phi, \delta \phi_1)$$

Noether Current

$$\mathcal{J}_x = \mathbb{W}(\phi, \mathcal{L}_x \phi) - X \cdot \underline{L}$$

$$d\mathcal{J}_x = -E \mathcal{L}_x \phi$$

$$\mathcal{J}_x = d \underbrace{Q_x}_{\substack{\uparrow \\ \text{Noether charge}}} + \underbrace{C_x}_{\substack{\uparrow \\ \text{Constraint}}}$$

Noether charge

Constraint

Iyer-Wald Formalism

First Variational Identity

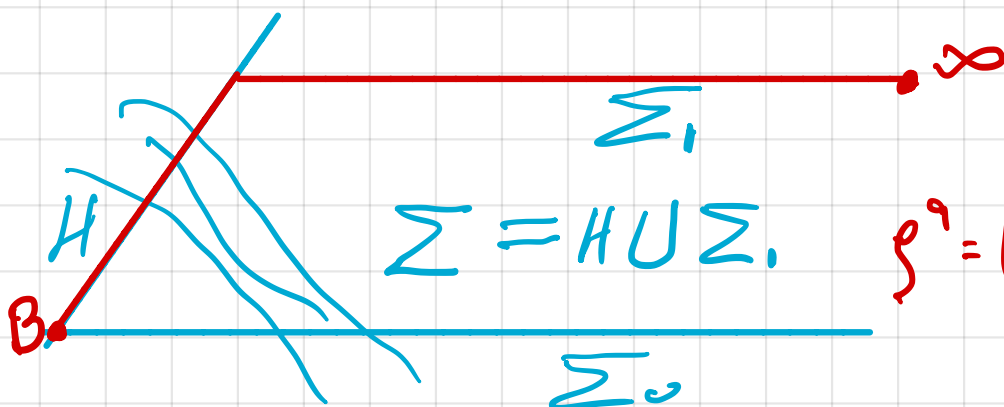
$$d[\delta Q_x - \chi \cdot \Theta(\phi, \delta\phi)] = \omega(\phi, \delta\phi, \mathcal{L}_x\phi) - \chi \cdot E \delta\phi - \delta C_x$$

Second Variational Identity

$$d[\delta^2 Q_\xi - \xi \cdot \delta\Theta(\phi, \delta\phi)] = \omega(\phi, \delta\phi, \mathcal{L}_\xi\delta\phi) - \xi \cdot \delta E \delta\phi - \delta^2 C_\xi$$

(ξ is a Killing field on the background with $E=0$)

Application to BHs



Iyer-Wald Formalism

Application to BHs

$$\delta M - \Omega^I \delta J_I =$$

$$\int_B [\delta Q_\xi - \xi \cdot \mathbb{W}(\phi, \delta\phi)] - \int_\Sigma \delta C_\xi$$

$$= - \int_H \delta C_\xi = \int_H \hat{\epsilon} \delta T_{ab} K^{ab} \geq 0$$

(null energy condition)

$$\delta^2 M - \Omega^I \delta^2 J_I = \mathcal{E}_\Sigma(\phi, \delta\phi) +$$

$$\int_B [\delta^2 Q_\xi - \xi \cdot \delta \mathbb{W}(\phi, \delta\phi)] - \int_\Sigma \xi \cdot \delta E \delta\phi - \int_\Sigma \delta^2 C_\xi$$

$$\text{Canonical Energy: } \mathcal{E}_\Sigma(\phi, \delta\phi) = \int_\Sigma \omega(\phi, \delta\phi, \mathcal{L}_\xi \delta\phi)$$

$$= - \int_H \delta^2 C_\xi + \mathcal{E}_H(\phi, \delta\phi) + \mathcal{E}_{\Sigma_1}(\phi, \delta\phi)$$

$$= \int_H \hat{\epsilon} \delta^2 T_{ab} K^{ab} + \frac{1}{4\pi} \int_H \hat{\epsilon} \delta \delta_{cd} \delta \sigma^{cd} \xi^a \nabla_a \mathcal{H} + \mathcal{E}_{\Sigma_1}(\phi, \delta\phi)$$

$$\geq \mathcal{E}_\Sigma(\phi, \delta\phi) = -\frac{\kappa}{8\pi} \delta^2 A_B^{MP} \text{ (trick!)}$$

5D Myers-Perry BHs

$$ds^2 = -dt^2 + \frac{\mu}{\Xi} (dt - a_1 \sin^2 \theta d\varphi^1 - a_2 \cos^2 \theta d\varphi^2)^2 + \frac{r^2}{\Pi - \mu r^2} dr^2 +$$

Myers and Perry
1986 Ann. Phys.

$$(r^2 + a_1^2) \sin^2 \theta (d\varphi^1)^2 + (r^2 + a_2^2) \cos^2 \theta (d\varphi^2)^2$$

$$\Xi = r^2 + a_1^2 \cos^2 \theta + a_2^2 \sin^2 \theta$$

$$\mu \geq (a_1 + a_2)^2 \text{ BHs}$$

$$\Pi = (r^2 + a_1^2)(r^2 + a_2^2)$$

$$\mu < (a_1 + a_2)^2 \text{ NSs}$$

Mass: $M = \frac{3\pi\mu}{8}$, Angular Momentum: $J_{\pm} = \frac{\pi\mu a_{\pm}}{4}$

Horizon radius: $r_h = \frac{\sqrt{\mu - (a_1 + a_2)^2} + \sqrt{\mu - (a_1 - a_2)^2}}{2}$

Area: $A = 2\pi^2 \mu r_h$

Angular velocity: $\Omega^{\pm} = \frac{a_{\pm}}{r_h^2 + a_{\pm}^2}$

Surface gravity: $\kappa = \frac{2r_h^2 + a_1^2 + a_2^2 - \mu}{\mu r_h}$

BH condition: $\alpha = \sqrt{32M^3 - 27\pi(J_1 + J_2)^2} \geq 0$

Conclusion

① Extremal 5D MP BHs can not be over-spun at the linear level

$$32M^3 - 27\pi (J_1 + J_2)^2 = 0$$

$$\delta M - \frac{9\pi (J_1 + J_2)}{16M^2} (\delta J_1 + \delta J_2) < 0$$

impossible because

$$\Omega^2 = \Omega \equiv \frac{9\pi (J_1 + J_2)}{16M^2} .$$

Conclusion

- ② Nearly extremal MP BHs
can be over-spun at the linear level
but can not be over-spun at the
quadratic level

$$f(\alpha) = 32M^3(\alpha) - 27\pi (J_1\alpha + J_2\alpha)^2$$

$$f(\alpha) = \underline{\alpha^2} + \gamma_1\alpha + \gamma_2\alpha^2 + O(\alpha^3, \alpha^2\alpha, \alpha^2, \alpha^3)$$

$$\gamma_1 = -6\sqrt{3\pi} \left(\sqrt{\frac{J_2}{J_1}} \delta J_1 + \sqrt{\frac{J_1}{J_2}} \delta J_2 \right) \alpha$$

$$\gamma_2 = 27\pi \left(\sqrt{\frac{J_2}{J_1}} \delta J_1 + \sqrt{\frac{J_1}{J_2}} \delta J_2 \right)^2$$

$$f(\alpha) = \underline{\left[\alpha - 3\sqrt{3\pi} \left(\sqrt{\frac{J_2}{J_1}} \delta J_1 + \sqrt{\frac{J_1}{J_2}} \delta J_2 \right) \right]^2}$$

+ high order terms

Conclusion

- ③ It is still possible to violate WCC around SD MP PMS by fully non-linear dynamics, although highly unlikely.

— Pan Figueras

Thank you for your listening!