

Jet TMDs

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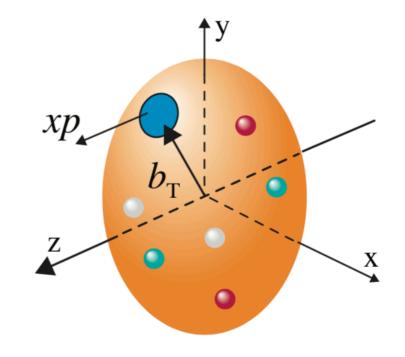
ICTS Seminar @ USTC

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TMD physics

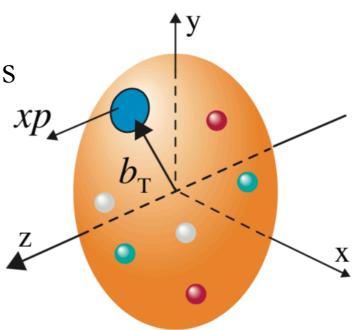
- Efforts in understanding nucleon structures using collisions
 - Collinear PDFs, FFs
 - Important for LHC physics
 - Universal formalism (somehow) wellestablished
 - But only longitudinal information

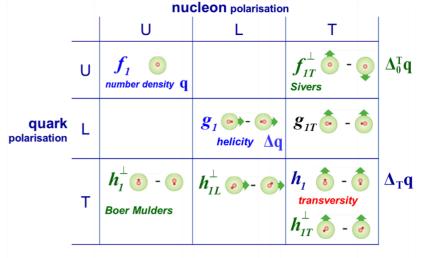
$$\frac{d\sigma}{dQ^2} \simeq \sum_{i,j=q,g} \int_0^1 dx_1 dx_2 \mathcal{H}_{ij}(Q^2,\mu^2) f_{i\leftarrow h}(x_1,\mu^2) f_{j\leftarrow h}(x_2,\mu^2)$$



TMD physics

- Efforts in understanding nucleon structures using collisions
 - Transvers Momentum Dependent PDFs, FFs
 - 3D imaging of Nucleon
 - Information on spin, orbital momentum
 - How to measure ? Force the system TM small enough
 - Non trivial extension of the collinear formalism, (Non-)universalities





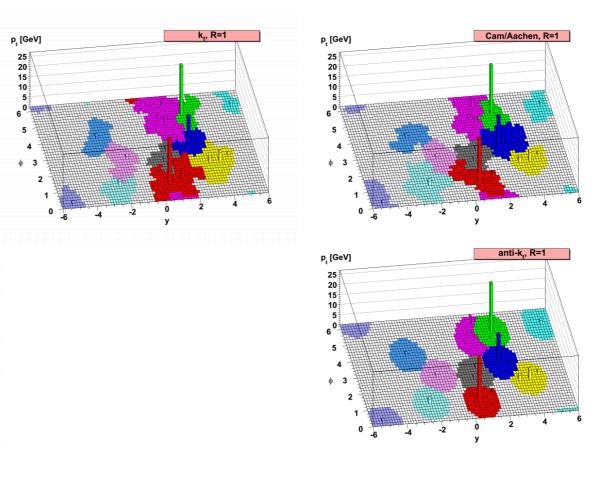
Jets

• A clustering of particles, very natural at the colliders

$$d_{ij} = \min(k_{i,T}^{2p}, k_{j,T}^{2p}) \frac{\Delta y^2 + \Delta \phi^2}{R^2}$$
$$d_{iB} = k_{i,T}^{2p}$$

- If d_{iB} is the smallest, then promote i as a jet and remove i from the event sample
- If d_{ij} is the smallest, then cluster i and j to ij
- Iterate

p=-1: anti-kT, p=1: kT, p=0 Cambridge/Archen



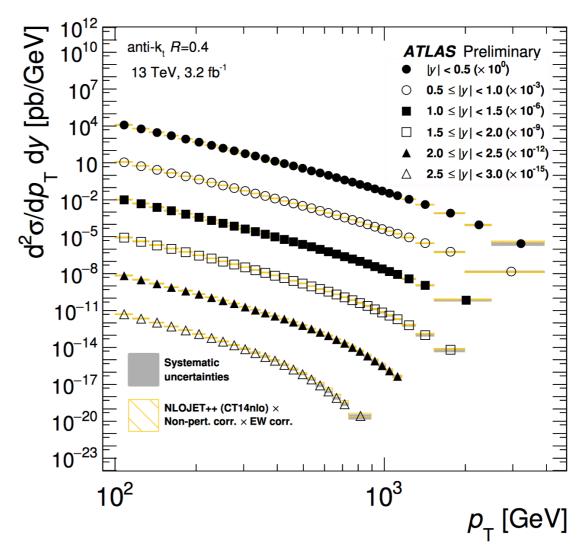
Cacciari, Salam, Soyez '08

Outline

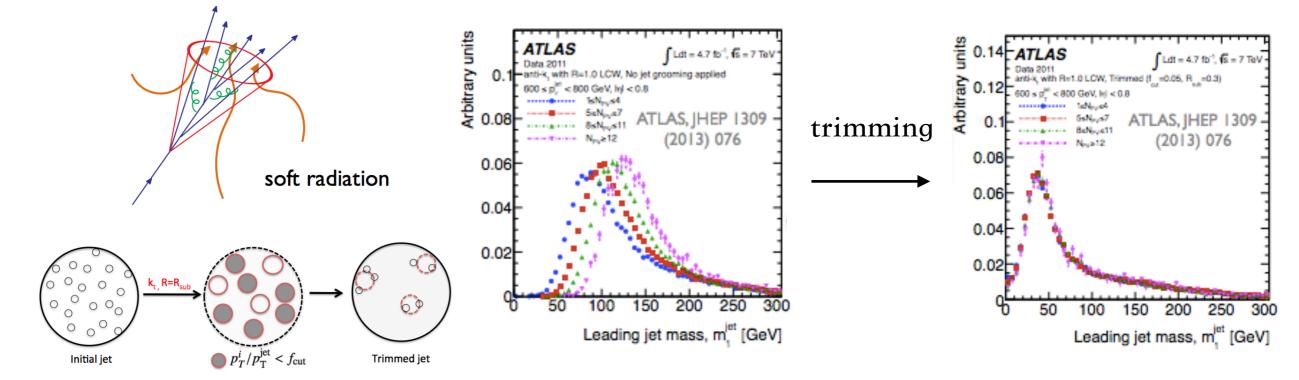
- Why jet physics?
 - Current status
 - Opportunity for TMDs
- Jet TMDs
 - In-jet TMDs
 - Jet imbalance
- Conclusions

- Most modern colliders are jet machines, LHC, EIC \cdots
 - Benchmark processes with large statistics and small errors
 - Probing QCD ranging over orders of magnitudes in scales

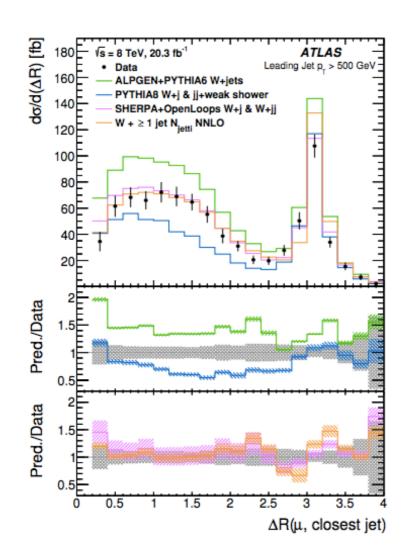
Global quantities probe scales of order 10GeV ~ 1TeV



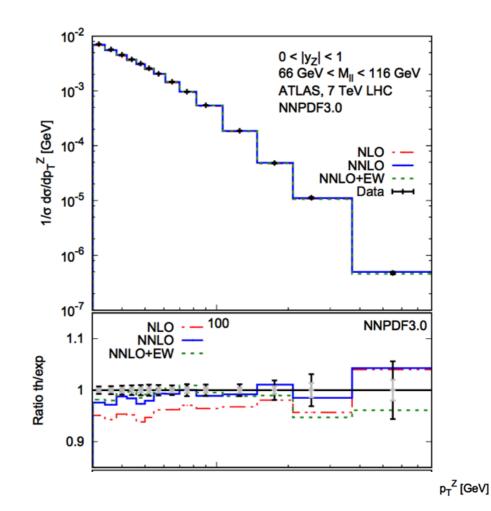
- Most modern colliders are jet machines, LHC, EIC \cdots
 - Benchmark processes with large statistics and small errors
 - Probing QCD ranging over orders of magnitudes in scales
 - Modern jet substructure techniques for extending the scale region, reducing contaminations, distinguish g/q jets …



- Precision first principle theoretical understanding
 - Exciting fixed order precision achievements
 - N²LO or N³LO predictions, V+J, J+J, DIS jet \cdots



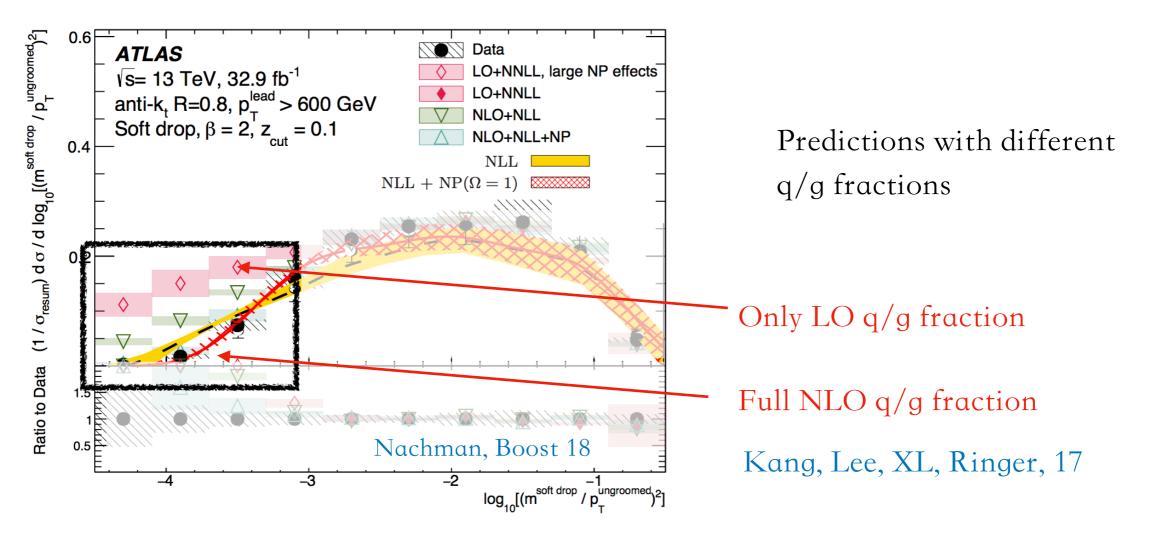
Boughezal, XL, Petriello +; Gehrmann, Glover, +; ···



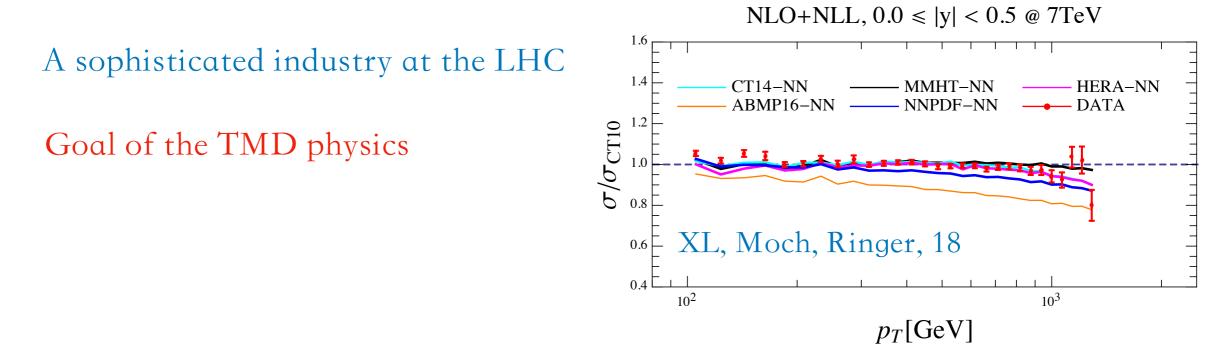
- Precision first principle theoretical understanding
 - Resummation frameworks for jet observables
 - Jet R
 - (un-)groomed jet shapes
 - • • •

Precision for jet physics, though challenging but feasible now

- Precision data + precision theory
 - Jet physics can be tested with high accuracy
 - Sensitive to minor differences in the theories



- Precision data + precision theory
 - Jet physics can be tested with high accuracy
 - Sensitive to minor differences in the theories
 - Interesting Non.Pert. quantities can be extracted with confidence, PDFs, (TMD-)FFs, Sivers Function…



- Jets for TMD physics
 - Conventional TMDs: DY, SIDIS and e⁺e⁻ thus far
 - Simple and well-established
 - Lots of data so far

- Jets for TMD physics
 - Conventional TMDs: DY, SIDIS and e⁺e⁻ thus far
 - Simple and well-established. Exciting opportunities for jet.
 - Lots of data so far. So are the jet processes, with more options

Modern jets will offer more

- Global TMDs
- In-jet TMDs

We will develop formalism for these, for both polarized and unpolarized TMDs

- TMD w.r.t different axis: beam/jet axis, photon …
- Probing the factorization breaking effects: Rogers-Mulders 10, Qiu-Collins 07.

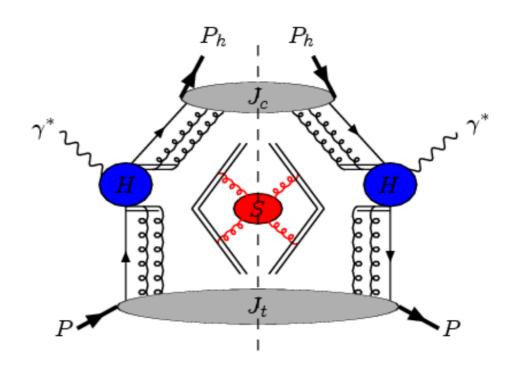
In-jet TMDs

- Conventional approach to TMD FFs
 - SIDSI, e⁺e⁻

For SIDIS

$$egin{aligned} F_{UU}(x_B,z_h,Q^2,P_{h\perp}) &= \sum_{q=u,d,s,\ldots} e_q^2 \int d^2 ec{k}_\perp d^2 ec{p}_\perp d^2 ec{\lambda}_\perp \ & imes q \left(x_B(k_\perp) \hat{q} \left(z_h, p_\perp
ight) \left(S(ec{\lambda}_\perp)
ight)^{-1}
ight. \ & imes H_{UU} \left(Q^2
ight) \delta^{(2)} \left(z_h ec{k}_\perp + ec{p}_\perp + ec{\lambda}_\perp - ec{P}_lpha_\perp
ight) \end{aligned}$$

Entanglement between TMD FFs and TMD PDFs Hard to probe gluon FFs



- Hadron in jet as an alternative
 - Possibility for gluon TMD FFs
 - With more experimental controls
 - Dis-entangle final TMDs from ISR

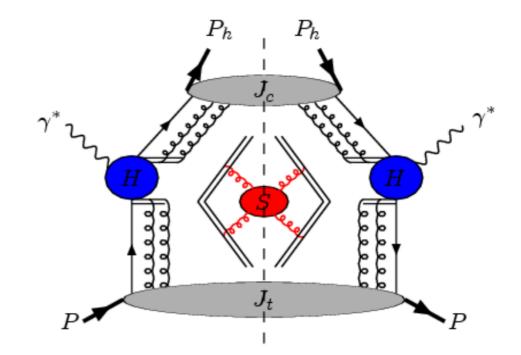
For SIDIS

$$egin{aligned} F_{UU}(x_B,z_h,Q^2,P_{h\perp}) &= \sum_{q=u,d,s,...} e_q^2 \int d^2 ec{k}_\perp d^2 ec{p}_\perp d^2 ec{\lambda}_\perp \ & imes q \left(x_B \left(ec{k}_\perp
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ight) \ & imes H_{UU} \left(Q^2
ight) \delta^{(2)} \left(z_h ec{k}_\perp + ec{p}_\perp + ec{\lambda}_\perp - ec{P}_{h\perp}
ight) \end{aligned}$$

For in-jet TMDs

The TMD PDFs will be absent!

Cleaner probe of the TMD FFs



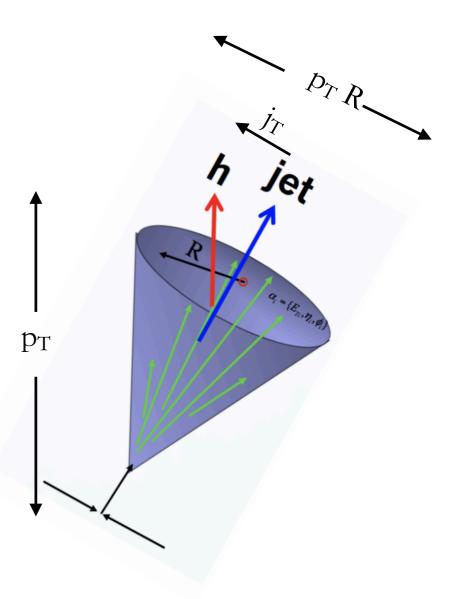
- Hadron TMDs in the inclusive jet production
 - A multi-scale problem

 $p_T >> p_T R >> j_T > \Lambda_{QCD}$

 $p_T \sim 20 - 1000 \text{GeV}, R \sim 0.2 - 0.5$ $j_T \sim 100 \text{MeV} - 5 \text{ GeV}$

Con:

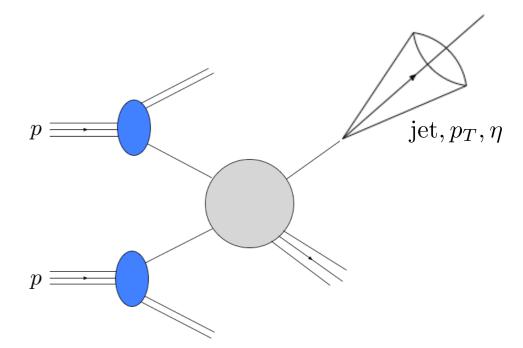
- Fixed order breaks down due to log(hierarchies)
 Pro:
- Factorization and resummation
- Universality



- Hadron TMDs in the inclusive jet production
 - Factorization

 $\frac{d\sigma^{pp \to (\text{jet }h)X}}{dp_T d\eta dz_h d^2 \boldsymbol{j}_\perp} = \sum_{a,b,c} f_a(x_a,\mu) \otimes f_b(x_b,\mu) \otimes H^c_{ab}(x_a,x_b,\eta,p_T/z,\mu) \otimes \mathcal{G}^h_c(z,z_h,\omega_J R,\boldsymbol{j}_\perp,\mu)$

H: partonic cross section for producing a parton(hadron), known to NLO (Vogelsang)f: collinear PDFsG: Universal semi inclusive jet function (analog of hadron FFs), all the j dependence encoded here. Same function for EIC.



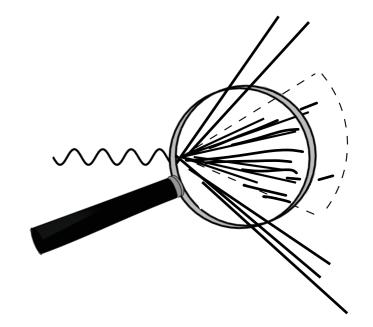
Fully differential in the jet kinematics

- Hadron TMDs in the inclusive jet production
 - Factorization

 $\frac{d\sigma^{pp \to (\text{jet }h)X}}{dp_T d\eta dz_h d^2 \boldsymbol{j}_{\perp}} = \sum_{a,b,c} f_a(x_a,\mu) \otimes f_b(x_b,\mu) \otimes H^c_{ab}(x_a,x_b,\eta,p_T/z,\mu) \otimes \mathcal{G}^h_c(z,z_h,\omega_J R,\boldsymbol{j}_{\perp},\mu)$

H: partonic cross section for producing a parton(hadron), known to NLO (Vogelsang)f: collinear PDFsG: Universal semi inclusive jet function (analog of hadron FFs), all the j dependence encoded here. Same function for EIC.

Valid for $p_T R \sim j$



- Hadron TMDs in the inclusive jet production
 - Factorization

$$egin{aligned} \mathcal{G}^h_c(z,z_h,\omega_J R,oldsymbol{j}_\perp,\mu) =& \mathcal{H}_{c o i}(z,\omega_J R,\mu) \int d^2oldsymbol{k}_\perp d^2oldsymbol{\lambda}_\perp \delta^2\left(z_holdsymbol{\lambda}_\perp+oldsymbol{k}_\perp-oldsymbol{j}_\perp
ight) \ & imes D_{h/i}(z_h,oldsymbol{k}_\perp,\mu,
u) S_i(oldsymbol{\lambda}_\perp,\mu,
u R) \;, \end{aligned}$$

H: radiations out-side the jet cone, not allowed into the jet

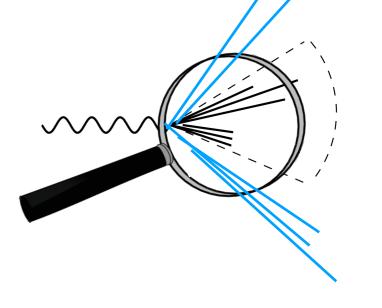
S: soft radiations

D: unsubtracted TMD FFs. Which is Non.Pert., if

 $j \sim \Lambda_{\rm QCD}$

S.D : standard TMD FFs as in SIDIS Universal

Valid for $j \leq p_T R$



- Hadron TMDs in the inclusive jet production
 - Resummation

perturbative Sudakov factor

$$\hat{\mathcal{D}}_{h/i}(z_h, \boldsymbol{j}_\perp; \mu_J) = rac{1}{z_h^2} \int rac{b \, db}{2\pi} J_0(j_\perp b/z) C_{j\leftarrow i} \otimes D_{h/j}(z_h, \mu_{b*}) e^{-S_{ ext{pert}}^i(b_*, \mu_j) - S_{ ext{NP}}^i(b, \mu_J)}$$

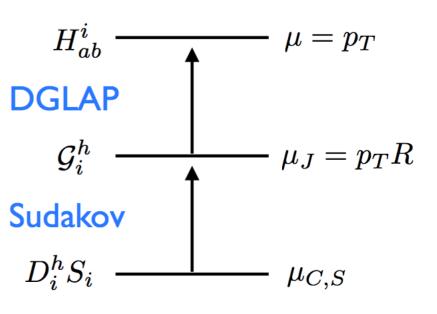
Perturbative matching

Non-perturbative input

RG evolution

$$S_{\text{pert}}^{i}(b_{*},\mu_{J}) = \int_{\mu_{b_{*}}}^{\mu_{J}} \frac{d\mu'}{\mu'} \left(\Gamma_{\text{cusp}}^{i} \ln\left(\frac{\mu_{J}^{2}}{\mu'^{2}}\right) + \gamma^{i}\right)$$

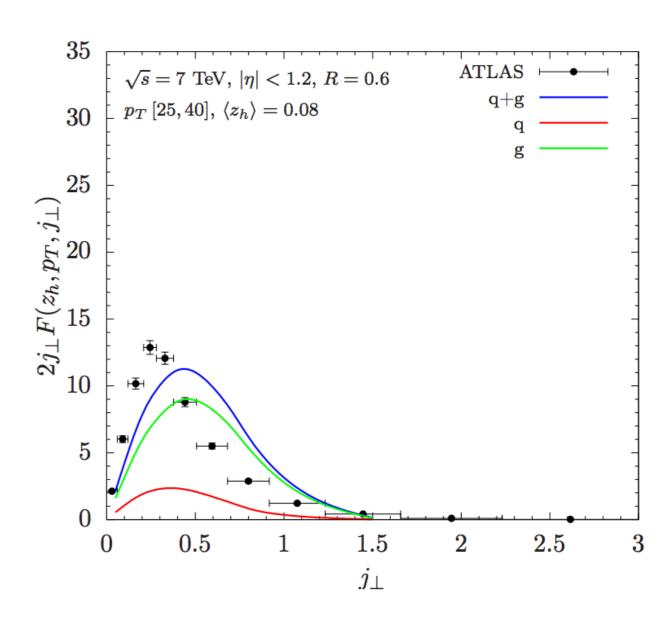
NP input from standard SIDIS/e+e-Sun, Kang, Prokudin, Yuan'16



- Hadron TMDs in the inclusive jet production
 - Phenomenology
 - No z_h bin measures but integrated over [0,1]
 - Would be great if RHIC has such a measurement

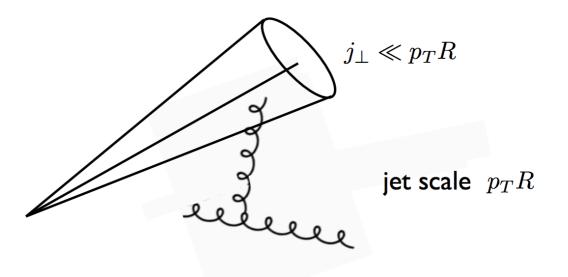
35 siTMDFF $\sqrt{s} = 7$ TeV, $|\eta| < 1.2, R = 0.6$ ATLAS $p_T [25, 40], < z_h >= 0.08$ 30 25ATLAS, Eur. Phys. J C71 (2011) 1795 $\begin{array}{ccc} 2j_{\perp}F(z_h,p_T,j_{\perp}) \\ 0 & 12 & 02 \\ 0 & 0 \end{array}$ 10 50 0.52 1.52.53 1 0 $j_{\perp} \, [{
m GeV}]$

- Hadron TMDs in the inclusive jet production
 - Phenomenology
 - Control pT to pick the g/q TMDFFs
 - Better opportunity for gluon TMDFFs



- Hadron TMDs in the inclusive jet production
 - Interesting to the HEP QCD
 - Non-global logarithms

$$d\sigma = \sum_{abcd} f_a f_b H^c_{ab} \mathcal{H}_{cd} \hat{\mathcal{D}}_d \times S_{d,\text{NGL}}$$

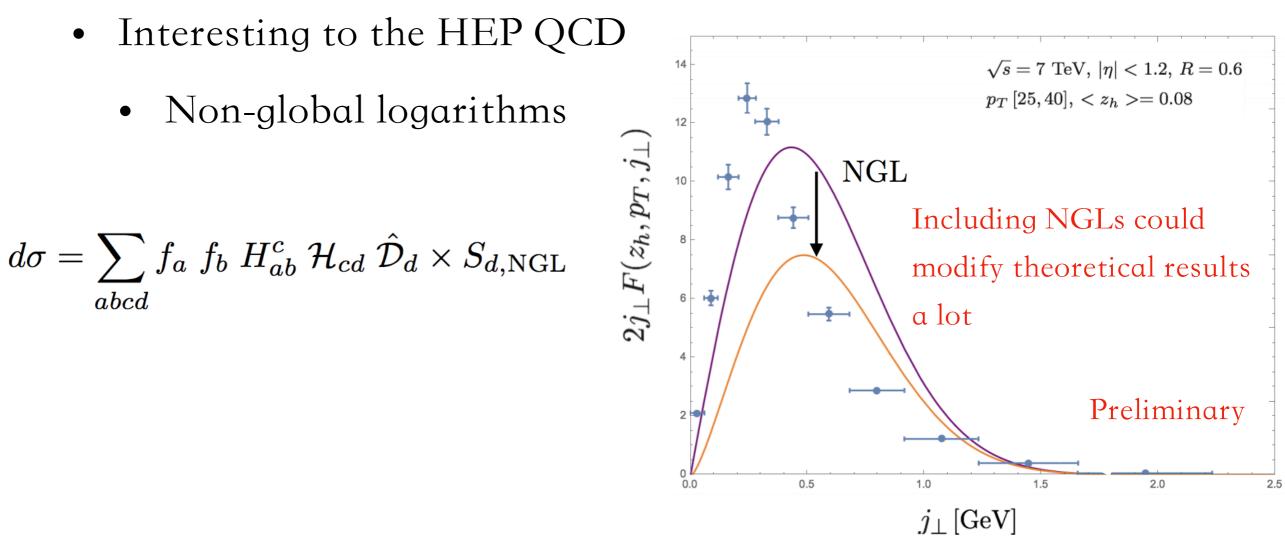


- Discovered by G. Salam et. al. (hep-ph/0104277), start at NNLO
- Can not be resummed using conventional CSS formalism
- Reduced for ratios (asymmetry \cdots)

NGLs are there, but NO direct phenomenology studies now.. too small

However, things may be different here since j is too small!

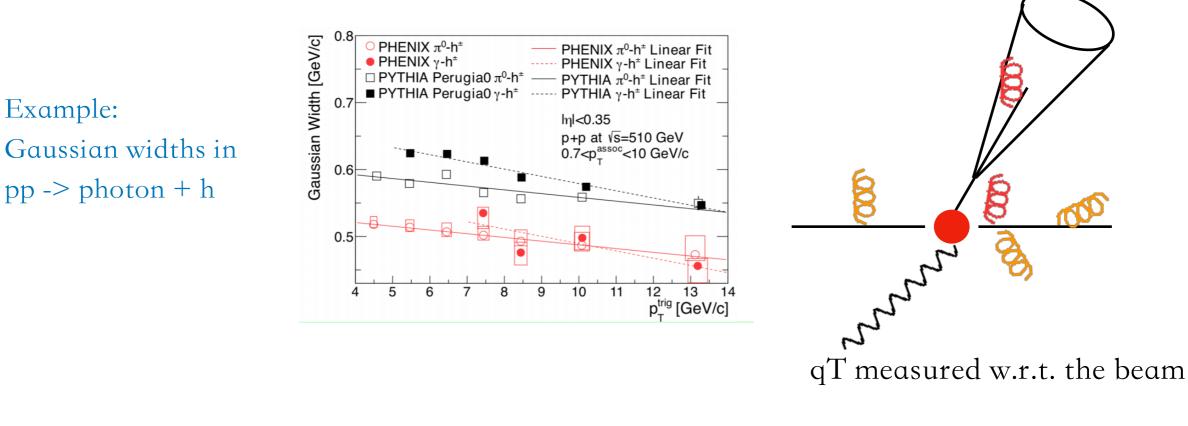
• Hadron TMDs in the inclusive jet production



In-jet TMD may provide the very first experimental demonstration of the NGLs Alternatively, groomed away NGLs, see Makris, Neill, Vaidya, 2017, but hard to relate to the normal TMD.

Jet Imbalance

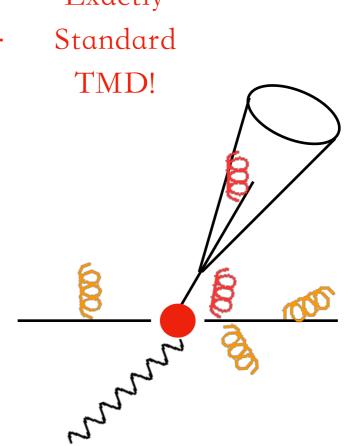
- For probing the factorization breaking effects
 - Theo. says the factorization breaks down, but how large is the effect?
 - exp. does measure the process or sth. similar, and no theory to compare with



• For probing the factorization breaking effects

$$\begin{aligned} \frac{d\sigma}{d\mathcal{PS}} &= \sum_{a,b,c} \int d\phi_J \int \prod_i^4 d^2 \vec{k}_{i\perp} \delta^{(2)} (\vec{q}_\perp - \sum_i^4 \vec{k}_{i\perp}) \\ &\times f_a^{\text{unsub}}(x_a, k_{1\perp}^2) f_b^{\text{unsub}}(x_b, k_{2\perp}^2) S_{n\bar{n}n_J}^{\text{global}}(\vec{k}_{3\perp}) & \longleftarrow \\ &\times S_{n_J}^{cs}(\vec{k}_{4\perp}, R) H_{ab \to c\gamma}(p_\perp) J_c(p_\perp R) \,, \end{aligned}$$
 Exactly Standard TMD!

- Everything are known to at least NLO
- Resummation both qT and R at NLL
- the prediction is precise
- Deviation from future data will shed light on the factorization breaking effects



• For probing the factorization breaking effects

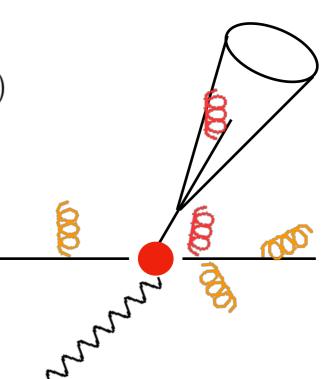
$$\frac{d\Delta\sigma}{d\mathcal{PS}} = \epsilon^{\alpha\beta} s^{\alpha}_{\perp} \sum_{a,b,c} \int d\phi_J \int \prod_i^4 d^2 \vec{k}_{i\perp} \delta^{(2)} (\vec{q}_{\perp} - \sum_i^4 \vec{k}_{i\perp})$$

$$\times \frac{k^{\beta}_{1\perp}}{M} f^{\perp \,\text{SIDIS}}_{1T,a} (x_a, k^2_{1\perp}) f^{\text{unsub}}_b (x_b, k^2_{2\perp})$$

$$\times S_{n\bar{n}n_J} (\vec{k}_{3\perp}) S^{cs}_{n_J} (\vec{k}_{4\perp}, R) H^{\text{Sivers}}_{ab \to c\gamma} (p_{\perp}) J_c (p_{\perp}R)$$

- Can be done for the polarized case
- Sensitive to the Sivers function in SIDIS

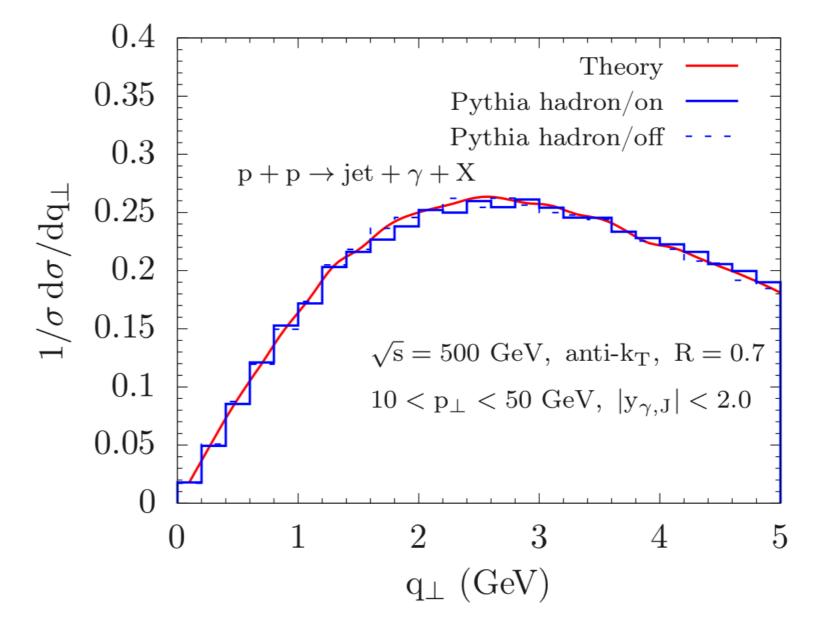
NB, can be readily extended to the hadron imbalance for the PHENIX results



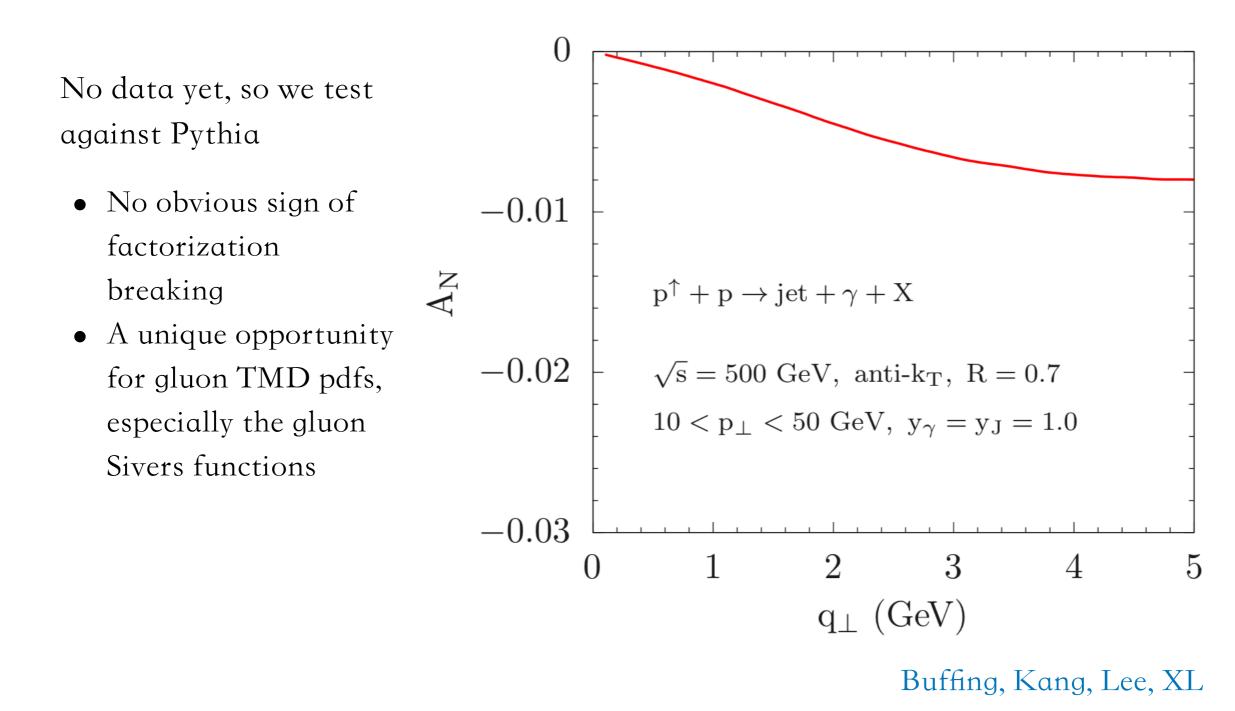
• For probing the factorization breaking effects

No data yet, so we test against Pythia

- No obvious sign of factorization breaking
- A unique opportunity for gluon TMD pdfs, especially the gluon Sivers functions



• For probing the factorization breaking effects



Conclusions

- We can control jets as good as DY, SIDIS, e⁺e⁻, both theoretically and experimentally.
- Jets provide different and thorough approaches to the TMD physics
 - In jet TMDs
 - Global TMDs
 - Precise phenomenological test

Thanks