

Quantum Mirror Map for Del Pezzo Geometries

arXiv:1908.11396

Yuji Sugimoto

(Univ. of Science and Technology of China)

with Sanefumi Moriyama (Osaka City Univ.)
Tomohiro Furukawa

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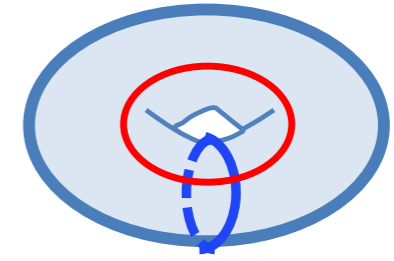
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Susy. Gauge Theory

encoded in algebraic curve

e.g.) D5 del Pezzo $0 = \left[\left(Q^{1/2} + Q^{-1/2} \right) \left(P^{1/2} + P^{-1/2} \right) \right]^2 - z/\alpha$



A-period

B-period

Mirror Map $\Pi_A(z) \sim c_1 z + c_2 z + \dots$
redefining the variables

Free energy relating to
BPS indices

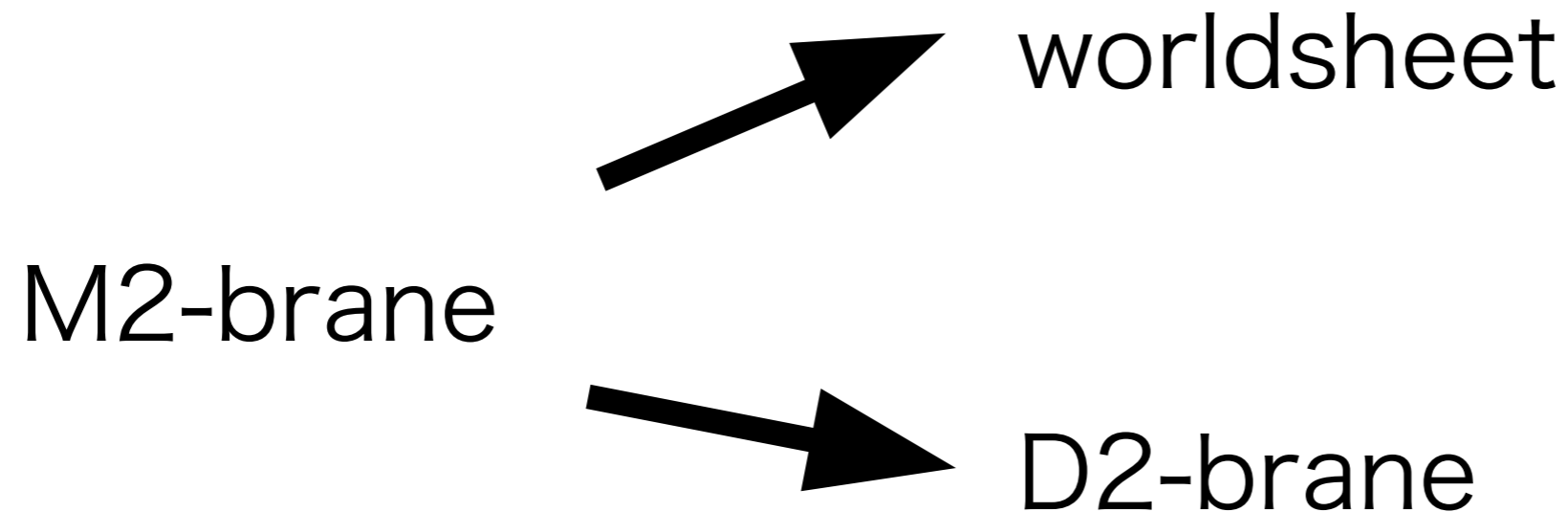
not well studied

well studied

Susy. Gauge Theory

e.g. Free energy of ABJM

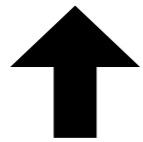
$$F = (\text{Worldsheet inst.}) + (\text{Membrane inst.}) \\ + (\text{bound state})$$



Susy. Gauge Theory

e.g. Free energy of ABJM

$$F = (\text{Worldsheet inst.}) + (\text{Membrane inst.})$$



A-period

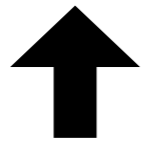
(mirror map)

- they are given by (refined) topological string
(analogous to topological string ver. of AdS/CFT)
- Well known example of AdS/CFT in this case is
duality between Chern-Simons and topological string

Susy. Gauge Theory

e.g. Free energy of ABJM

$$F = (\text{Worldsheet inst.}) + (\text{Membrane inst.})$$



A-period

(mirror map)

A-period plays a key role to determine the structure of the theory.



We want to understand the physical meaning of A-period.

Susy. Gauge Theory

e.g. Free energy of ABJM

$$F = (\text{Worldsheet inst.}) + (\text{Membrane inst.})$$



A-period

(mirror map)

Recent progress

- Some kinds of ABJM has group theoretical structure coming from symmetry of curve.

(Kubo-Moriyama-Nosaka(2018))

- Free energy corresponds to **B-period**.

Then, what about A-period ?

Generalized ABJM theory [Honda, Moriyama, 2014]

$$(U(N)_k \times U(N)_{-k} \times U(N)_k \times U(N)_{-k})$$

encoded in D_5 del Pezzo with quantization

$$\hat{H} = \left[\left(\hat{Q}^{1/2} + \hat{Q}^{-1/2} \right) \left(\hat{P}^{1/2} + \hat{P}^{-1/2} \right) \right]^2 - z/\alpha$$
$$\hat{Q}\hat{P} = q\hat{P}\hat{Q}$$

$$\Pi_A(z)$$



$$\Pi_A(z, \hbar)$$
$$(q = e^{i\hbar})$$

Quantum Mirror Map

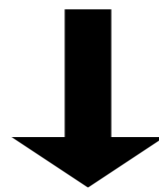
Generalized ABJM theory [Honda, Moriyama, 2014]

$(U(N)_k \times U(N)_{-k} \times U(N)_k \times U(N)_{-k})$

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$$\hat{H} = \left[\left(\hat{Q}^{1/2} + \hat{Q}^{-1/2} \right) \left(\hat{P}^{1/2} + \hat{P}^{-1/2} \right) \right]^2 - z/\alpha$$
$$\hat{Q}\hat{P} = q\hat{P}\hat{Q} \quad [\text{SU}(2)]^3 \text{ sym.}$$

- $[\text{SU}(2)]^3$ comes from breaking of D_5 sym.
- The sym. breaking is observed in **B-period**
(Kubo-Moriyama-Nosaka(2018))



We consider quantum mirror map of D_5 del Pezzo

Result

A-period is similar to **B-period**

- expressed by Weyl characters
- coef. of characters = Integer
- has same reps. as those in **B-period**

Plan

A-period is similar to **B-period**

① ● expressed by Weyl characters

② [● coef. of characters = Integer

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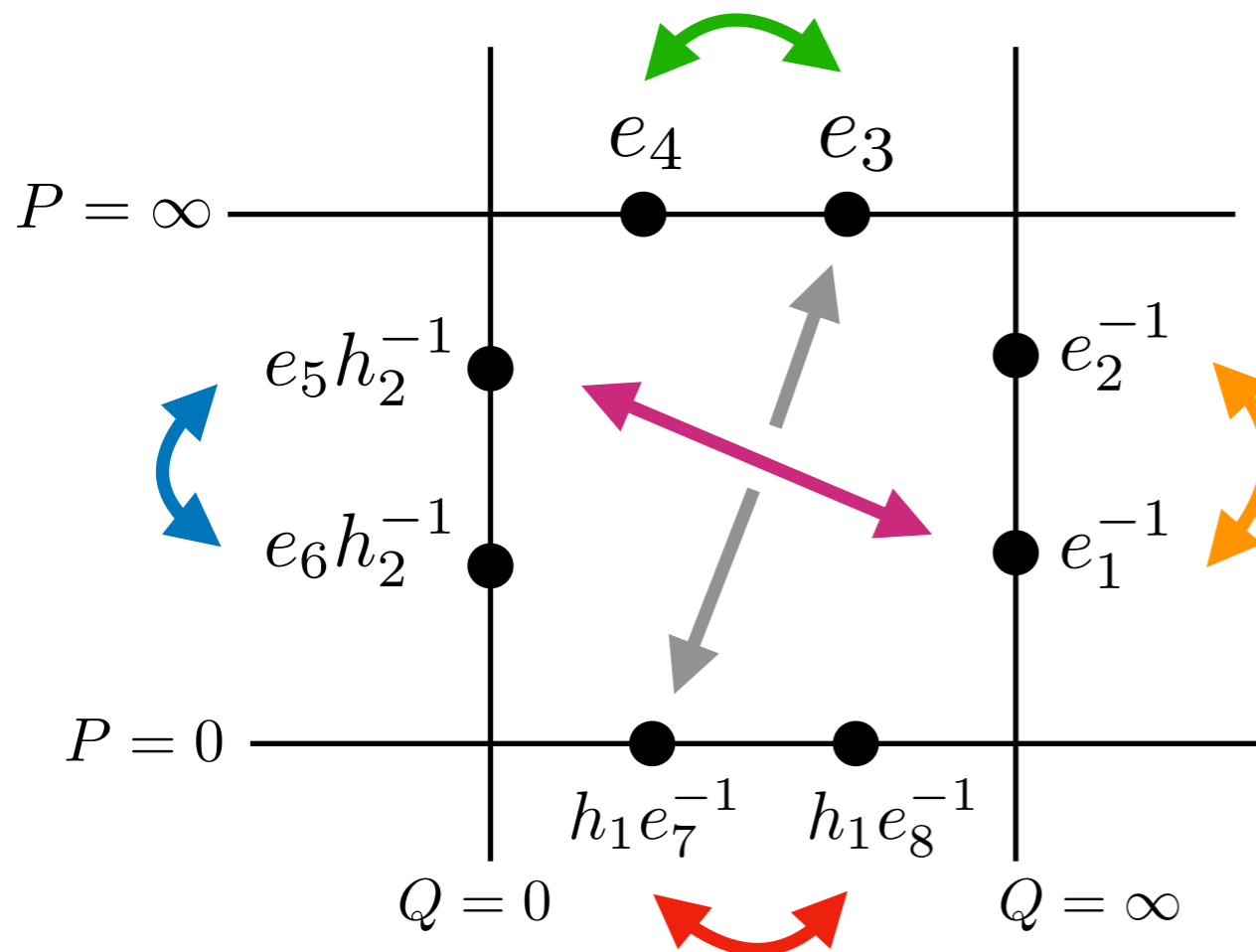
D₅ Del Pezzo geometry

Kubo-Moriyama-Nosaka(2018)

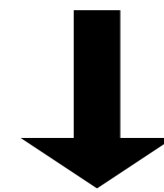
$$\begin{aligned}
 0 = & e_3 e_4 Q^{-1} P & - (e_3 + e_4) P & + Q P \\
 & - h_2^{-1} e_3 e_4 (e_5 + e_6) Q^{-1} & - \frac{z}{\alpha} & - (e_1^{-1} + e_2^{-1}) Q \\
 & + h_1^2 (e_1 e_2 e_7 e_8)^{-1} Q^{-1} P^{-1} & - h_1 (e_1 e_2)^{-1} (e_7^{-1} + e_8^{-1}) P^{-1} & + (e_1 e_2)^{-1} Q P^{-1}
 \end{aligned}$$

● $h_1^2 h_2^2 = \prod_{i=1}^8 e_i$

classical curve



invariant under exchanging of asymptotic values



D₅ Weyl sym.

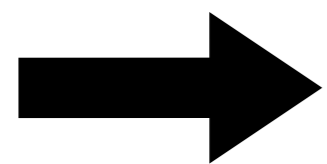
D₅ Del Pezzo geometry Kubo-Moriyama-Nosaka(2018)

$$\frac{\widehat{H}}{\alpha} + \frac{z}{\alpha} = e_3 e_4 \widehat{Q}^{-1} \widehat{P} - (e_3 + e_4) \widehat{P} + \widehat{Q} \widehat{P} - h_2^{-1} e_3 e_4 (e_5 + e_6) \widehat{Q}^{-1} + \frac{E}{\alpha} - (e_1^{-1} + e_2^{-1}) \widehat{Q} + h_1^2 (e_1 e_2 e_7 e_8)^{-1} \widehat{Q}^{-1} \widehat{P}^{-1} - h_1 (e_1 e_2)^{-1} (e_7^{-1} + e_8^{-1}) \widehat{P}^{-1} + (e_1 e_2)^{-1} \widehat{Q} \widehat{P}^{-1}$$

quantum curve

10 parameters - (2 + 2 + 1) parameters

- 8 asymptotic values determine curve 2
- $(\widehat{P}, \widehat{Q}) \sim (A\widehat{P}, B\widehat{Q})$ 2
- $h_1^2 h_2^2 = \prod_{i=1}^8 e_i$ 1



5 parameters $(e_1, e_3, e_5, \bar{h}_1, \bar{h}_2)$
 $(\bar{h}_1 = qh_1, \bar{h}_2 = q^{-1}h_2)$

Quantum mirror map

Aganagic, Cheng, Dijkgraaf Krefl, Vafa(2011)

$$\Pi_A(z, \hbar) \sim \oint \frac{\log P[X]}{X} dX = Ez^{-1} + \left(-\frac{E^2}{2} - A_2 \right) z^{-2} + \dots$$

$$\left[\begin{array}{l} P[X] = \frac{\Psi[q^{-1}X]}{\Psi[X]}, \quad \hat{Q}\Psi[X] = X\Psi[X] \quad \left[\frac{\hat{H}}{\alpha} + \frac{z}{\alpha} \right] \Psi[X] = 0 \\ \text{solve Schrödinger eq. order by order} \end{array} \right]$$

$$\frac{A_2}{\alpha^2 q^{-1}} = \frac{e_3}{e_1} + \frac{\bar{h}_1}{e_1} + \frac{\bar{h}_1 e_3}{e_1} + \frac{e_3 e_5}{\bar{h}_2^2} + \frac{e_3 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3^2 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3}{\bar{h}_2} + \frac{e_3}{\bar{h}_2 e_1} + \frac{e_3 e_5}{\bar{h}_2} + \frac{e_3 e_5}{\bar{h}_2 e_1}$$

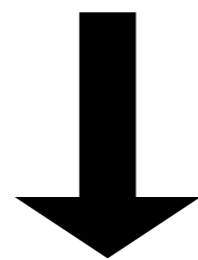
- invariant under Weyl transf.
- contains 10 terms.

Quantum mirror map

$$\Pi_A(z, \hbar) \sim \oint \frac{\log P[X]}{X} dX = Ez^{-1} + \left(-\frac{E^2}{2} - A_2 \right) z^{-2} + \dots$$

$$\left[\begin{array}{l} P[X] = \frac{\Psi[q^{-1}X]}{\Psi[X]}, \quad \hat{Q}\Psi[X] = X\Psi[X] \quad \left[\frac{\hat{H}}{\alpha} + \frac{z}{\alpha} \right] \Psi[X] = 0 \\ \text{solve Schrödinger eq. order by order} \end{array} \right]$$

$$\frac{A_2}{\alpha^2 q^{-1}} = \frac{e_3}{e_1} + \frac{\bar{h}_1}{e_1} + \frac{\bar{h}_1 e_3}{e_1} + \frac{e_3 e_5}{\bar{h}_2^2} + \frac{e_3 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3^2 e_5}{\bar{h}_1 \bar{h}_2^2} + \frac{e_3}{\bar{h}_2} + \frac{e_3}{\bar{h}_2 e_1} + \frac{e_3 e_5}{\bar{h}_2} + \frac{e_3 e_5}{\bar{h}_2 e_1}$$



$$\alpha = q^{1/2} \bar{h}_2^{1/2} e_1^{1/4} e_3^{-1/2} e_5^{-1/4}$$

$$A_2 = \chi_{10}$$

Plan

A-period is similar to **B-period**



- expressed by Weyl characters

②



- coef. of characters = Integer
- has same reps. as those in **B-period**

Multi-covering & BPS

$$\Pi_A(z, \hbar) \sim \sum_{l=1}^{\infty} (-1)^{l+1} A_l z^{-l} \quad A_4 \ni \frac{3}{2} \chi_{54} + \frac{5}{2} \chi_{45} + \frac{11}{2} \chi_1$$

fractional

➔ Does this have multi-covering structure ?

- A-period of A_1 geometry has **multi-covering structure**.
(ABJM theory)
- By this structure, coefficients are **integers**.

[Hatsuda-Marino-Moriyama-Okuyama(2013)]

multi-covering structure

$$(\text{coef. of } n\text{-th order}) = \sum_{j \leq n} (\text{coef. of } j\text{-th order})$$

Multi-covering & BPS

$$\Pi_A(z, \hbar) =: \log z_{\text{eff}} - \log z$$

$$\begin{aligned} \longrightarrow \log z &\sim -A_1 z_{\text{eff}}^{-1} + (A_2 - A_1^2) z_{\text{eff}}^{-2} - \left(A_3 - 3A_1 A_2 + \frac{3}{2} A_1^3 \right) z_{\text{eff}}^{-3} \\ &+ \left(A_4 - 2A_2^2 - 4A_1 A_3 - 8A_1^2 A_2 - \frac{8}{3} A_1^4 \right) z_{\text{eff}}^{-4} + \dots \end{aligned}$$

consider inverse function

$$\begin{cases} A_1 = 0 \\ A_2 = \chi_{10} \\ A_3 = (q^{1/2} + q^{-1/2}) \chi_{16} \\ A_4 = (q^2 + q^{-2}) \chi_1 + (q^{3/2} + q^{-3/2}) (\chi_{45} + 3\chi_1) + \frac{3\chi_{54} + 5\chi_{45} + 11\chi_1}{2} \end{cases}$$

still fractional

Multi-covering & BPS

$$\Pi_A(z, \hbar) =: \log z_{\text{eff}} - \log z$$

$$\begin{aligned} \longrightarrow \log z \sim & \underbrace{-A_1 z_{\text{eff}}^{-1}} + \underbrace{(A_2 - A_1^2) z_{\text{eff}}^{-2}} - \underbrace{\left(A_3 - 3A_1 A_2 + \frac{3}{2} A_1^3 \right) z_{\text{eff}}^{-3}} \\ & + \underbrace{\left(A_4 - 2A_2^2 - 4A_1 A_3 - 8A_1^2 A_2 - \frac{8}{3} A_1^4 \right) z_{\text{eff}}^{-4}} + \dots \end{aligned}$$

$$= \epsilon_1(q, e_i, \bar{h}_i)$$

$$= -\epsilon_2(q, e_i, \bar{h}_i) + \frac{\epsilon_1(q^2, e_i^2, \bar{h}_i^2)}{2}$$

$$= \epsilon_3(q, e_i, \bar{h}_i) + \frac{\epsilon_1(q^3, e_i^3, \bar{h}_i^3)}{3}$$

$$= -\epsilon_4(q, e_i, \bar{h}_i) + \frac{\epsilon_2(q^2, e_i^2, \bar{h}_i^2)}{2} + \frac{\epsilon_1(q^4, e_i^4, \bar{h}_i^4)}{4}$$

multi-covering
structure

$$\epsilon_1 = 0, \quad -\epsilon_2 = \chi_{10},$$

$$\epsilon_3 = \left(q^{1/2} + q^{-1/2} \right) \chi_{16}, \quad -\epsilon_4 = \left(q^2 + q^{-2} \right) \chi_1 + \left(q + q^{-1} \right) \left(\chi_{45} + 3\chi_1 \right) + 4\chi_1$$

integer !

Multi-covering & BPS

[Huang-Klemm-Poretschkin(2013)]

[Moriyama-Nosaka-Yano(2017)]

d	(j_L, j_R)	BPS	$(-1)^{d-1} \sum_{ d =1} (N_{j_L, j_R}^d)_{d^+ - d^-}$	representations
1	(0; 0)	16	$\delta_{+1} + \delta_{-1}$	16
2	$(0, \frac{1}{2})$	10	$1_{+2} + \delta_0 + 1_{-2}$	10
3	(0, 1)	16	$\delta_{+1} + \delta_{-1}$	16
4	$(0, \frac{1}{2})$	1	1_0	1
	$(0, \frac{3}{2})$	45	$\delta_{+2} + 29_0 + \delta_{-2}$	45
	$(\frac{1}{2}, 2)$	1	1_0	1

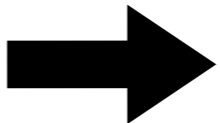
A-period

$$\left[\begin{array}{l}
 \epsilon_1 = 0, \\
 -\epsilon_2 = \underline{\chi_{10}}, \\
 \epsilon_3 = (q^{\frac{1}{2}} + q^{-\frac{1}{2}}) \underline{\chi_{16}}, \\
 -\epsilon_4 = (q^2 + q^{-2}) \underline{\chi_1} + (q + q^{-1}) (\underline{\chi_{45}} + 3\underline{\chi_1}) + \underline{4\chi_1},
 \end{array} \right.$$

B-period

same reps. as B-period for each order

Physical meaning ??



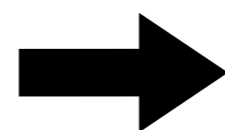
Future work

Summary

A-period is similar to **B-period**

- expressed by Weyl characters
- coef. of characters = Integer
- has same reps. as those in **B-period**

Physical meaning ??



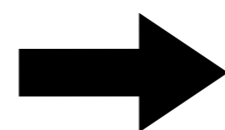
Future work

Summary

A-period is similar to **B-period**

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Physical meaning ??



Future work

Appendix

before using SU(2) character

$$\begin{aligned}\epsilon_7 = & (q^{\frac{11}{2}} + q^{-\frac{11}{2}})\chi_{16} + (q^{\frac{9}{2}} + q^{-\frac{9}{2}})(\chi_{144} + 4\chi_{16}) + (q^{\frac{7}{2}} + q^{-\frac{7}{2}})(\chi_{560} + 4\chi_{144} + 13\chi_{16}) \\ & + (q^{\frac{5}{2}} + q^{-\frac{5}{2}})(\chi_{720} + 4\chi_{560} + 9\chi_{144} + 25\chi_{16}) + (q^{\frac{3}{2}} + q^{-\frac{3}{2}})(3\chi_{560} + 8\chi_{144} + 27\chi_{16}) \\ & + (q^{\frac{1}{2}} + q^{-\frac{1}{2}})(\chi_{720} + 3\chi_{560} + 9\chi_{144} + 27\chi_{16}),\end{aligned}$$

after using SU(2) character

$$\begin{aligned}\epsilon_7 = & \chi_{\frac{11}{2}}\chi_{16} + \chi_{\frac{9}{2}}(\chi_{144} + 3\chi_{16}) + \chi_{\frac{7}{2}}(\chi_{560} + 3\chi_{144} + 9\chi_{16}) \\ & + \chi_{\frac{5}{2}}(\chi_{720} + 3\chi_{560} + 5\chi_{144} + 12\chi_{16}) + \chi_{\frac{3}{2}}(-\chi_{720} - \chi_{560} - \chi_{144} + 2\chi_{16}) \\ & + \chi_{\frac{1}{2}}(\chi_{720} + \chi_{144}),\end{aligned}$$

χ_n SU(2) character

much simpler for q-dependent term & coef.

but negative coef,

without overall minus sign

do not appear

$$\epsilon_4 = (q^2 + q^{-2}) \chi_1 + (q + q^{-1}) (\chi_{45} + 3\chi_1) + (-\chi_{54} + \chi_{45} + 3\chi_1)$$

with overall minus sign

$$-\epsilon_4 = (q^2 + q^{-2})\chi_1 + (q + q^{-1})(\chi_{45} + 3\chi_1) + 4\chi_1$$

4	$(0, \frac{1}{2})$	1	1_0	1
	$(0, \frac{3}{2})$	45	$8_{+2} + 29_0 + 8_{-2}$	45
	$(\frac{1}{2}, 2)$	1	1_0	1

multi-covering structure

$$(\text{coef. of } n\text{-th order}) = \sum_{j \leq n} (\text{coef. of } j\text{-th order})$$

- $(e_1, e_3, e_5, \bar{h}_1, \bar{h}_2) = (1, 1, 1, q, q^{-1})$ (2,2) model
 $= (q^{-1/2}, q^{1/2}, q^{-1/2}, q, q^{-1})$ (1,1,1,1) model

$$\mu = \log z, \quad \mu_{\text{eff}} = \log z_{\text{eff}}$$

(effective) chemical pot.

= log [(effective) complex modulus]

$$\Xi_k(z) = \sum_{N=0}^{\infty} z^N \underline{Z_k(N)} = \text{Det} \left[1 + z \hat{H}^{-1} \right],$$

Partition fn. of N M2-branes

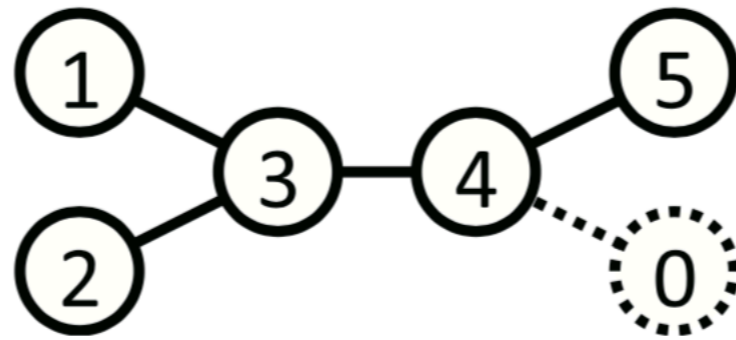
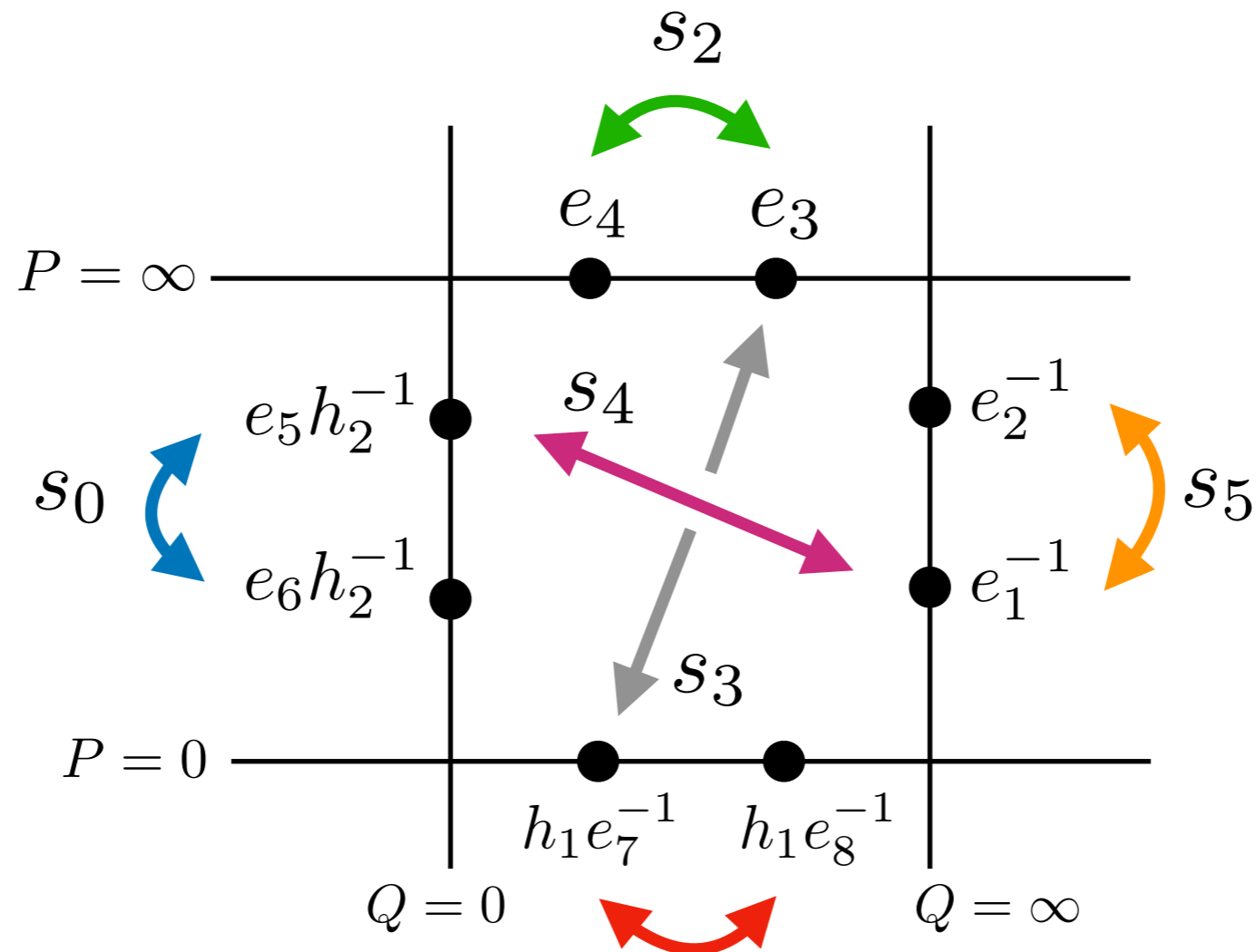


Figure 3: The Dynkin diagram of the D_5 algebra.

from [Kubo,Moriyama,Nosaka(2018)]



$$\begin{aligned}
s_1 : (\bar{h}_1, \bar{h}_2, e_1, e_3, e_5; \alpha) &\mapsto \left(\frac{e_1 e_3 e_5}{\bar{h}_1 \bar{h}_2}, \bar{h}_2, e_1, e_3, e_5; \alpha \right), \\
s_2 : (\bar{h}_1, \bar{h}_2, e_1, e_3, e_5; \alpha) &\mapsto \left(\frac{\bar{h}_1}{e_3}, h_2, e_1, \frac{1}{e_3}, e_5; e_3 \alpha \right), \\
s_3 : (\bar{h}_1, \bar{h}_2, e_1, e_3, e_5; \alpha) &\mapsto \left(\bar{h}_1, \frac{e_1 e_5}{\bar{h}_1 \bar{h}_2}, e_1, \frac{e_1 e_3 e_5}{\bar{h}_1 \bar{h}_2}, e_5; \alpha \right), \\
s_4 : (\bar{h}_1, \bar{h}_2, e_1, e_3, e_5; \alpha) &\mapsto \left(\frac{\bar{h}_1 \bar{h}_2}{e_1 e_5}, \bar{h}_2, \frac{\bar{h}_2}{e_5}, e_3, \frac{\bar{h}_2}{e_1}; \alpha \right), \\
s_5 : (\bar{h}_1, \bar{h}_2, e_1, e_3, e_5; \alpha) &\mapsto \left(\bar{h}_1, \frac{\bar{h}_2}{e_1}, \frac{1}{e_1}, e_3, e_5; \frac{\alpha}{e_1} \right),
\end{aligned}$$

Susy. Gauge Theory

e.g. Free energy of ABJM

$$F = \underbrace{\begin{array}{l} \text{(Worldsheet inst.)} \\ \text{poles} \\ + \text{(bound state)} \\ \text{poles} \end{array}}_{\text{No pole}} + \underbrace{\text{(Membrane inst.)}}_{\text{poles}}$$

e.g. Quantum dilogarithm

$$i \log \Phi \sim \underbrace{\sum_{d=1}^{\infty} \frac{(-1)^d Q^d}{2d \sin\left(\frac{dk}{2}\right)}}_{\text{poles}} + \sum_{d=1}^{\infty} \frac{(-1)^d Q^{2\pi d/k}}{2d \sin\left(\frac{2\pi^2 d}{k}\right)}_{\text{poles}}$$

No pole