

# Gravitational Wave Echoes

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ICTS Workshop, Yichang, 24 November 2019

Zhang & **SYZ**, 1709.07503

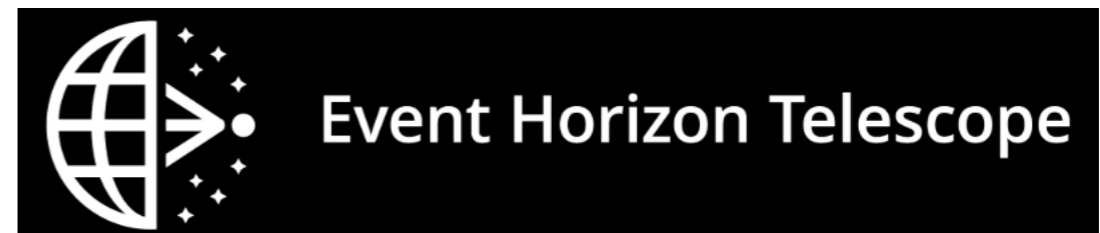
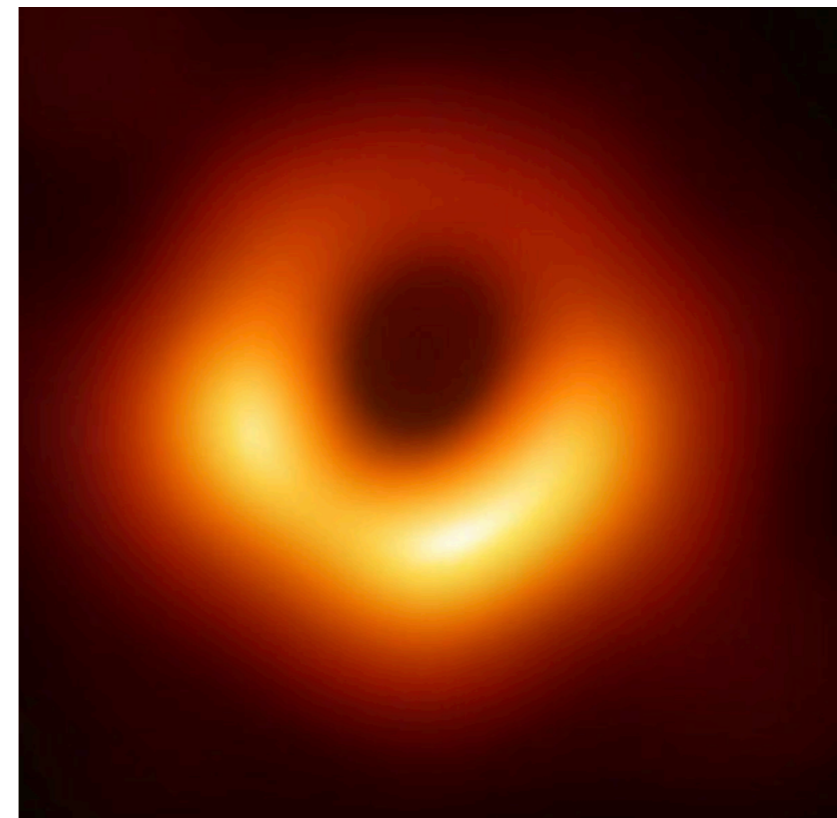
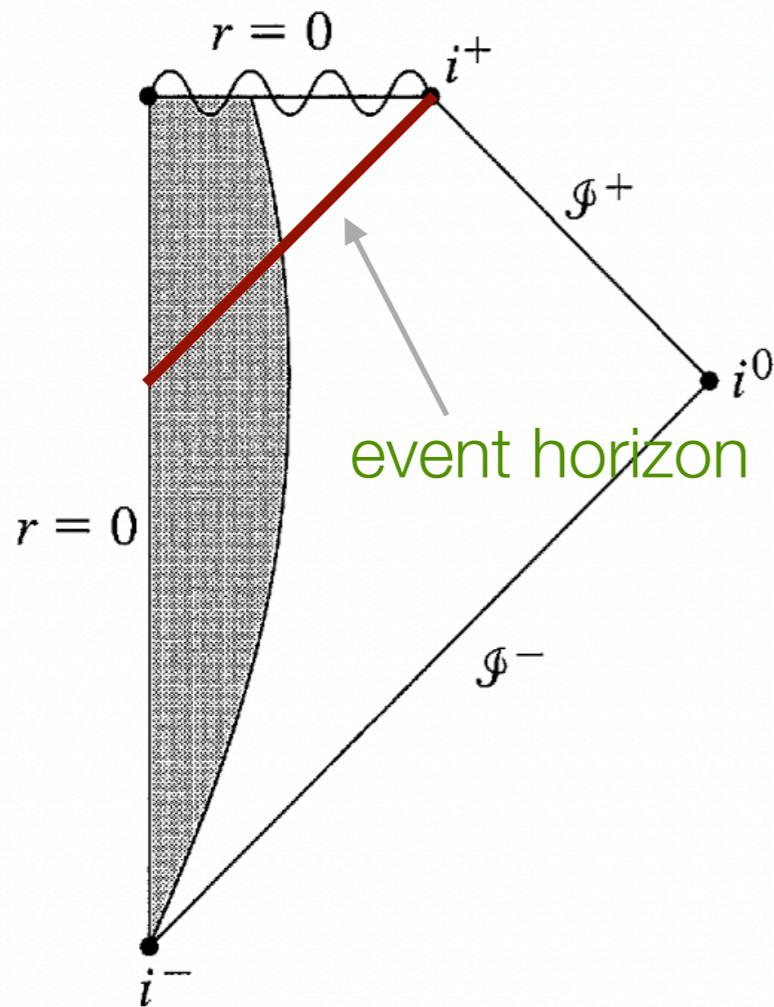
Huang, Xu & **SYZ**, 1908.00189

Wang, Zhang, **SYZ** & Piao, 1802.02003, 1904.00212

# The photo of a black hole

What defines a black hole?

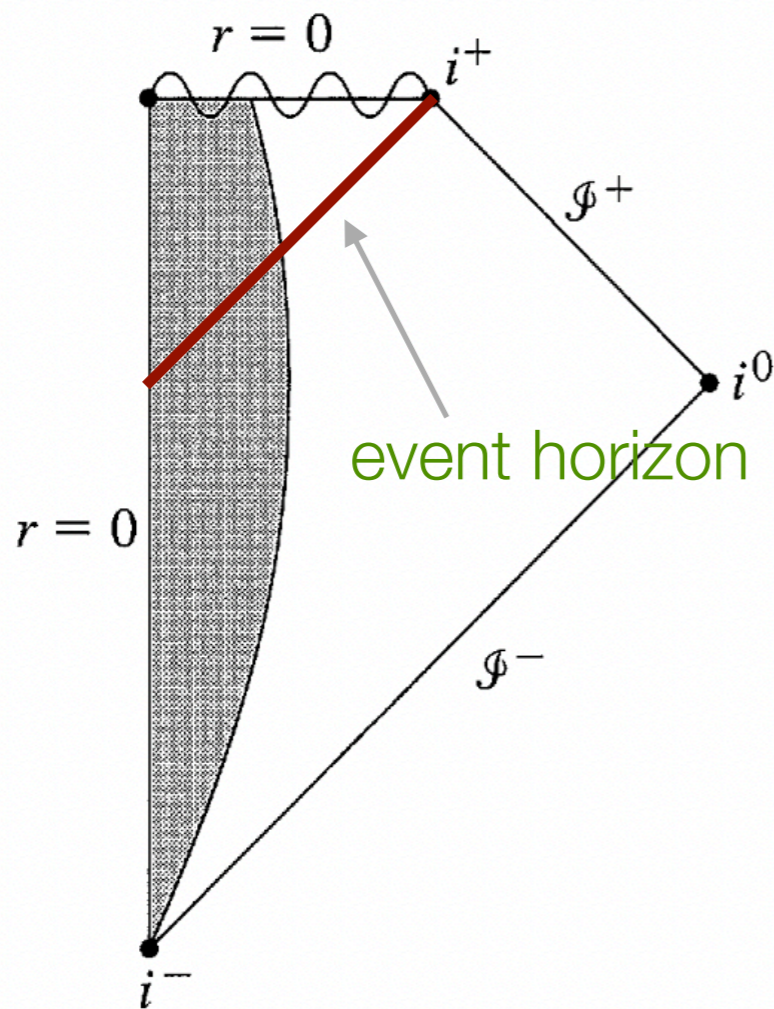
M87\*



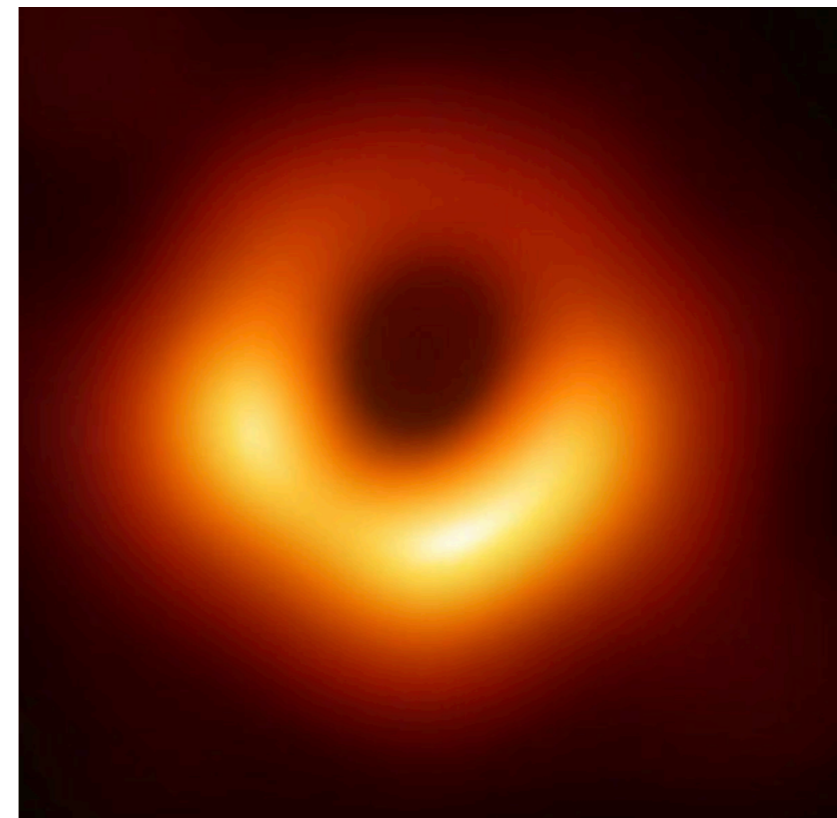
good PR but atrocious, grotesque abuse!

# The photo of a black hole

What defines a black hole?



M87\*



good PR but atrocious, grotesque abuse!

# The sound of a black hole merger

PRL **116**, 061102 (2016) week ending  
12 FEBRUARY 2016

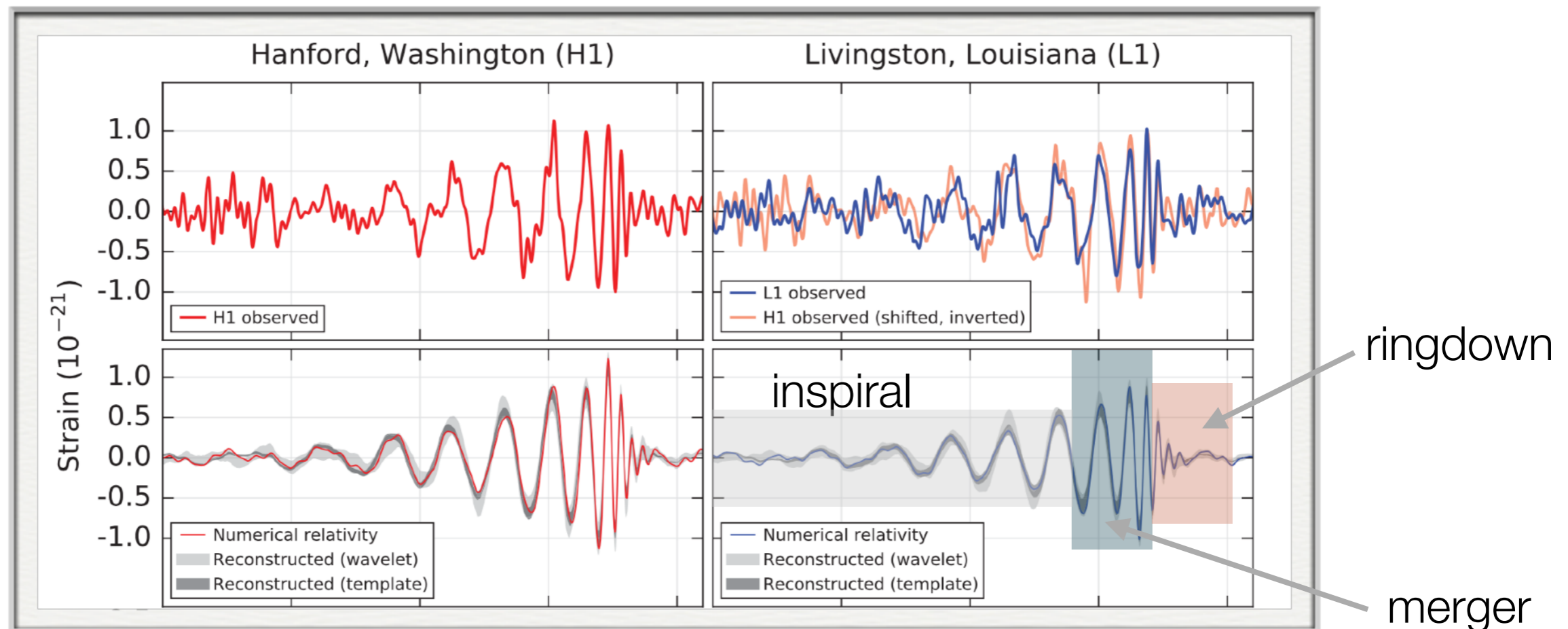
Selected for a **Viewpoint** in *Physics*  
PHYSICAL REVIEW LETTERS

☞

**Observation of Gravitational Waves from a Binary Black Hole Merger**

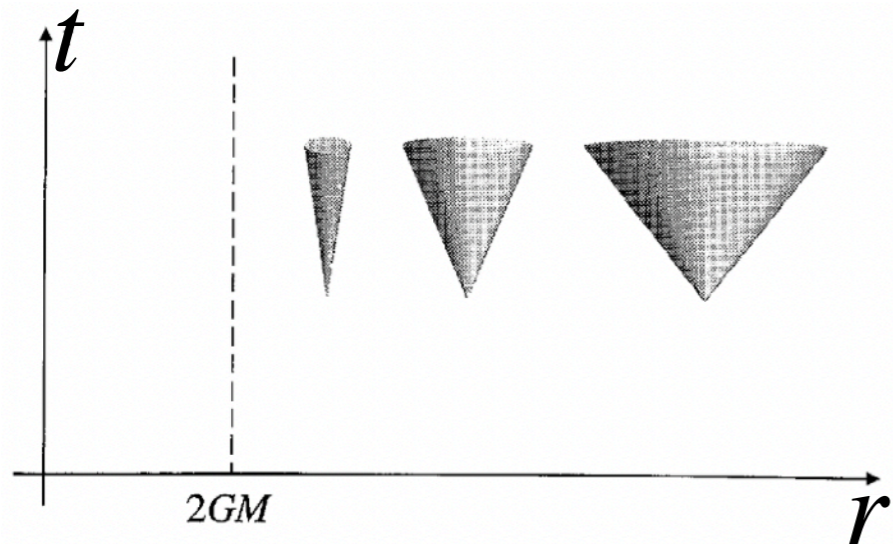
B. P. Abbott *et al.*<sup>\*</sup>  
(LIGO Scientific Collaboration and Virgo Collaboration)  
(Received 21 January 2016; published 11 February 2016)

GW150914



# But have we observed an event horizon?

light cones near horizon

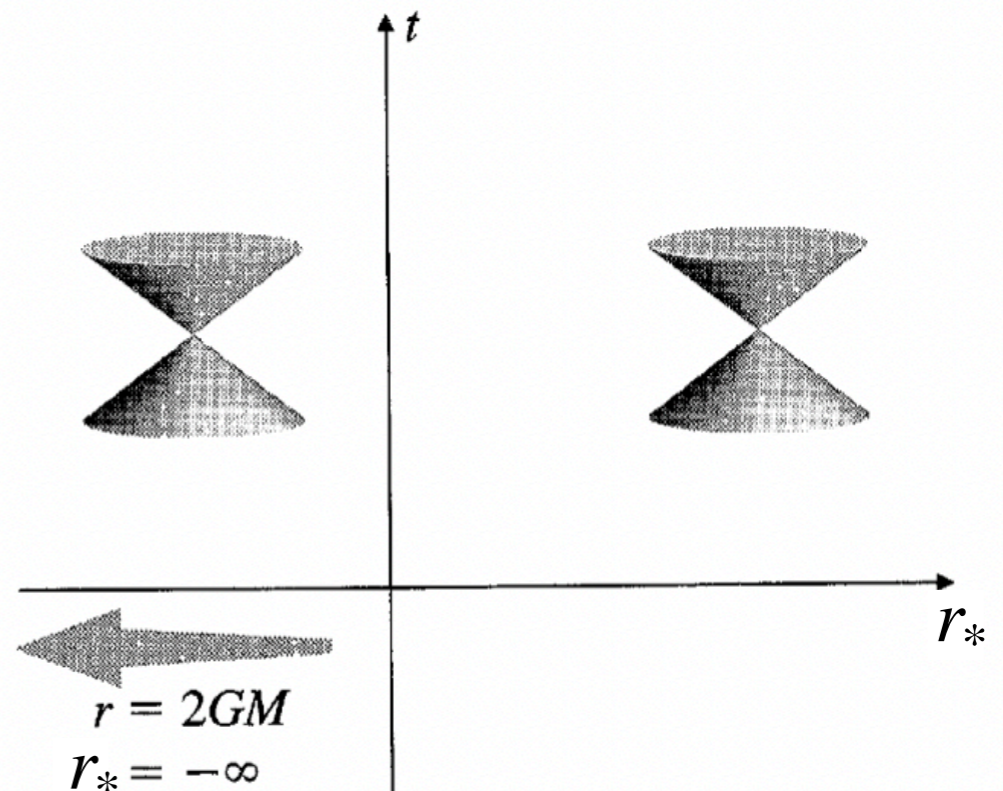
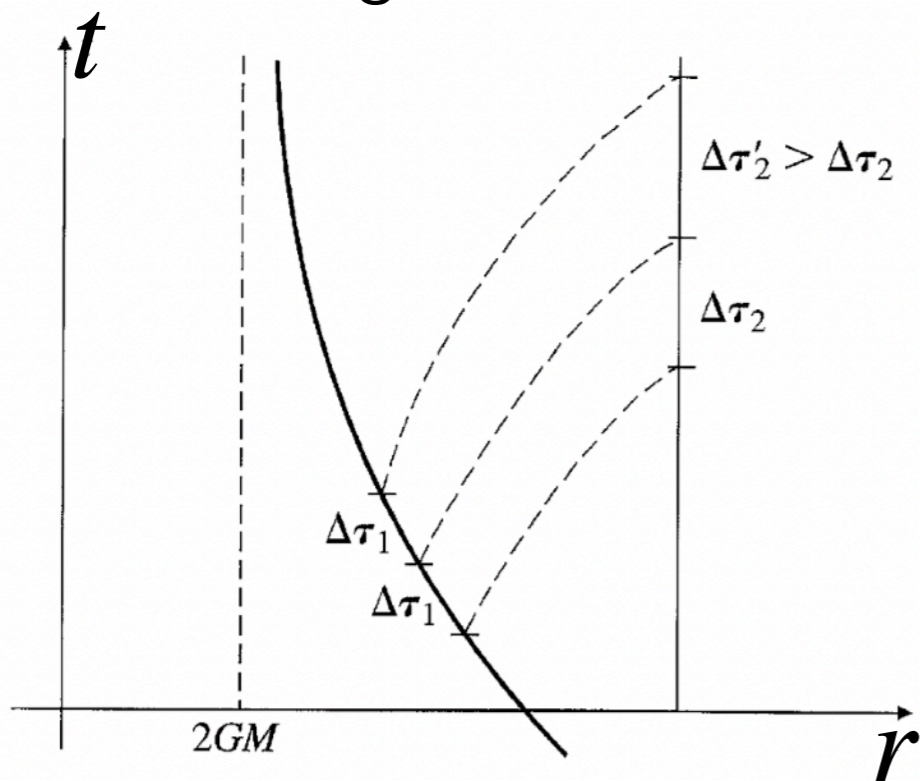


Tortoise coordinate

$$r_* = r + 2GM \ln \left( \frac{r}{2GM} - 1 \right)$$

$$ds^2 = \left( 1 - \frac{2GM}{r} \right) (-dt^2 + dr_*^2) + r^2 d\Omega^2$$

beacon falling into black hole



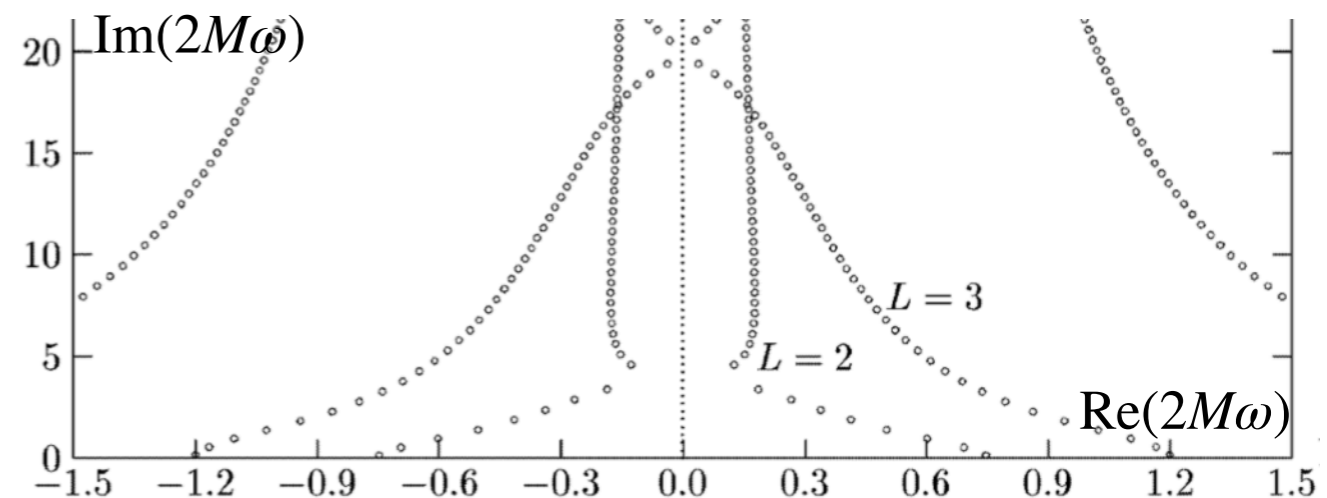
# The magic of quasi-normal modes?

- QNMs of BH: characteristic oscillation modes of BH spacetime

Regge-Wheeler-Zerilli equation

$$\frac{d^2\psi_l(r_*)}{dr_*^2} + [\omega^2 - V_l(r_*)] \psi_l(r) = S_l$$

- damped:  $\omega = \omega_R + i\omega_I$
- ringdown dominated by QNMs
- QNM spectrum determined by BH's  $M$  and  $a$  (Kerr)
- test of no-hair theorem



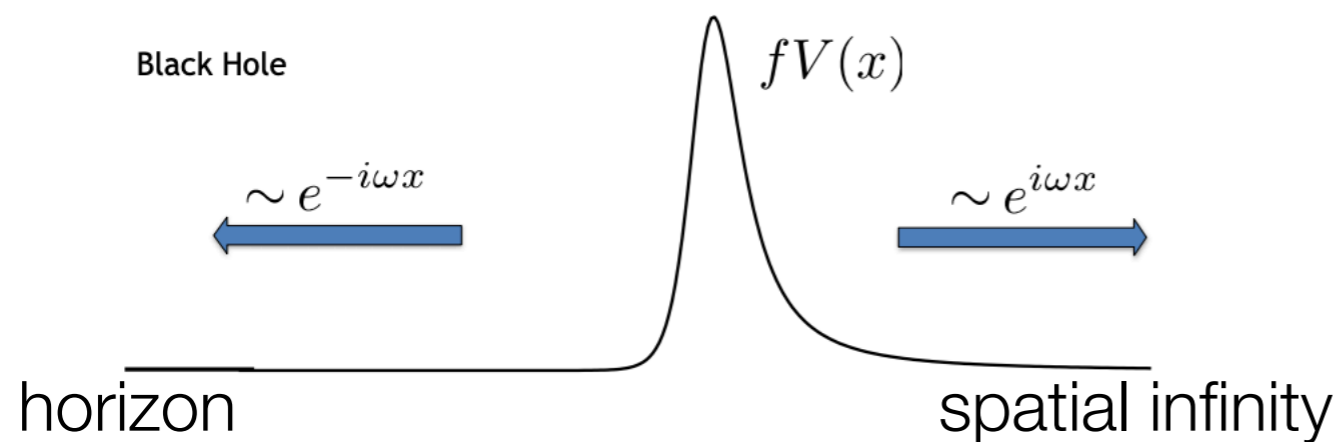
Ringdown: only conclusive test of BH(?)

# Not necessarily!

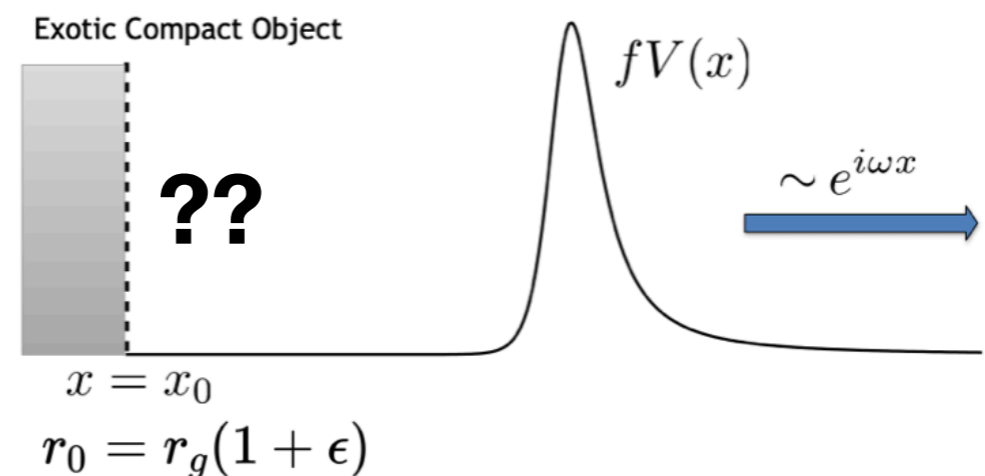
- This is based on assumption: **Ringdown  $\simeq$  QNMs**
- Ringdown simply determined by photosphere

frequency & damping time  $\sim$  circular photo orbit

- QNMs are sensitive to boundary conditions



Ringdown  $\simeq$  QNMs



**Ringdown  $\neq$  QNMs!**

# Exotic compact object (ECO): an example

Featured in Physics

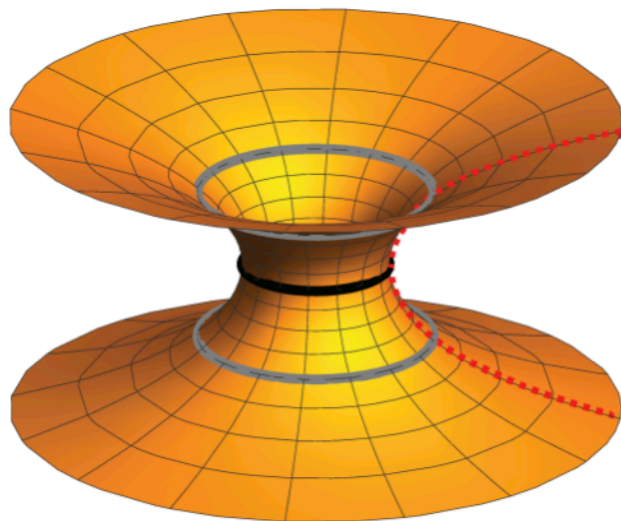
Editors' Suggestion

## Is the Gravitational-Wave Ringdown a Probe of the Event Horizon?

Vitor Cardoso, Edgardo Franzin, and Paolo Pani

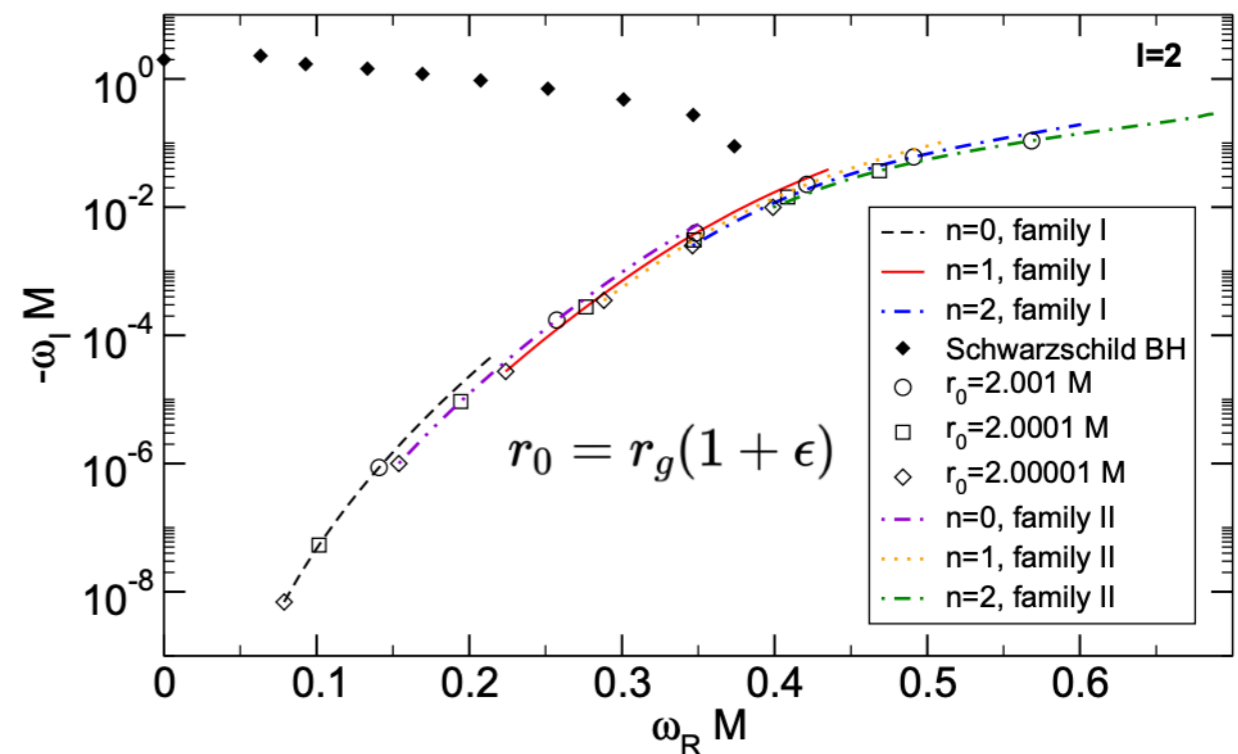
Phys. Rev. Lett. **116**, 171101 – Published 27 April 2016; Erratum [Phys. Rev. Lett. \*\*117\*\*, 089902 \(2016\)](#)

### Schwarzschild Wormhole



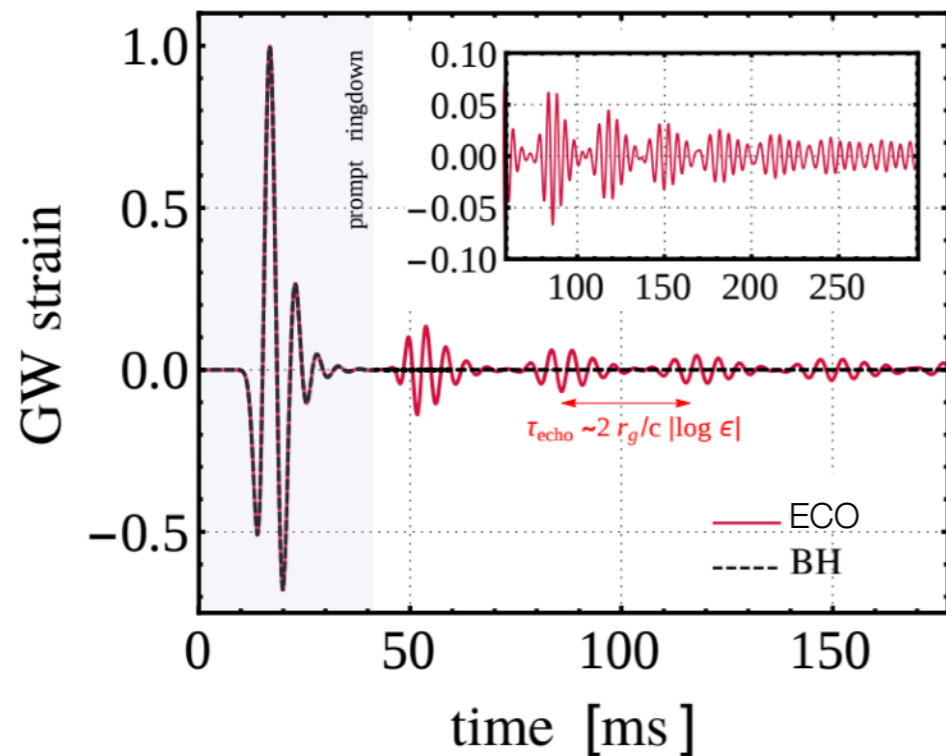
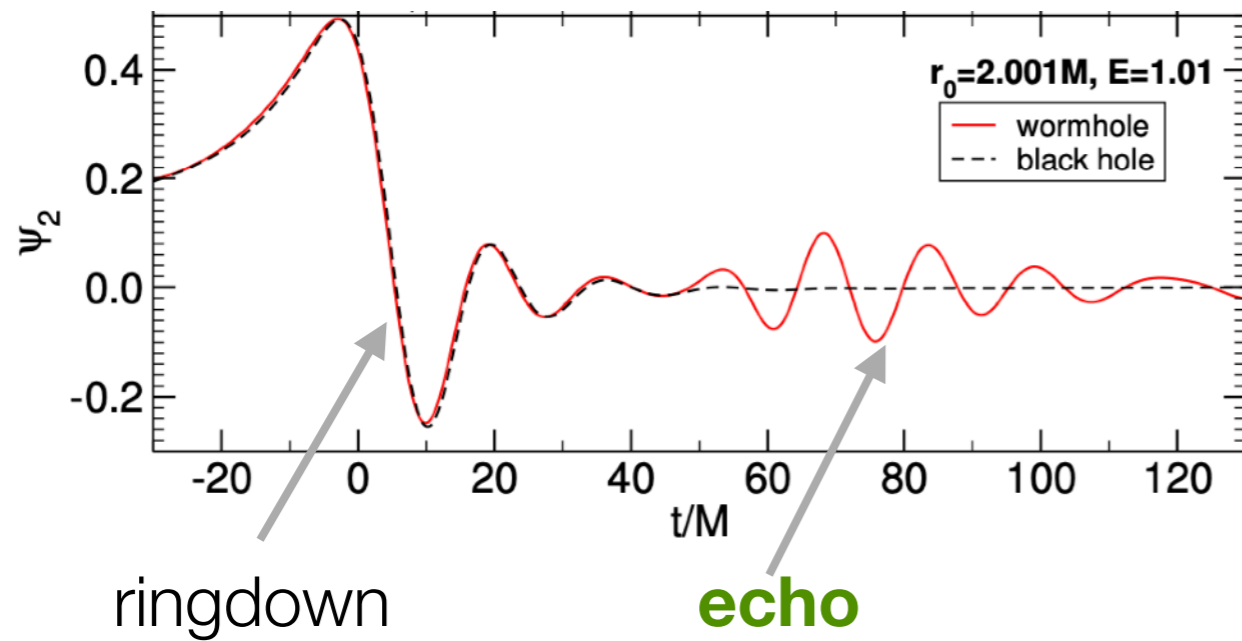
“almost” Schwarzschild BH

### QNMs

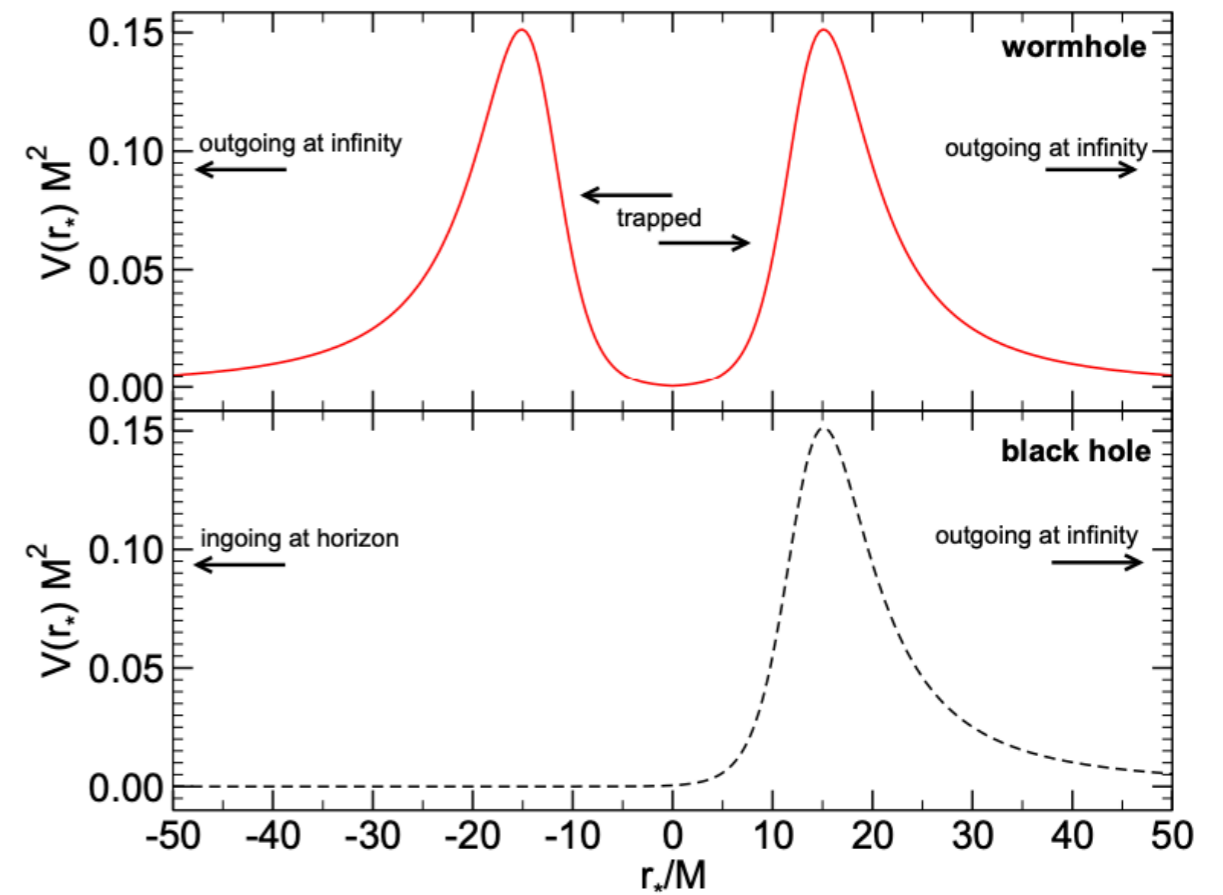




## GW waveform



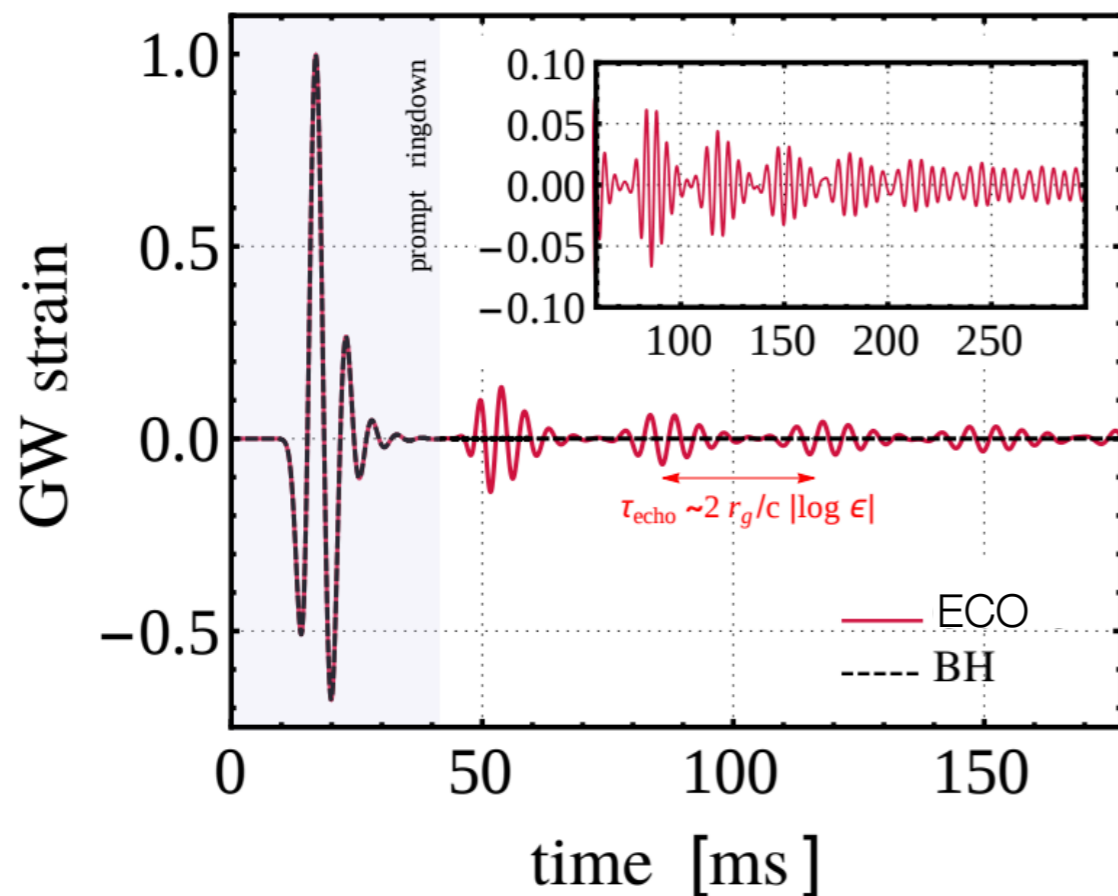
## Effective potential



For an ECO

1. Ringdown is same as BH
2. There are extra echoes

# Smoking gun for new physics: echoes



New physics radius

$$r_0 = r_g(1 + \epsilon)$$

Echo interval

$$\tau_{\text{echo}} \sim \frac{|r_*(1.5r_g) - r_*(r_0)|}{c} \sim \frac{2r_g}{c} |\log(\epsilon)|$$

What kind of new physics?

black hole mimickers: fuzzball, gravastar, boson stars, ...

quantum gravity effects: firewall

Planck length  $\epsilon \sim 10^{-40} \rightarrow \tau_{\text{echo}}(60M_{\odot}) \sim 50 \text{ ms}$

# Echoes from MVMG black hole

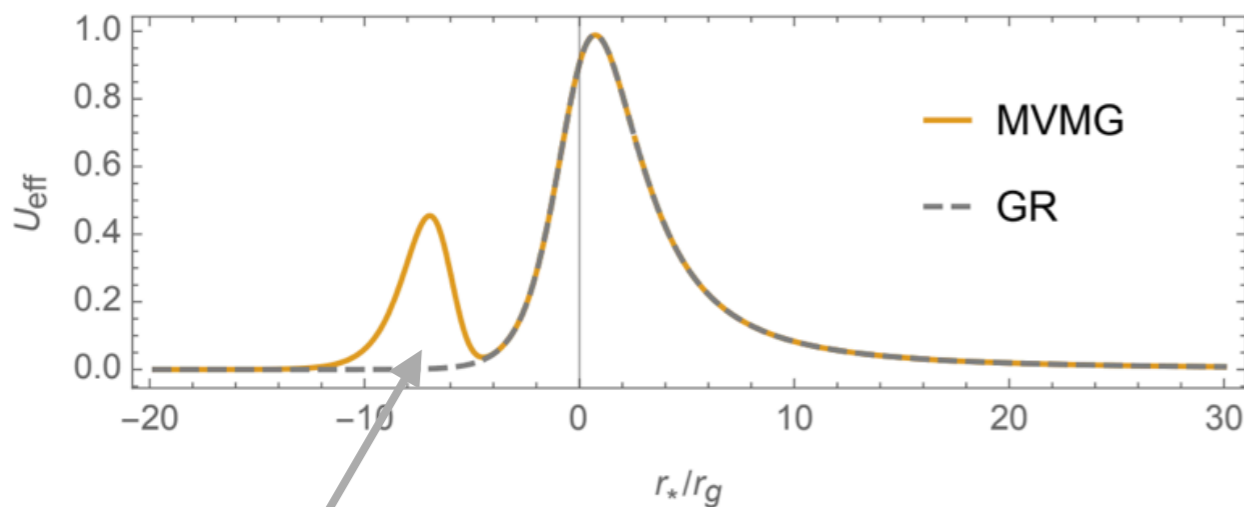
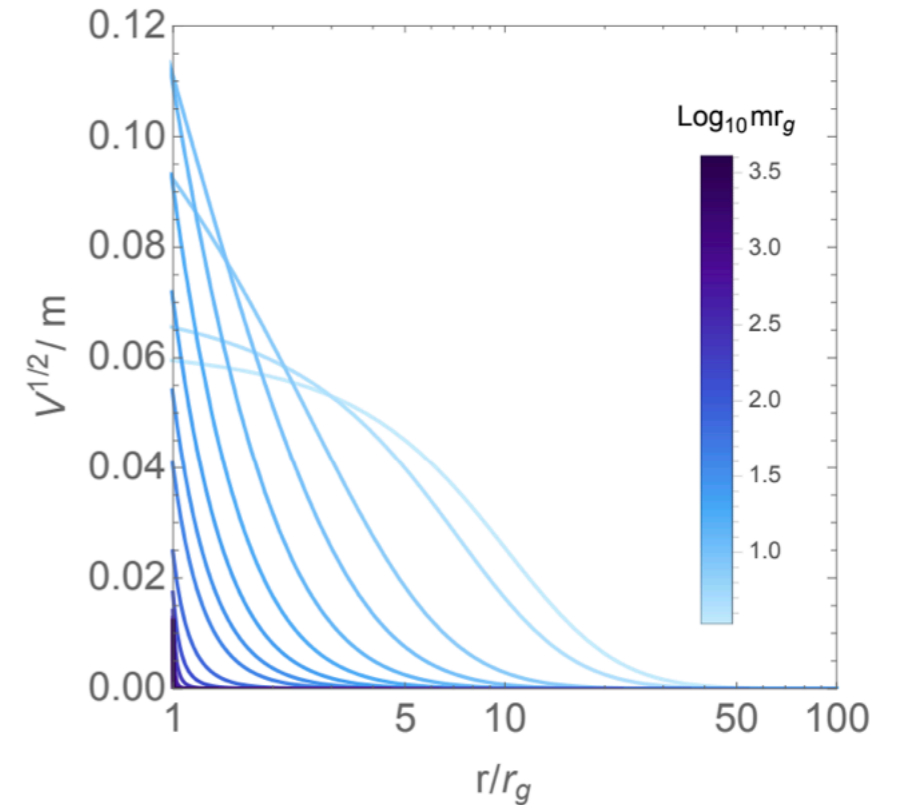
Zhang & SYZ, 1709.07503

## Mass-varying massive gravity

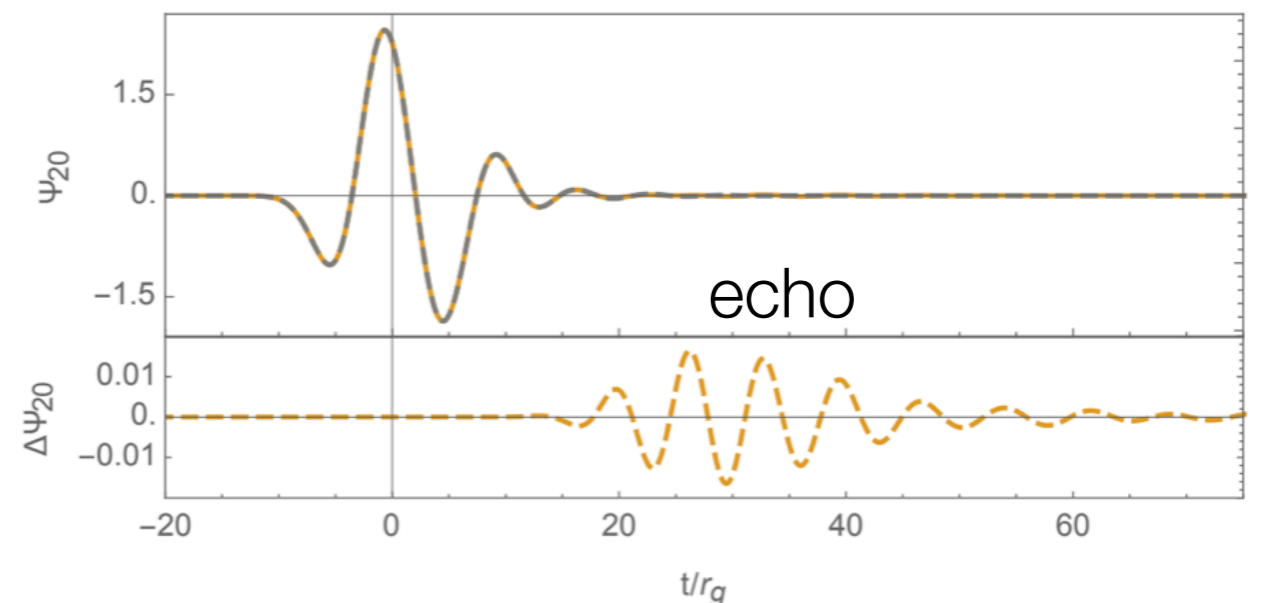
$$S = M_P^2 \int d^4x \sqrt{-g} \left[ \frac{R}{2} + V(\sigma) \mathcal{U} - \frac{1}{2} (\partial\sigma)^2 - W(\sigma) \right]$$

## Hairy black hole

graviton mass becomes very big  
close to the horizon



graviton mass barrier



# How to characterize GW echoes?

Mark, Zimmerman, Du & Chen, 1709.07503

Parametrize new physics with a reflective boundary

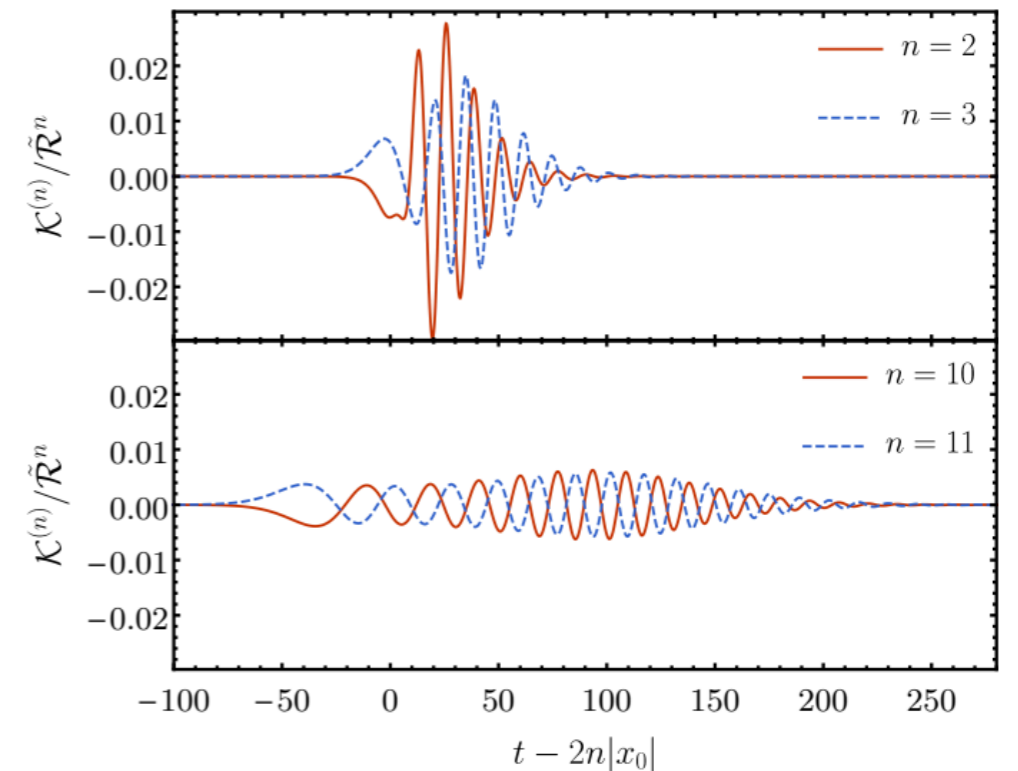
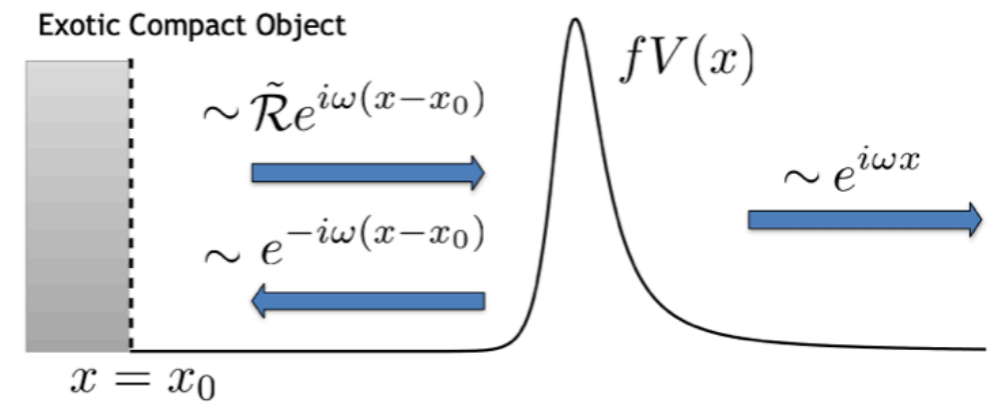
$$\psi \propto e^{-i\omega(x-x_0)} + \tilde{\mathcal{R}}(\omega)e^{i\omega(x-x_0)}, \quad x \rightarrow x_0$$

$$\psi(x) = \int_{-\infty}^{\infty} dx' \tilde{g}_{\text{ref}}(x, x') \tilde{S}(x')$$

$$\tilde{g}_{\text{ref}}(x, x') = \underbrace{\tilde{g}_{\text{BH}}(x, x')}_{\text{GR part}} + \underbrace{\tilde{\mathcal{K}} \frac{\tilde{\psi}_{\text{up}}(x)\tilde{\psi}_{\text{up}}(x')}{W_{\text{BH}}}}_{\text{echo part}}$$

$$\tilde{\mathcal{K}} = \tilde{\mathcal{T}}_{\text{BH}} \tilde{\mathcal{R}} e^{-2i\omega x_0} \sum_{n=1}^{\infty} (\tilde{\mathcal{R}}_{\text{BH}} \tilde{\mathcal{R}})^{(n-1)} e^{-2i(n-1)\omega x_0}$$

$n$ -th term corresponds to  $n$ -th echo



identity individual echoes

# The Fredholm approach to echoes

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Huang, Xu & **SYZ**, 1908.00189

Regge-Wheeler-Zerilli equation

$$\phi''(\omega, x) + (\omega^2 - V(x))\phi(\omega, x) = \mathcal{I}(\omega, x)$$

Fredholm integral equation

$$\phi(x) = f(x) + \lambda \int dy K(x, y)\phi(y)$$

$$K(x, y) = G(x, y)V(y) \quad \left(\frac{d^2}{dx^2} + \omega^2\right)G(x, y) = \delta(x - y) \quad f(x) = \int dy G(x, y)\mathcal{I}(y)$$

Separate kernel function

$$K(x, y) = \sum_{j=1}^{\infty} \alpha_j(x)\beta_j(y)$$

# Solution

$$\phi(x) = f(x) - \lambda \int dy f(y) \frac{\Delta_K(x, y, \lambda)}{\Delta(\lambda)}$$

$$\Delta(\lambda) = \det(c)$$

$$c_{ij} = \delta_{ij} - \lambda \int dx \beta_i(x) \alpha_j(x)$$

$$\Delta_K(x, y, \lambda) = \begin{vmatrix} 0 & \alpha_1(x) & \alpha_2(x) & \cdots \\ \beta_1(y) & c_{11} & c_{12} & \cdots \\ \beta_2(y) & c_{21} & c_{22} & \cdots \\ \cdots & \cdots & \cdots & \ddots \end{vmatrix}$$

Example: BH boundary condition with potential  $V(x)$

$$G(x, y) = \frac{e^{i\omega|x-y|}}{2i\omega}$$

Observer at large  $x$

$$K(x, y) = \alpha(x)\beta(y) \quad \alpha(x) = \frac{e^{i\omega x}}{2i\omega}, \quad \beta(y) = e^{-i\omega y} V(y).$$

$$\phi(x) = f(x) + \int dx_1 \frac{e^{i\omega(x-x_1)}}{2i\omega} \frac{V(x_1)f(x_1)}{1 - \frac{1}{2i\omega} \int dy V(y)}$$

# Perturbation scheme

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If separation is difficult, expand in  $K(x, y)$

$$\phi(x) = f(x) + \lambda \int dy f(y) \frac{\sum_{n=0}^{\infty} \frac{(-\lambda)^n}{n!} A_n(x, y)}{\sum_{n=0}^{\infty} \frac{(-\lambda)^n}{n!} D_n}$$

$$A_n(x, y) = \int \prod_{i=1}^n dx_i \begin{vmatrix} K(x, y) & K(x, x_1) & K(x, x_2) & \cdots & K(x, x_n) \\ K(x_1, y) & K(x_1, x_1) & K(x_1, x_2) & \cdots & K(x_1, x_n) \\ \cdots & \cdots & \cdots & \ddots & \cdots \\ K(x_n, y) & K(x_n, x_1) & K(x_n, x_2) & \cdots & K(x_n, x_n) \end{vmatrix}$$

$$D_n = \int \prod_{i=1}^n dx_i \begin{vmatrix} K(x_1, x_1) & K(x_1, x_2) & K(x_1, x_3) & \cdots & K(x_1, x_n) \\ K(x_2, x_1) & K(x_2, x_2) & K(x_2, x_3) & \cdots & K(x_2, x_n) \\ \cdots & \cdots & \cdots & \ddots & \cdots \\ K(x_n, x_1) & K(x_n, x_2) & K(x_n, x_3) & \cdots & K(x_n, x_n) \end{vmatrix}$$

# Convergence and error estimate

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- Approximate solution

Huang, Xu & **SYZ**, 1908.00189

$$\phi_{\text{app}}(x) = f(x) + \int_a^b \frac{\sum_{n=0}^N \frac{(-\lambda)^n}{n!} A_n(x, y)}{\sum_{n=0}^N \frac{(-\lambda)^n}{n!} D_n} f(y) dy.$$

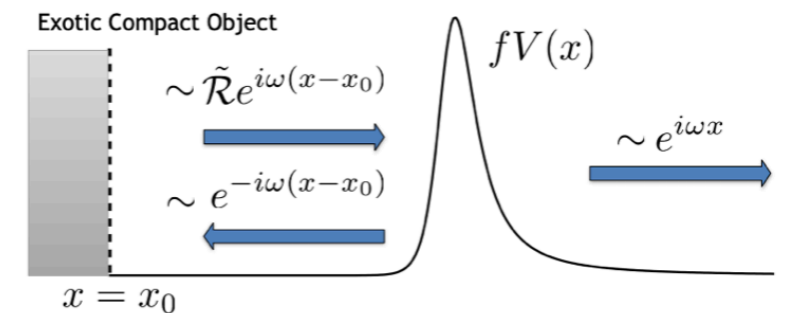
- absolute convergent if  $V(x)$  is normalizable
- For Schwarzschild potential
  - estimated errors from truncating  $N$
  - estimated errors from neglecting  $(-L, a) \cup (b, +\infty)$



# Echoes in Fredholm formalism

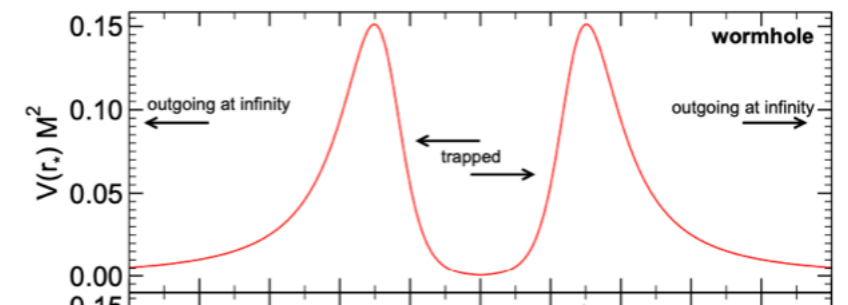
## 1. Reflective “mirror”

$$G(x, y) = \frac{e^{i\omega|x-y|}}{2i\omega} + R(\omega) \frac{e^{i\omega(x+y)}}{2i\omega}$$



## 2. Extra potential barrier

$$V(x) = \bar{V}(x) + R(\omega)V_Q(x)$$



For either case

$$K(x, y) = \bar{K}(x, y) + R(\omega)Q(x, y)$$

GR kernel

echo kernel

$$\phi(x) = f(x) + \lambda \int dy f(y) \frac{\sum_{n=0}^{\infty} \frac{(-\lambda)^n}{n!} A_n(x, y)}{\sum_{n=0}^{\infty} \frac{(-\lambda)^n}{n!} D_n}$$

$$D_n = \sum_{j=0}^n R^j D_n^j, \quad A_n(x, y) = \sum_{j=0}^{n+1} R^j A_n^j(x, y)$$

$$D_n^j = \sum_{1 \leq i_1 < \dots < i_j \leq n} \int \prod_{i=1}^n dx_i \begin{vmatrix} \bar{K}(x_1, x_1) & \dots & Q(x_1, x_{i_1}) & \dots & Q(x_1, x_{i_j}) & \dots & \bar{K}(x_1, x_n) \\ & \dots & & \dots & & \vdots & \\ \bar{K}(x_n, x_1) & \dots & Q(x_n, x_{i_1}) & \dots & Q(x_n, x_{i_j}) & \dots & \bar{K}(x_n, x_n) \end{vmatrix},$$

$$A_n^j(x, y) = \sum_{1 \leq i_1 < \dots < i_j \leq n} \int \prod_{i=1}^n dx_i \begin{vmatrix} \bar{K}(x, y) & \bar{K}(x, x_1) & \dots & Q(x, x_{i_1}) & \dots & Q(x, x_{i_j}) & \dots & \bar{K}(x, x_n) \\ \bar{K}(x_1, y) & \bar{K}(x_1, x_1) & \dots & Q(x_1, x_{i_1}) & \dots & Q(x_1, x_{i_j}) & \dots & \bar{K}(x_1, x_n) \\ & \dots & & \dots & & & \vdots & \\ \bar{K}(x_n, y) & \bar{K}(x_n, x_1) & \dots & Q(x_n, x_{i_1}) & \dots & Q(x_n, x_{i_j}) & \dots & \bar{K}(x_n, x_n) \end{vmatrix}$$

# Diagrammatica for perturbations

Huang, Xu & SYZ, 1908.00189

$$K(x, y) = \text{---}\ominus\text{---}, \quad \int dx K(x, x) = \text{---}\bigcirc\ominus\text{---}, \quad \int dx K(x_1, x) K(x, x_2) = \text{---}\ominus\text{---}\ominus\text{---}$$

$$D_4 = (\text{---}\bigcirc\ominus\text{---})^4 - 6 \text{---}\bigcirc\ominus\text{---} \text{---}\bigcirc\ominus\text{---} + 3 (\text{---}\bigcirc\ominus\text{---})^2 + 3 (\text{---}\bigcirc\ominus\text{---})^2 + 8 \text{---}\bigcirc\ominus\text{---} \text{---}\bigcirc\ominus\text{---} - 3! \text{---}\bigcirc\ominus\text{---}$$

$$A_3(x, y) = \text{---}\ominus\text{---} (\text{---}\bigcirc\ominus\text{---})^3 - 3 \text{---}\ominus\text{---}\ominus\text{---} (\text{---}\bigcirc\ominus\text{---})^2 - 3 \text{---}\ominus\text{---} \text{---}\bigcirc\ominus\text{---} \text{---}\bigcirc\ominus\text{---} + 3 \text{---}\ominus\text{---}\ominus\text{---} \text{---}\bigcirc\ominus\text{---} \\ + 6 \text{---}\ominus\text{---}\ominus\text{---}\ominus\text{---} \text{---}\bigcirc\ominus\text{---} + 2 \text{---}\ominus\text{---} \text{---}\bigcirc\ominus\text{---} - 6 \text{---}\ominus\text{---}\ominus\text{---}\ominus\text{---}\ominus\text{---}$$

$$\phi(x) = f(x) + \lambda \int dy f(y) \frac{\sum_{n=0}^{\infty} \frac{(-\lambda)^n}{n!} A_n(x, y)}{\sum_{n=0}^{\infty} \frac{(-\lambda)^n}{n!} D_n}$$

$$= f(x) - \lambda \int dy \frac{\text{cut exp}(-\lambda \text{---}\bigcirc\ominus\text{---} - \frac{\lambda^2}{2} \text{---}\bigcirc\ominus\text{---} \text{---}\bigcirc\ominus\text{---} - \dots)}{\text{exp}(-\lambda \text{---}\bigcirc\ominus\text{---} - \frac{\lambda^2}{2} \text{---}\bigcirc\ominus\text{---} \text{---}\bigcirc\ominus\text{---} - \dots)} f(y)$$

For echoes, add  $R(\omega)Q(x, y) = \text{---}\bullet\text{---}$ .

$$D_1 = \text{---}\ominus\text{---} + \text{---}\bullet\text{---},$$

$$D_2 = (\text{---}\ominus\text{---})^2 - \text{---}\ominus\text{---}\text{---}\ominus\text{---}$$

$$+ 2(\text{---}\ominus\text{---}\text{---}\bullet\text{---} - \text{---}\ominus\text{---}\text{---}\bullet\text{---})$$

$$+ ((\text{---}\bullet\text{---})^2 - \text{---}\bullet\text{---}\text{---}\bullet\text{---}),$$

$$A_0(x, y) = \text{cut}D_1 = \text{---}\ominus\text{---} + \text{---}\bullet\text{---},$$

$$A_1(x, y) = \text{cut}D_2 = \text{---}\ominus\text{---} \text{---}\ominus\text{---} - \text{---}\ominus\text{---}\text{---}\ominus\text{---}$$

$$+ (\text{---}\ominus\text{---}\text{---}\bullet\text{---} + \text{---}\bullet\text{---}\text{---}\ominus\text{---} - 2 \text{---}\ominus\text{---}\text{---}\bullet\text{---})$$

$$+ (\text{---}\bullet\text{---}\text{---}\bullet\text{---} - \text{---}\bullet\text{---}\text{---}\bullet\text{---}),$$

**For  $n$ -th echo, collect all vacuum and 2-point diagrams  
with  $n$  solid vertices**

# Diagrammatical rules

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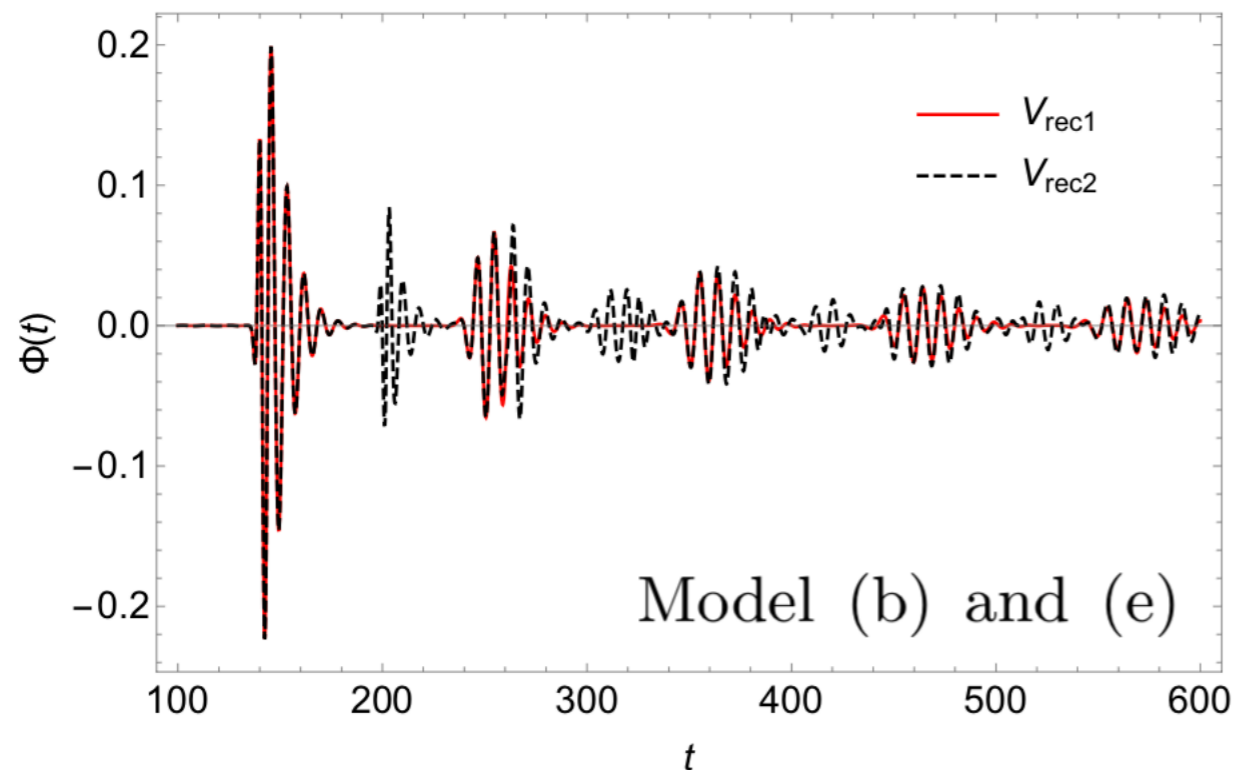
- $\bar{K}(x_i, x_j)$  is represented as a (2-point) circle vertex;
- $R(\omega)Q(x_i, x_j)$  is represented as a (2-point) solid vertex;
- Diagrams with  $n$  solid vertices contribute to the  $n$ -th echo waveform;
- $\int dx$  is represented as a “propagator”;
- For each diagram with an odd number of loops, assign a minus sign;
- $D_n$  includes all loop diagrams with  $n$  vertices;
- $A_n$  is obtained by cutting open one of the loops in  $D_{n+1}$ .

# Numerical results

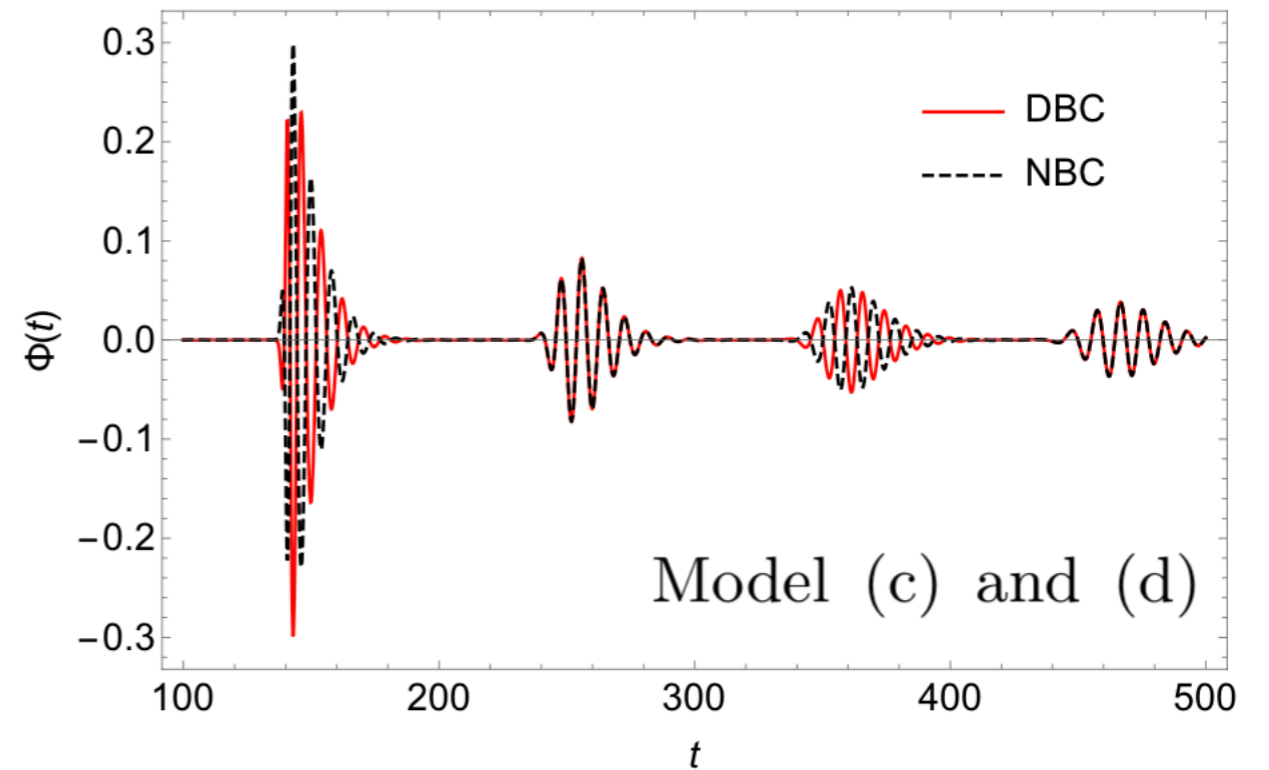
Huang, Xu & **SYZ**, 1908.00189

Model	Potential $V(x)$	Left Boundary
(a)	$V_{\text{gr}}(x)$	open
(b)	$V_{\text{gr}}(x) + V_{\text{rec1}}(x)$	open
(c)	$V_{\text{gr}}(x)$	Dirichlet $x = -20$
(d)	$V_{\text{gr}}(x)$	Neumann $x = -20$
(e)	$V_{\text{gr}}(x) + V_{\text{rec2}}(x)$	open

extra barrier

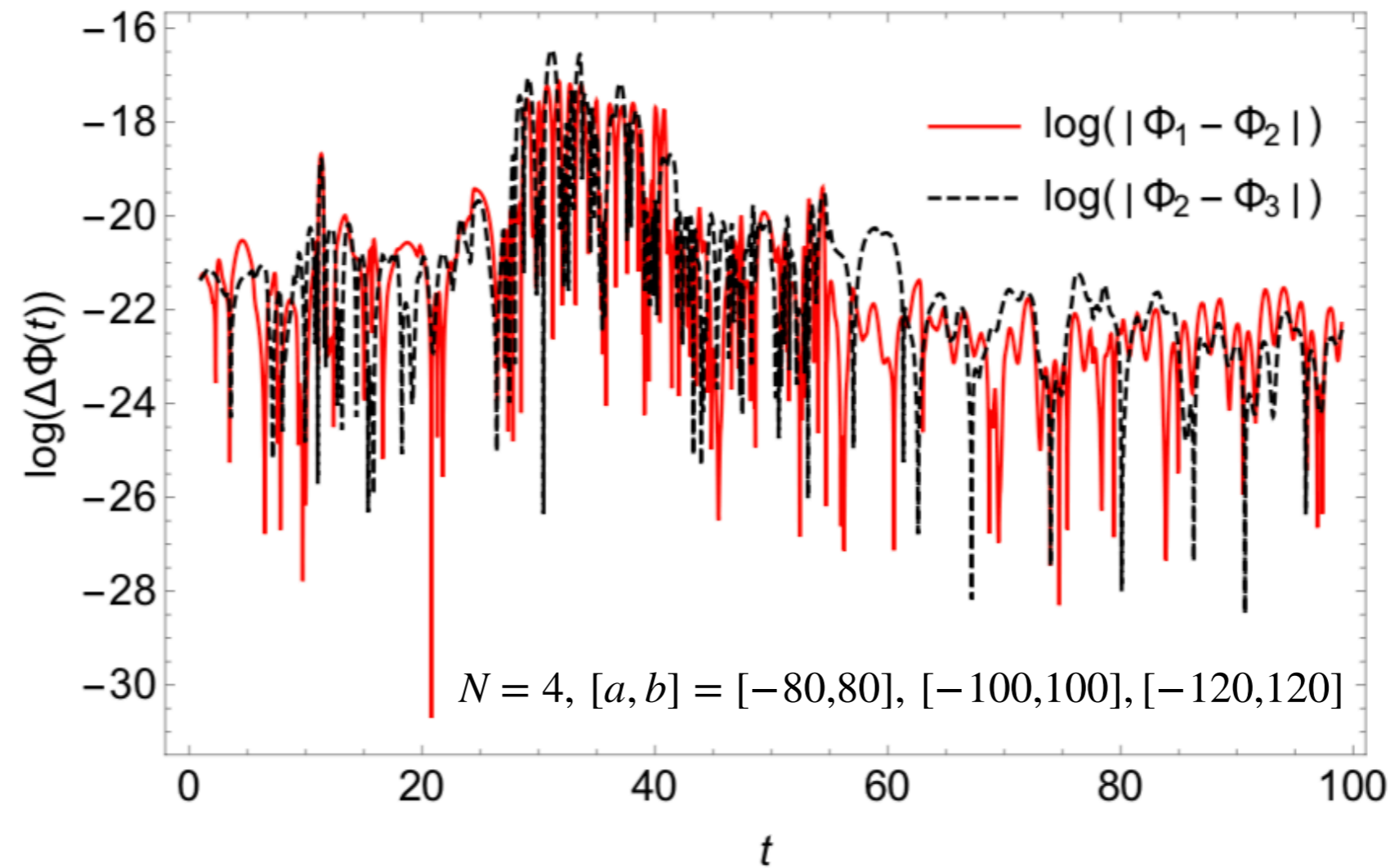


reflective boundary



# Numerical accuracy

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# Echoes from rotating (Kerr-like) ECOs

## Teukolsky master equation

Huang, Xu & **SYZ**, 1908.00189

$$\Delta^{-s} \frac{d}{dr} \left( \Delta^{s+1} \frac{dR(r)}{dr} \right) + U(\omega, r) R(r) = \mathcal{I}_1(\omega, r)$$



$$\phi = \Delta^{\frac{s}{2}} (r^2 + a^2)^{\frac{1}{2}} R, \quad \frac{d}{dx} = \frac{\Delta}{r^2 + a^2} \frac{d}{dr}$$

$$\frac{d^2 \phi}{dx^2} + (\omega^2 - V(\omega, x)) \phi = \mathcal{I}(\omega, x)$$

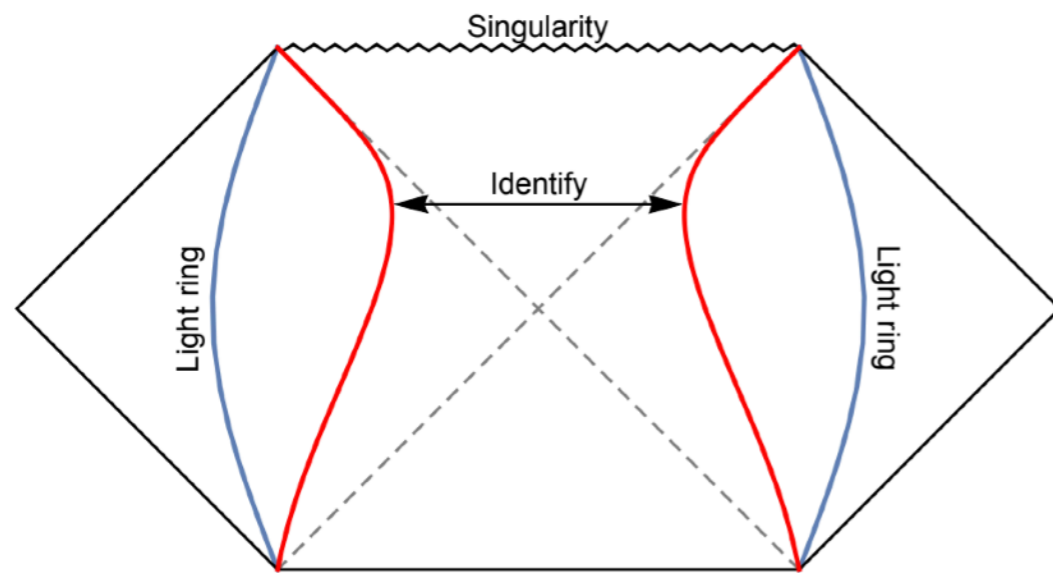
$$V(\omega, x) = \frac{3}{4} \frac{\Delta^2 r}{(r^2 + a^2)^5} - \frac{2\Delta}{(r^2 + a^2)^3} (\Delta + 2r^2) + \frac{a^2 m^2}{(r^2 + a^2)^2} - \frac{2am\lambda}{(r^2 + a^2)} + \frac{(4i\omega r - \lambda)\Delta}{(r^2 + a^2)^2} - \frac{s\Delta + s^2(r - M)^2}{(r^2 + a^2)^2}$$

complex, frequency-dependent effective potential  
but Fredholm formalism still applies

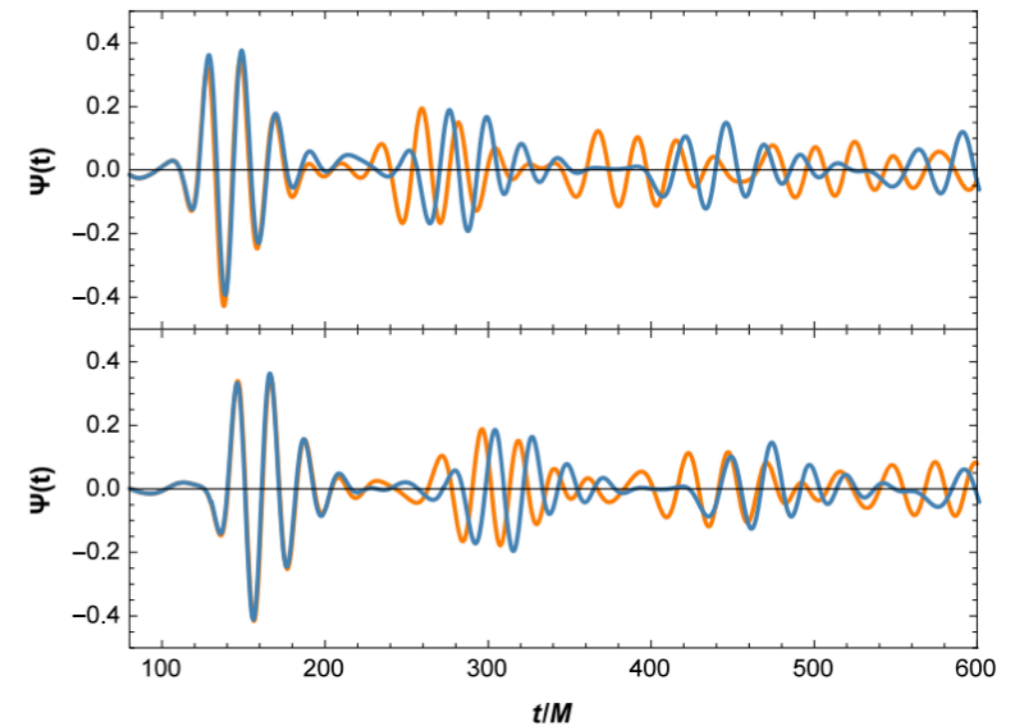


# Unequal interval echoes?

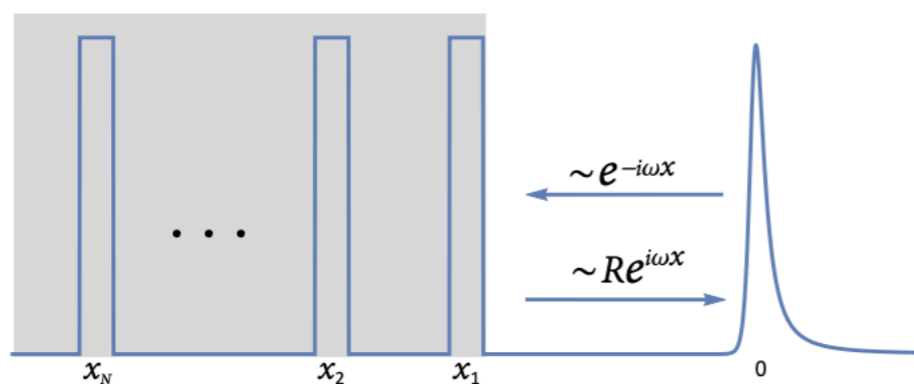
## Slowly pinch-off worm hole



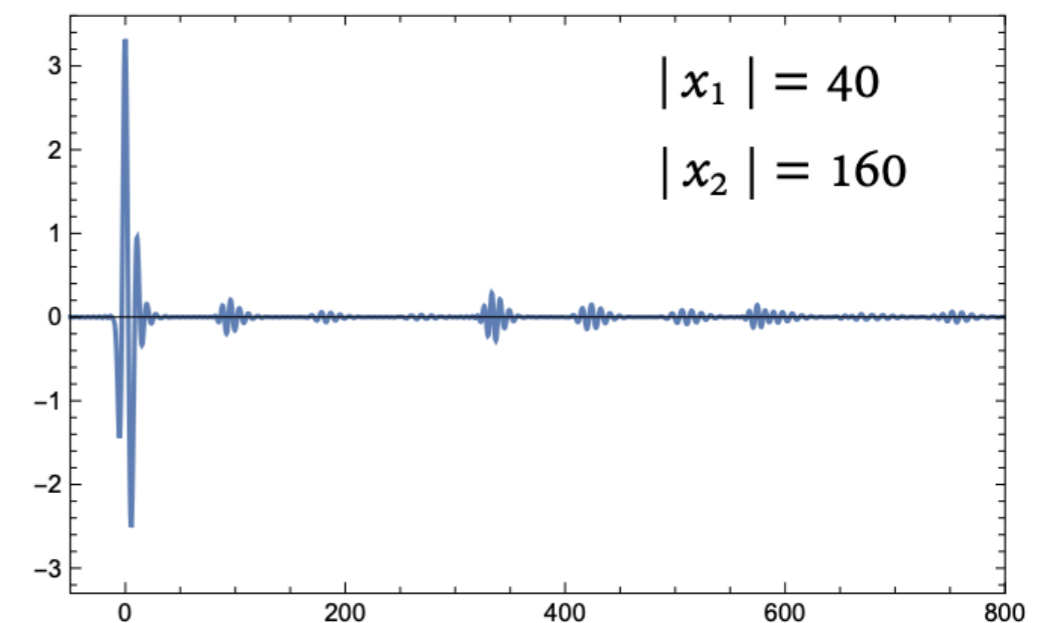
Wang, Zhang, **SYZ** & Piao, 1802.02003



## Multiple extra barriers



Li & Piao, 1904.05652



# Model selection: equal vs unequal

Wang, Zhang, **SYZ** & Piao, 1904.00212

Equal interval echoes

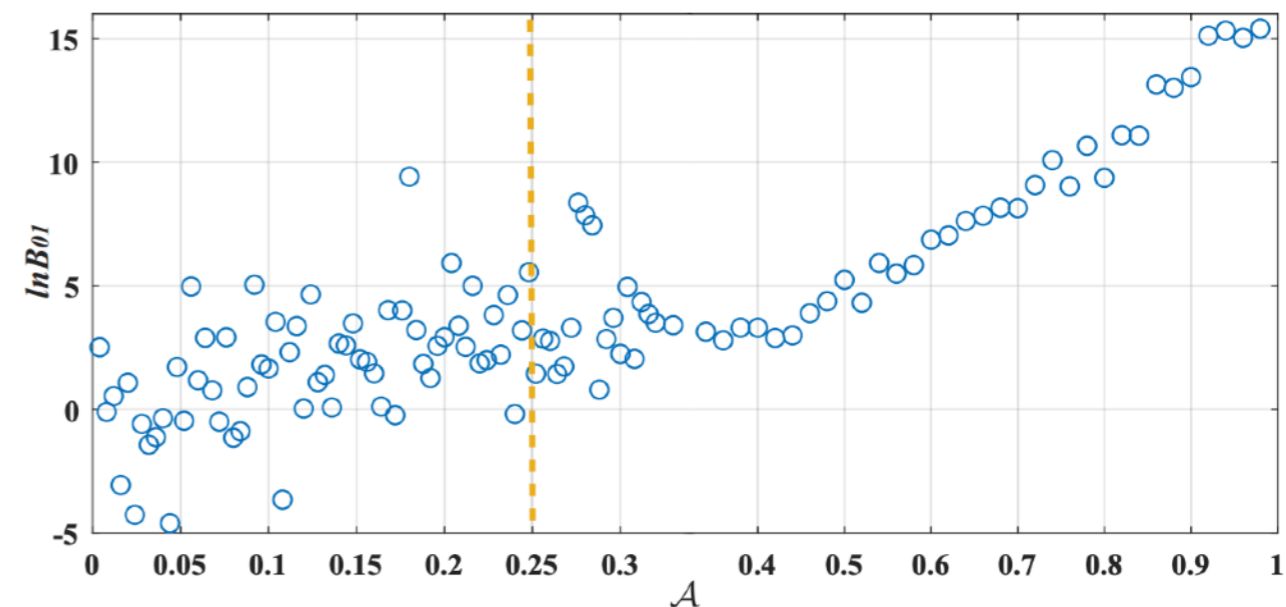
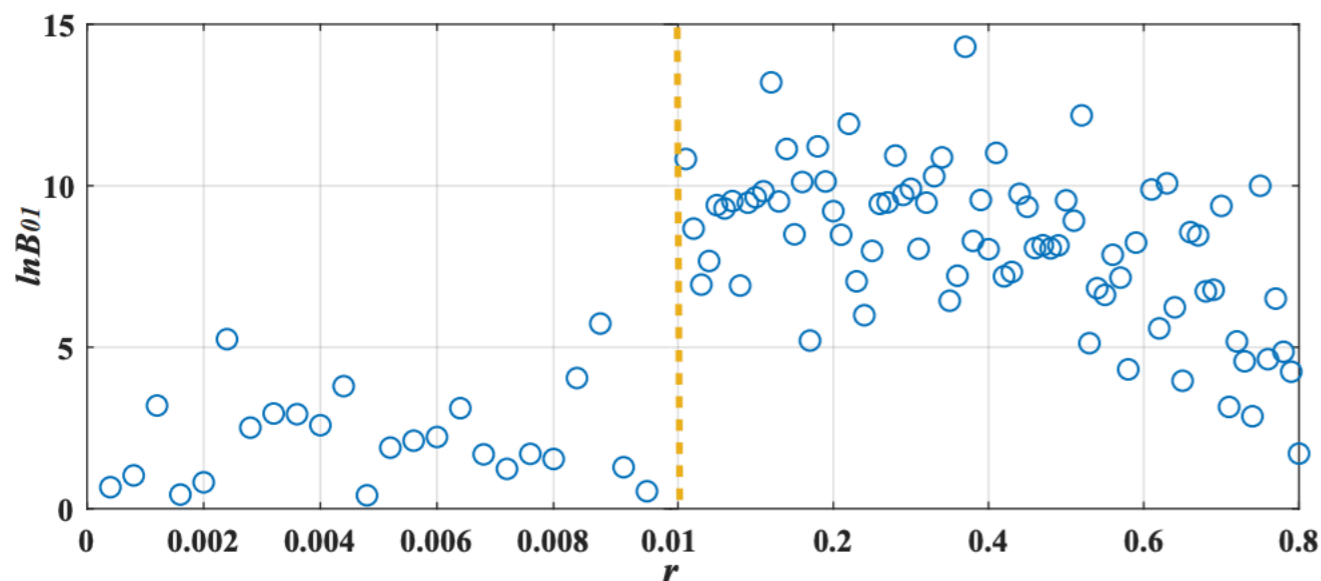
$$x_n = t - t_{echo} - n\Delta t_{echo}$$

Unequal interval echoes

$$x_n = t - t_{echo} - n\Delta t_{echo} - \frac{n(n+1)}{2}r\Delta t_{echo}$$

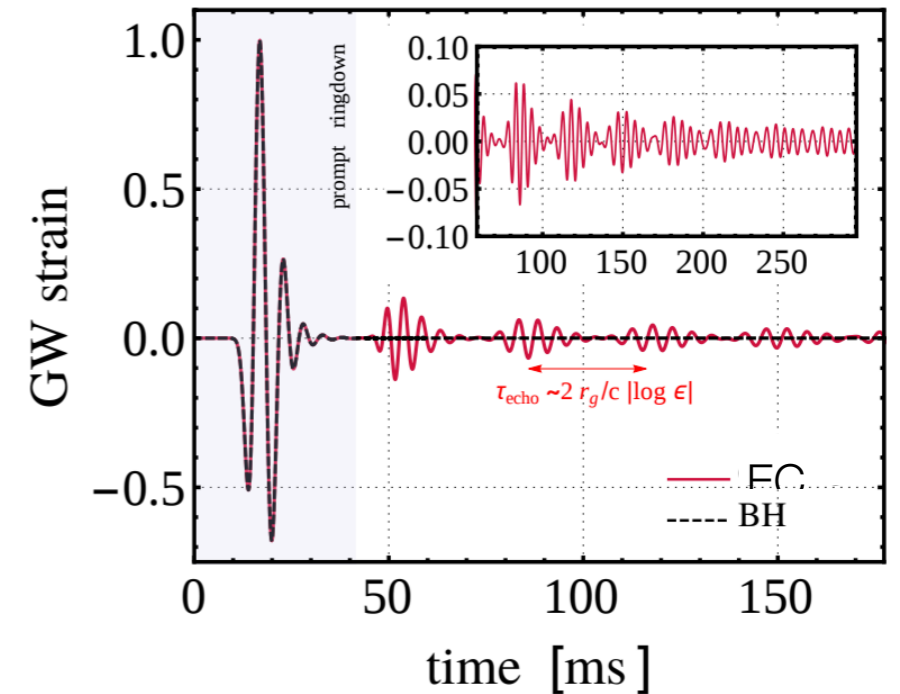
use unequal echo templates if variation of echo intervals is greater than statistical errors of intervals

Bayes factor  $\ln B_{01} = \ln p(d|\mathcal{H}_1, I) - \ln p(d|\mathcal{H}_0, I)$



# Summary

- We have not observed BH horizons.
- GW echoes encode new physics.
- Massive graviton  $\rightarrow$  Echoes
- Possibility of unequal echo intervals
- We use the Fredholm approach to re-process GW echoes.
  - This approach can be presented diagrammatically.
  - Numerically, the Fredholm formalism is very accurate.



Thank you

