

Analytic Two-Loop Higgs Amplitudes and Maximal Transcendentality Principle

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Based on the work with Qingjun Jin (靳庆军)
arXiv:1804.04653 and in progress

中科大交叉中心
October 11, 2018



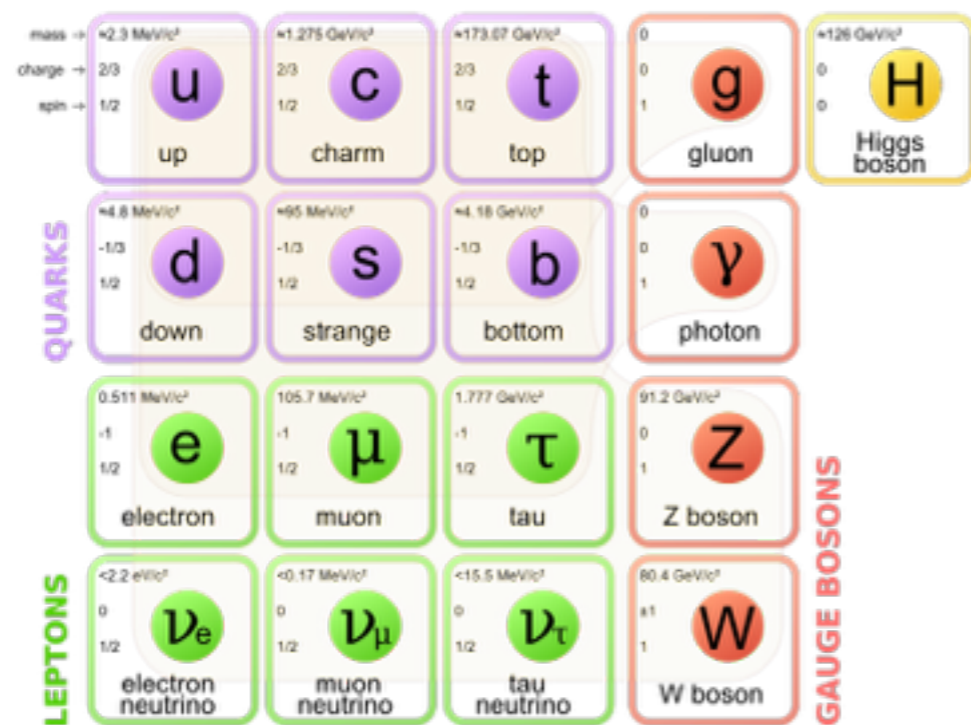
Content

- Motivation
- Computation
- Results
- Summary and outlook



Quantum Field theory (QFT)

QFT is the foundation of modern theoretical physics: particle physics, condensed matter, gravity and cosmology, etc.



Standard Model (SM)
of Particle Physics

Physics 2013



Photo: Pnicolet via
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François Englert



Photo: G-M Greuel
via Wikimedia
Commons
Peter W. Higgs

SM-like Higgs
boson discovered

Challenges

Experimental

Efficient
perturbative
methods

Theoretical

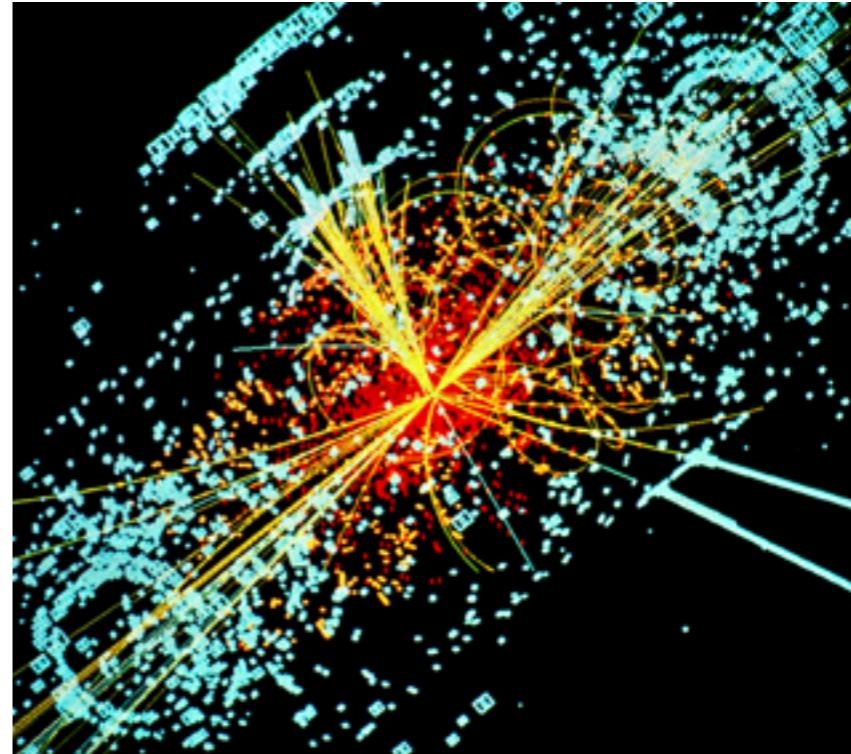
Non-perturbative
method



Large Hadron Collider (LHC)

Precise test of SM

New physics?!



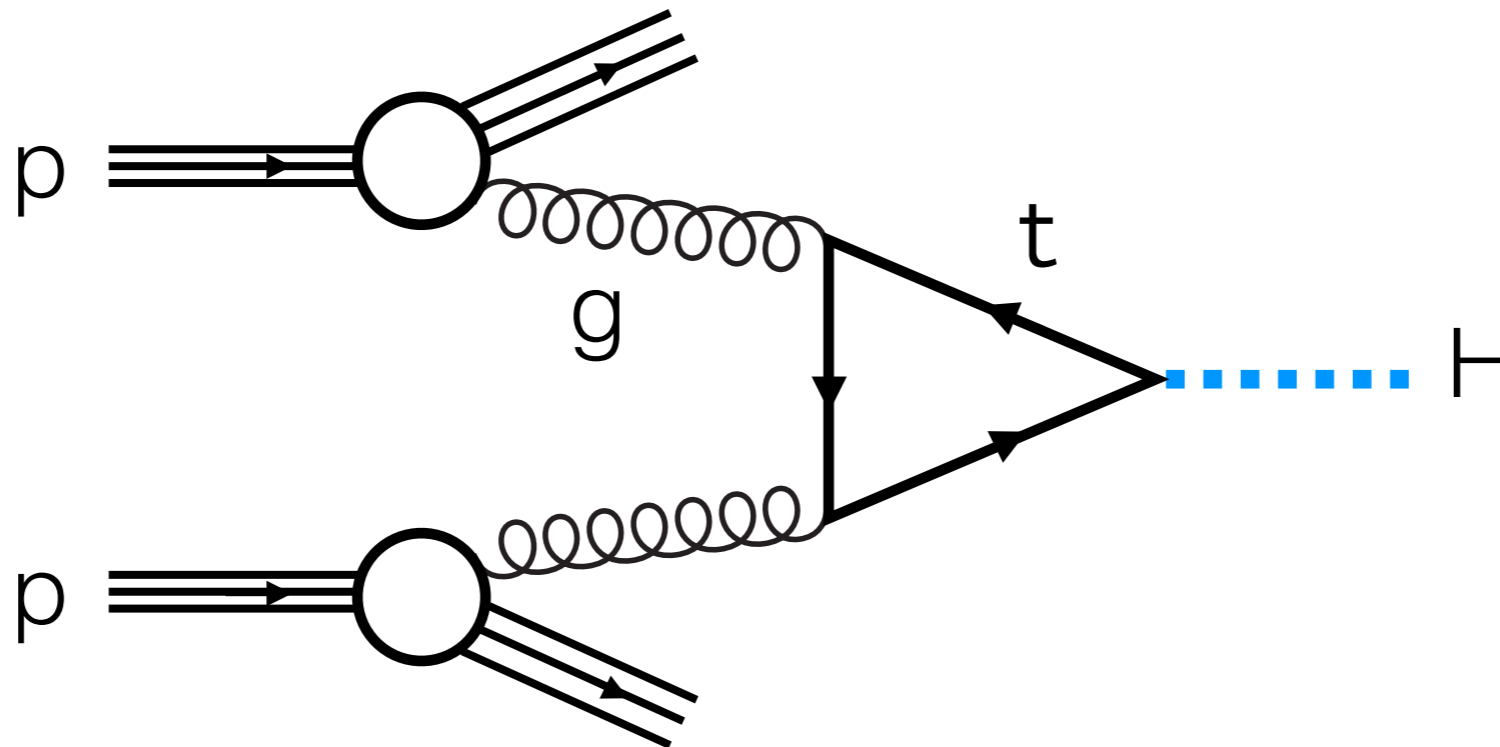
Simulated CMS event

Higher energy and luminosity -> increasing precision

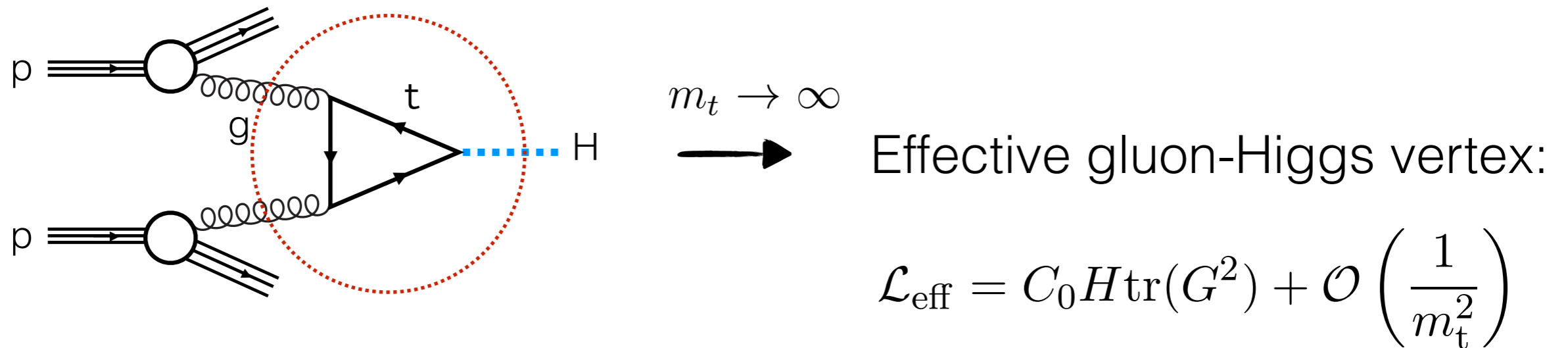
Precise theoretical prediction — at two or higher loops —
of the scattering processes is mandatory.

Higgs boson @ LHC

The dominant production mechanism is the gluon fusion through a top quark loop.



Effective Field Theory (EFT)



There have been computations for inclusive Higgs production to N³LO orders in the heavy quark limit.

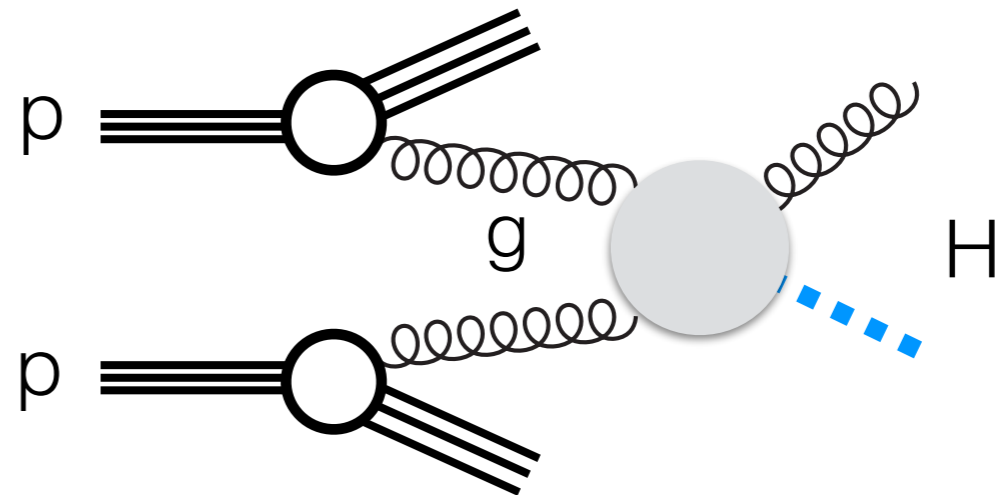
[Anastasiou, Duhr, Dulat, Furlan, Gehrmann, Herzog, Lazopoulos, Mistlberger 2016]



Effective Field Theory (EFT)

Higgs plus jet production is sensitive to new physics.

EFT description is not good when $p_T \sim 2m_t$



High dimension operators contribution are important.

$$\mathcal{L}_{\text{eff}} = C_0 O_0 + \frac{1}{m_t^2} \sum_{i=1}^4 C_i O_i + \mathcal{O}\left(\frac{1}{m_t^4}\right)$$

Dimension-7 operators

$$O_1 = H \text{Tr}(G_\mu^\nu G_\nu^\rho G_\rho^\mu),$$

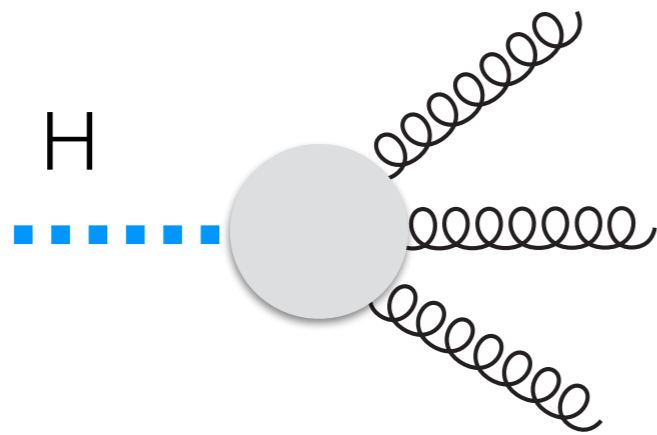
$$O_2 = H \text{Tr}(D_\rho G_{\mu\nu} D^\rho G^{\mu\nu}),$$

$$O_3 = H \text{Tr}(D^\rho G_{\rho\mu} D_\sigma G^{\sigma\mu}),$$

$$O_4 = H \text{Tr}(G_{\mu\rho} D^\rho D_\sigma G^{\sigma\mu}).$$

Goal

Compute two-loop Higgs amplitudes with dim-7 operators



$$\mathcal{L}_{\text{eff}} = C_0 O_0 + \frac{1}{m_t^2} \sum_{i=1}^4 C_i O_i + \mathcal{O}\left(\frac{1}{m_t^4}\right)$$

$$O_0 = H \text{tr}(G^2)$$

$$O_1 = H \text{Tr}(G_\mu^\nu G_\nu^\rho G_\rho^\mu),$$

$$O_2 = H \text{Tr}(D_\rho G_{\mu\nu} D^\rho G^{\mu\nu}),$$

$$O_3 = H \text{Tr}(D^\rho G_{\rho\mu} D_\sigma G^{\sigma\mu}),$$

$$O_4 = H \text{Tr}(G_{\mu\rho} D^\rho D_\sigma G^{\sigma\mu}).$$

for pure YM

This provides the two-loop virtual amplitudes for the top mass correction in EFT.

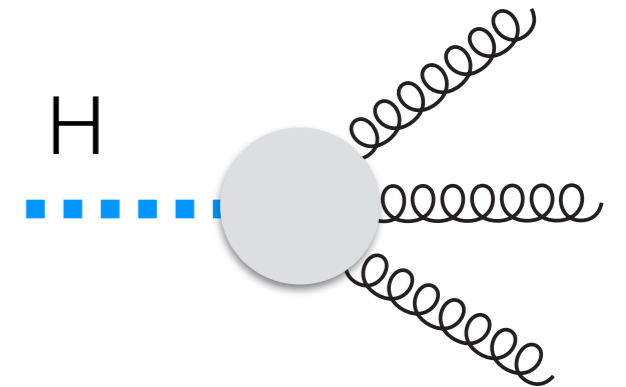
Form factors

Higgs amplitudes are equivalent to **form factors**:

$$F_{\mathcal{O}_i,n} = \int d^4x e^{-iq \cdot x} \langle p_1, \dots, p_n | \mathcal{O}_i(x) | 0 \rangle$$

Linear relation:

$$\mathcal{O}_2 = \frac{1}{2} \partial^2 \mathcal{O}_0 - 4 g_{\text{YM}} \mathcal{O}_1 + 2 \mathcal{O}_4 \quad \longrightarrow \quad F_{\mathcal{O}_2} = \frac{1}{2} q^2 F_{\mathcal{O}_0} - 4 g_{\text{YM}} F_{\mathcal{O}_1}$$



$$\mathcal{O}_0 = \text{tr}(G_{\mu\nu} G^{\mu\nu}).$$

$$\mathcal{O}_1 = \text{tr}(G_{\mu}^{\nu} G_{\nu}^{\rho} G_{\rho}^{\mu}),$$

$$\mathcal{O}_2 = \text{tr}(D_{\rho} G_{\mu\nu} D^{\rho} G^{\mu\nu}).$$



Theoretical motivations

Feynman diagram?



Feynman diagram method works in principle, but the complexity grows extremely fast with increasing number of external legs / loops.

n-gluon tree amplitudes:

n	2	3	4	5	6	7	8
# of diagrams	4	25	220	2485	34300	559405	10525900

Surprising simplicity

MHV (Maximally-helicity-violating) amplitudes:

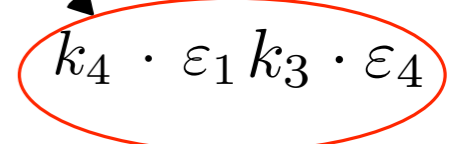
[Parke, Taylor '86]

$$A_n^{\text{tree}}(1^+, \dots, i^-, \dots, j^-, \dots, n^+) = \frac{\langle ij \rangle^4}{\langle 12 \rangle \dots \langle n1 \rangle}$$

Comparing with result of Feynman diagrams:

The image shows a grid of small Feynman diagrams and mathematical expressions, likely representing the expansion of the MHV amplitude into Feynman diagrams. The diagrams are arranged in a grid, with some containing mathematical terms like $k_4 \cdot \epsilon_1 k_3 \cdot \epsilon_4$. A red circle highlights the expression $k_4 \cdot \epsilon_1 k_3 \cdot \epsilon_4$ in the bottom right corner.

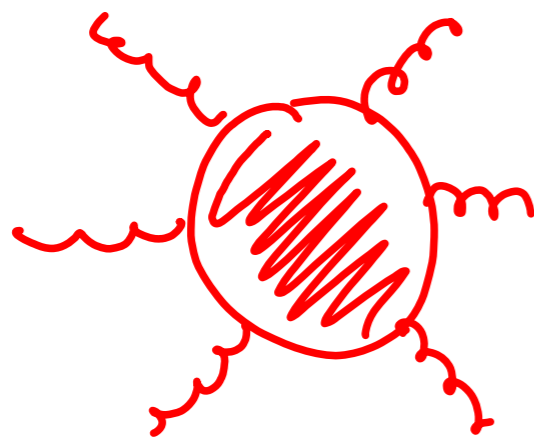
[Bern '93]



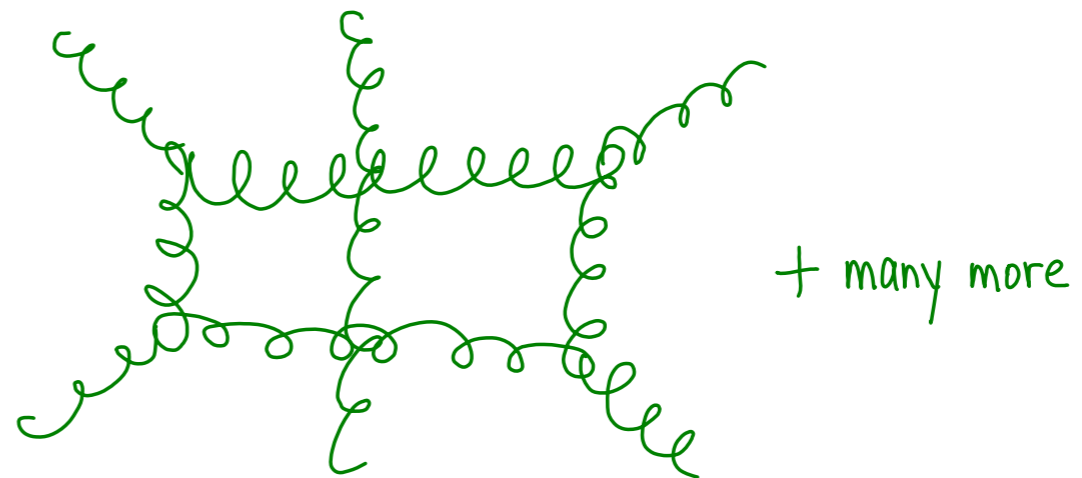
Surprising simplicity

A more non-trivial example of two-loop amplitudes:

Six-gluon MHV amplitudes in N=4 SYM



=



[Del Duca, Duhr, Smirnov 2010]

a heroic analytical computation

$$\begin{aligned}
& \frac{3}{4}\mathcal{G}\left(v_{231}, 1, \frac{1}{1-u_2}; 1\right) H(0; u_3) + \frac{3}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-u_2}, 1; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{312}, 1, \frac{1}{1-u_3}; 1\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{312}, \frac{1}{1-u_3}, 1; 1\right) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{321}, 1, \frac{1}{1-u_3}; 1\right) H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{321}, \frac{1}{1-u_3}, 1; 1\right) H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(0; u_1) H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right) H(0; u_1) H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(0; u_1) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, u_{231}; 1\right) H(0; u_1) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right) H(0; u_1) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0; u_1) H(0; u_3) + \\
& \frac{5}{24}\pi^2 H(0; u_1) H(0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right) H(0; u_2) H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right) H(0; u_2) H(0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, u_{123}; 1\right) H(0; u_2) H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0; u_2) H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0; u_2) H(0; u_3) + \frac{5}{24}\pi^2 H(0; u_2) H(0; u_3) + \\
& 3H(0; u_2) H(0, 0; u_1) H(0; u_3) + 3H(0; u_1) H(0, 0; u_2) H(0; u_3) + \\
& \frac{1}{4}H(0; u_2) H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) H(0; u_3) + \frac{1}{2}H(0; u_1) H(0, 1; (u_1+u_3)) H(0; u_3) + \\
& \frac{1}{4}H(0; u_1) H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) H(0; u_3) + \frac{1}{2}H(0; u_2) H(0, 1; (u_2+u_3)) H(0; u_3) + \\
& \frac{3}{4}H(0; u_2) H(1, 0; u_1) H(0; u_3) + \frac{3}{4}H(0; u_1) H(1, 0; u_2) H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right) H(0, 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0, 0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(0, 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right) H(0, 0; u_1) - \frac{23}{24}\pi^2 H(0, 0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0, 0; u_2) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{312}; 1\right) H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right) H(0, 0; u_2) - \\
& \frac{25}{4}H(0, 0; u_1) H(0, 0; u_2) - \frac{23}{24}\pi^2 H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right) H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right) H(0, 0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right) H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right) H(0, 0; u_3) + 3H(0; u_1) H(0; u_2) H(0, 0; u_3) - \\
& \frac{25}{4}H(0, 0; u_1) H(0, 0; u_3) - \frac{25}{4}H(0, 0; u_2) H(0, 0; u_3) - \frac{23}{24}\pi^2 H(0, 0; u_3) + \frac{1}{12}\pi^2 H(0, 1; u_1) + \\
& \frac{1}{12}\pi^2 H(0, 1; u_2) - \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) + \frac{1}{2}H(0; u_1) H(0; u_2) H(0, 1; (u_1+u_2)) +
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{12}\pi^2 H(0, 1; (u_1+u_2)) + \frac{1}{12}\pi^2 H(0, 1; u_3) + \frac{1}{4}H(0; u_1) H(0; u_2) H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) - \\
& \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) + \frac{1}{12}\pi^2 H(0, 1; (u_1+u_3)) - \frac{1}{24}\pi^2 H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) + \\
& \frac{1}{12}\pi^2 H(0, 1; (u_2+u_3)) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, \frac{u_1-1}{u_1+u_3-1}; 1\right) H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_1) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, u_{312}; 1\right) H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_2) H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_3) H(1, 0; u_1) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) H(1, 0; u_1) - \frac{1}{3}\pi^2 H(1, 0; u_1) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right) H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1+u_2}; 1\right) H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_2) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, u_{123}; 1\right) H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_1) H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_3) H(1, 0; u_2) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) H(1, 0; u_2) - \frac{1}{4}H(1, 0; u_1) H(1, 0; u_2) - \frac{1}{3}\pi^2 H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{3}\pi^2 H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right) H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right) H(1, 0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, u_{231}; 1\right) H(1, 0; u_3) + \\
& \frac{3}{4}H(0; u_1) H(0; u_2) H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_1) H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_2) H(1, 0; u_3) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_1) H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_2) H(1, 0; u_3) + \\
& \frac{1}{24}\pi^2 H(1, 1; u_1) + \frac{1}{24}\pi^2 H(1, 1; u_2) + \frac{1}{24}\pi^2 H(1, 1; u_3) + \frac{1}{2}H(0; u_2) H(0, 0; u_1) + \\
& \frac{1}{2}H(0; u_3) H(0, 0; u_2) + \frac{1}{2}H(0; u_1) H(0, 0; u_3) - \frac{1}{2}H(0; u_2) H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - \\
& \frac{1}{2}H(0; u_3) H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - H(0; u_1) H(0, 0, 1; (u_1+u_2)) - \\
& H(0; u_2) H(0, 0, 1; (u_1+u_2)) - \frac{1}{2}H(0; u_1) H\left(0, 0, 1; \frac{u_1+u_3-1}{u_1-1}\right) -
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{u_{312}}\right)+\frac{1}{4}H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{123}}\right)-\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{123}}\right)- \\
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{132}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{132}}\right)+\frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{213}}\right)- \\
& \frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{213}}\right)-\frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{231}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(0,1,1;\frac{1}{v_{231}}\right)+ \\
& \frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{312}}\right)-\frac{1}{4}H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{312}}\right)-\frac{1}{4}H(0;u_1)\mathcal{H}\left(0,1,1;\frac{1}{v_{321}}\right)+ \\
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(0,1,1;\frac{1}{v_{321}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(1,0,1;\frac{1}{u_{123}}\right)+\frac{1}{4}H(0;u_1)\mathcal{H}\left(1,0,1;\frac{1}{u_{231}}\right)+ \\
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(1,0,1;\frac{1}{u_{312}}\right)+\frac{1}{4}H(0;u_2)\mathcal{H}\left(1,0,1;\frac{1}{v_{123}}\right)-\frac{1}{4}H(0;u_3)\mathcal{H}\left(1,0,1;\frac{1}{v_{123}}\right)- \\
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(1,0,1;\frac{1}{v_{132}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(1,0,1;\frac{1}{v_{132}}\right)+\frac{1}{4}H(0;u_1)\mathcal{H}\left(1,0,1;\frac{1}{v_{213}}\right)- \\
& \frac{1}{4}H(0;u_3)\mathcal{H}\left(1,0,1;\frac{1}{v_{213}}\right)-\frac{1}{4}H(0;u_1)\mathcal{H}\left(1,0,1;\frac{1}{v_{231}}\right)+\frac{1}{4}H(0;u_3)\mathcal{H}\left(1,0,1;\frac{1}{v_{231}}\right)+ \\
& \frac{1}{4}H(0;u_1)\mathcal{H}\left(1,0,1;\frac{1}{v_{312}}\right)-\frac{1}{4}H(0;u_2)\mathcal{H}\left(1,0,1;\frac{1}{v_{312}}\right)-\frac{1}{4}H(0;u_1)\mathcal{H}\left(1,0,1;\frac{1}{v_{321}}\right)+ \\
& \frac{1}{4}H(0;u_2)\mathcal{H}\left(1,0,1;\frac{1}{v_{321}}\right)+H(0;u_2)\mathcal{H}\left(1,1,1;\frac{1}{v_{123}}\right)-H(0;u_3)\mathcal{H}\left(1,1,1;\frac{1}{v_{123}}\right)- \\
& H(0;u_1)\mathcal{H}\left(1,1,1;\frac{1}{v_{231}}\right)+H(0;u_3)\mathcal{H}\left(1,1,1;\frac{1}{v_{231}}\right)+H(0;u_1)\mathcal{H}\left(1,1,1;\frac{1}{v_{312}}\right)- \\
& H(0;u_2)\mathcal{H}\left(1,1,1;\frac{1}{v_{312}}\right)-\frac{3}{2}\mathcal{H}\left(0,0,0,1;\frac{1}{u_{123}}\right)-\frac{3}{2}\mathcal{H}\left(0,0,0,1;\frac{1}{u_{231}}\right)- \\
& \frac{3}{2}\mathcal{H}\left(0,0,0,1;\frac{1}{u_{312}}\right)-3\mathcal{H}\left(0,0,0,1;\frac{1}{v_{132}}\right)-3\mathcal{H}\left(0,0,0,1;\frac{1}{v_{213}}\right)-3\mathcal{H}\left(0,0,0,1;\frac{1}{v_{321}}\right)- \\
& \frac{1}{2}\mathcal{H}\left(0,0,1,1;\frac{1}{u_{123}}\right)-\frac{1}{2}\mathcal{H}\left(0,0,1,1;\frac{1}{u_{231}}\right)-\frac{1}{2}\mathcal{H}\left(0,0,1,1;\frac{1}{u_{312}}\right)- \\
& \frac{1}{2}\mathcal{H}\left(0,1,0,1;\frac{1}{u_{123}}\right)-\frac{1}{2}\mathcal{H}\left(0,1,0,1;\frac{1}{u_{231}}\right)-\frac{1}{2}\mathcal{H}\left(0,1,0,1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{4}\mathcal{H}\left(0,1,1,1;\frac{1}{v_{123}}\right)+\frac{1}{4}\mathcal{H}\left(0,1,1,1;\frac{1}{v_{132}}\right)+\zeta_3H(0;u_1)+\zeta_3H(0;u_2)+\zeta_3H(0;u_3)+ \\
& \frac{5}{2}\zeta_3H(1;u_1)+\frac{5}{2}\zeta_3H(1;u_2)+\frac{5}{2}\zeta_3H(1;u_3)+\frac{1}{2}\zeta_3\mathcal{H}\left(1;\frac{1}{u_{123}}\right)+\frac{1}{2}\zeta_3\mathcal{H}\left(1;\frac{1}{u_{231}}\right)+ \\
& \frac{1}{2}\zeta_3\mathcal{H}\left(1;\frac{1}{u_{312}}\right)-\frac{1}{2}\mathcal{H}\left(1,0,0,1;\frac{1}{u_{123}}\right)-\frac{1}{2}\mathcal{H}\left(1,0,0,1;\frac{1}{u_{231}}\right)-\frac{1}{2}\mathcal{H}\left(1,0,0,1;\frac{1}{u_{312}}\right)+ \\
& \frac{1}{4}\zeta_3\mathcal{H}\left(1;\frac{1}{v_{123}}\right)+\frac{1}{4}\zeta_3\mathcal{H}\left(1;\frac{1}{v_{132}}\right)+\frac{1}{4}\zeta_3\mathcal{H}\left(1;\frac{1}{v_{213}}\right)+\frac{1}{4}\zeta_3\mathcal{H}\left(1;\frac{1}{v_{231}}\right)+\frac{1}{4}\zeta_3\mathcal{H}\left(1;\frac{1}{v_{312}}\right)+ \\
& \frac{1}{4}\zeta_3\mathcal{H}\left(1;\frac{1}{v_{321}}\right)+\frac{1}{4}\mathcal{H}\left(0,1,1,1;\frac{1}{v_{213}}\right)+\frac{1}{4}\mathcal{H}\left(0,1,1,1;\frac{1}{v_{231}}\right)+\frac{1}{4}\mathcal{H}\left(0,1,1,1;\frac{1}{v_{312}}\right)+ \\
& \frac{1}{4}\mathcal{H}\left(0,1,1,1;\frac{1}{v_{321}}\right)+\frac{1}{4}\mathcal{H}\left(1,0,1,1;\frac{1}{v_{123}}\right)+\frac{1}{4}\mathcal{H}\left(1,0,1,1;\frac{1}{v_{132}}\right)+\frac{1}{4}\mathcal{H}\left(1,0,1,1;\frac{1}{v_{213}}\right)+ \\
& \frac{1}{4}\mathcal{H}\left(1,0,1,1;\frac{1}{v_{231}}\right)+\frac{1}{4}\mathcal{H}\left(1,0,1,1;\frac{1}{v_{312}}\right)+\frac{1}{4}\mathcal{H}\left(1,0,1,1;\frac{1}{v_{321}}\right)+\frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{123}}\right)+ \\
& \frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{132}}\right)+\frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{213}}\right)+\frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{231}}\right)+\frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{312}}\right)+ \\
& \frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{321}}\right)
\end{aligned}$$

$$\frac{1}{4}\mathcal{H}\left(1,1,0,1;\frac{1}{v_{321}}\right)+\frac{3}{2}\mathcal{H}\left(1,1,1,1;\frac{1}{v_{123}}\right)+\frac{3}{2}\mathcal{H}\left(1,1,1,1;\frac{1}{v_{231}}\right)+\frac{3}{2}\mathcal{H}\left(1,1,1,1;\frac{1}{v_{312}}\right)$$

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Result can be remarkably simple

17 pages =

[Goncharov, Spradlin, Vergu, Volovich 2010]

$$\sum_{i=1}^3 \left(L_4(x_i^+, x_i^-) - \frac{1}{2} \text{Li}_4(1 - 1/u_i) \right) - \frac{1}{8} \left(\sum_{i=1}^3 \text{Li}_2(1 - 1/u_i) \right)^2 + \frac{1}{24} J^4 + \frac{\pi^2}{12} J^2 + \frac{\pi^4}{72}$$

which is a single line of only classical polylogarithms!

→ require advanced mathematical tools: "Symbol"

Why so simple?

Many examples show that the final result can be put in a form which is far simpler than the intermediate steps !

Surprising simplicity



Hidden structure

“Theoretical experiment”:

looking into the theoretical data, and try to find hidden structures

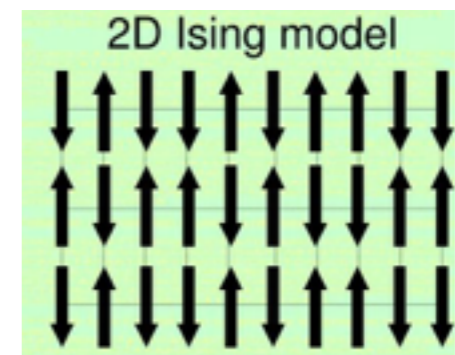
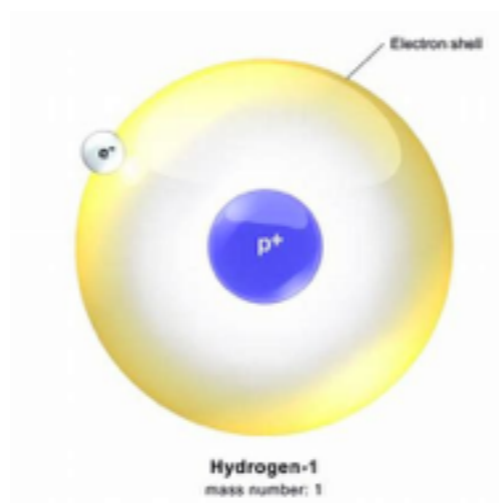
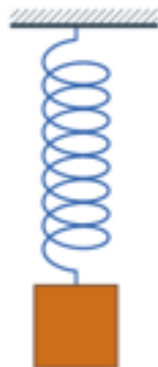
A toy model

N=4 SYM theory : -> QCD's maximally supersymmetric cousin

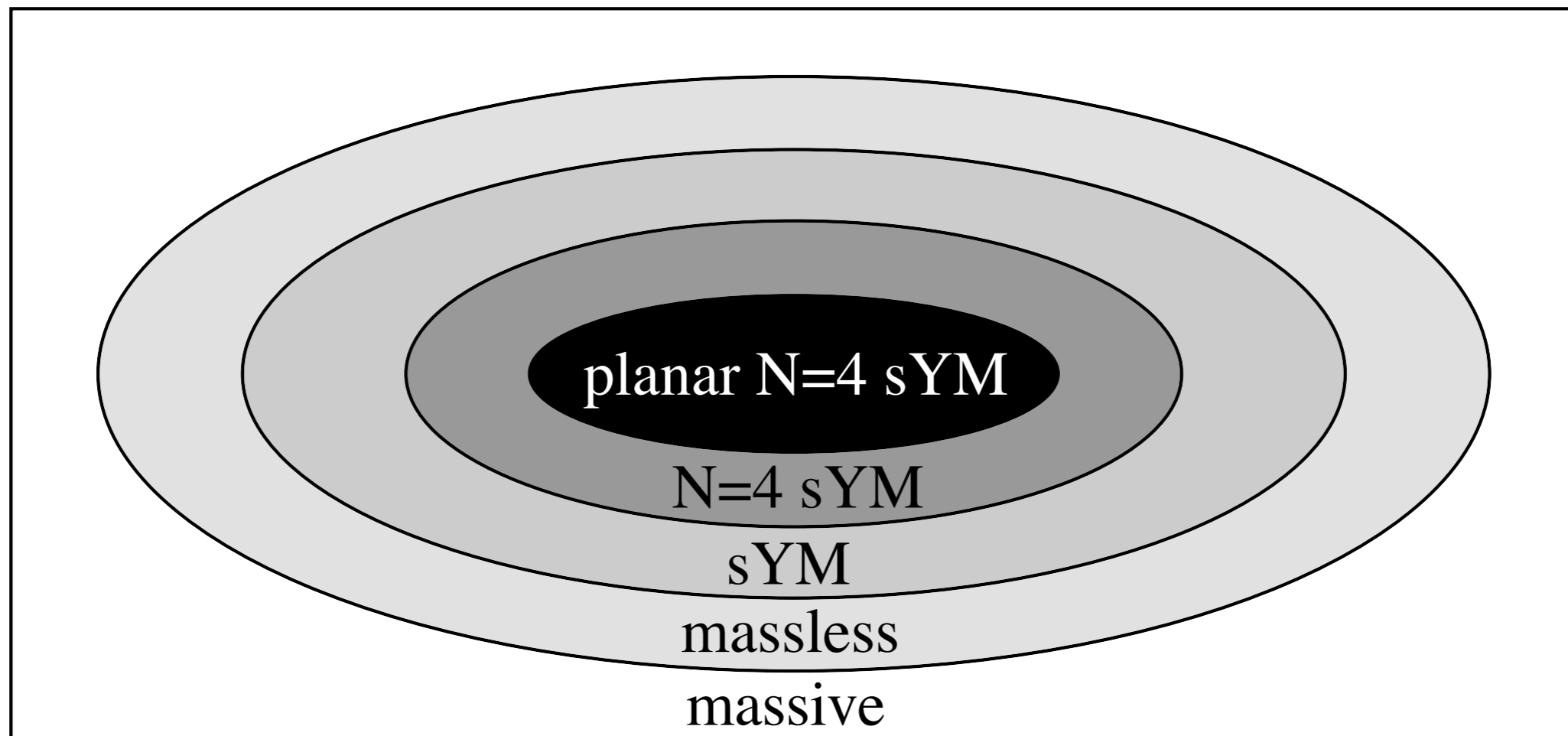
$$\mathcal{L} = -\frac{1}{g_{\text{YM}}^2} \text{Tr}(F_{\mu\nu} F^{\mu\nu}) + \text{fermions} + \text{scalars}$$

where all fields are the in the adjoint representation of the gauge group $SU(N_c)$.

Exactly solvable in planar limit!



Hierarchy of simplicity



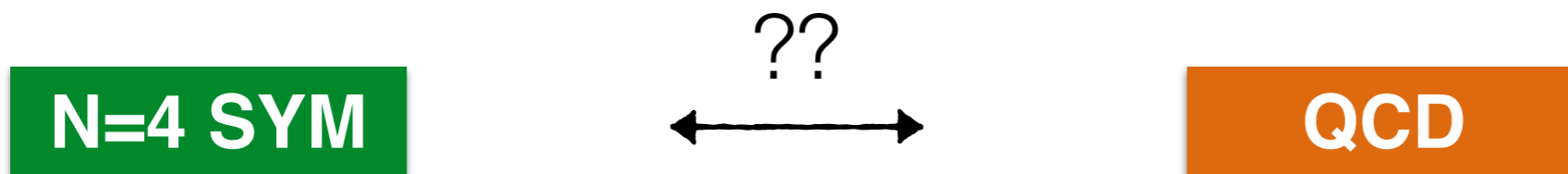
Lance Dixon 1105.0771



N=4 SYM

Techniques first developed by studying this toy model are used in general theories such as QCD, e.g.:
BCFW recursion relations, unitarity on-shell method.

Are there direct connections between the two theories?



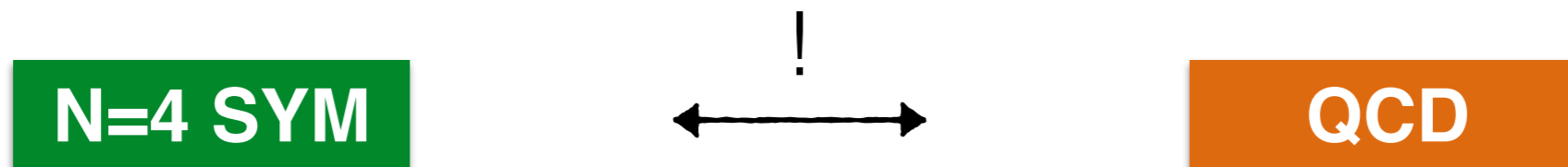
Maximal Transcendentality Principle



Maximally transcendental parts are equal between two theories!?

Number	Function	Transcendentality degree
$2/3, \sqrt{2}$	rational function	0
π	Log(x)	1
Riemann zeta value $\zeta(n)$	Polylog function $\text{Li}_n(x)$	n

Known examples



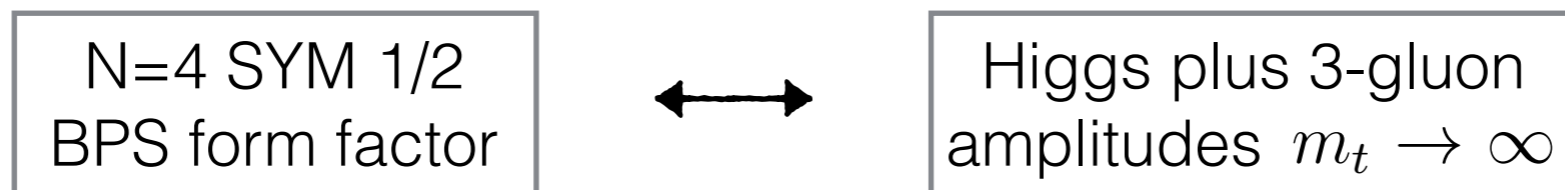
- Anomalous dimension of twist-2 operators

$$\gamma^{\mathcal{N}=4}(j) = \gamma^{\text{QCD}}(j)|_{\text{max. trans}}$$

[Kotikov, Lipatov, Onishchenko, Velizhanin 2004]

- Two-loop remainder function (kinematic dependent functions!):

[Brandhuber, Travaglini, GY 2012]



- Also for certain Wilson lines [Li, Manteuffel, Schabinger, Zhu 2014]

Two-loop Higgs to 3-gluon

$$\begin{aligned}
 & -2G(0,0,1,0,u) + G(0,0,1-v,1-v,u) + 2G(0,0,-v,1-v,u) - G(0,1,0,1-v,u) + 4G(0,1,1,0,u) - G(0,1,1-v,0,u) + G(0,1-v,0,1-v,u) \\
 & + G(0,1-v,1-v,0,u) - G(0,1-v,-v,1-v,u) + 2G(0,-v,0,1-v,u) + 2G(0,-v,1-v,0,u) - 2G(0,-v,1-v,1-v,u) - 2G(1,0,0,1-v,u) \\
 & - 2G(1,0,1-v,0,u) + 4G(1,1,0,0,u) - 4G(1,1,1,0,u) - 2G(1,1-v,0,0,u) + G(1-v,0,0,1-v,u) - G(1-v,0,1,0,u) - 2G(-v,1-v,1-v,u)H(0,v) \\
 & - 2G(1-v,1,0,0,u) + 2G(1-v,1,0,1-v,u) + 2G(1-v,1,1-v,0,u) + G(1-v,1-v,0,0,u) + 2G(1-v,1-v,1,0,u) - 2G(1-v,1-v,-v,1-v,u) \\
 & - G(1-v,-v,1-v,0,u) + 4G(1-v,-v,-v,1-v,u) - 2G(-v,0,1-v,1-v,u) - 2G(-v,1-v,0,1-v,u) - 2G(-v,1-v,1-v,0,u) + 4G(1,0,1,0,u) \\
 & + 4G(-v,-v,1-v,1-v,u) - 4G(-v,-v,-v,1-v,u) - G(0,0,1-v,u)H(0,v) - G(0,1,0,u)H(0,v) - G(0,1-v,0,u)H(0,v) + G(0,1-v,1-v,u)H(0,v) \\
 & - G(0,-v,1-v,u)H(0,v) - 2G(1,0,0,u)H(0,v) + G(1,0,1-v,u)H(0,v) + G(1,1-v,0,u)H(0,v) + G(1-v,0,0,u)H(0,v) - G(1-v,0,1-v,u)H(0,v) \\
 & - G(1-v,1,0,u)H(0,v) - G(1-v,1-v,0,u)H(0,v) - G(1-v,-v,1-v,u)H(0,v) + G(-v,0,1-v,u)H(0,v) + G(-v,1-v,0,u)H(0,v) + H(1,0,0,1,v) \\
 & - G(0,0,1-v,u)H(1,v) - G(0,0,-v,u)H(1,v) + G(0,1,0,u)H(1,v) - G(0,1-v,0,u)H(1,v) + G(0,1-v,-v,u)H(1,v) - 2G(0,-v,0,u)H(1,v) \\
 & + 2G(0,-v,1-v,u)H(1,v) + 2G(1,0,0,u)H(1,v) - G(1-v,0,0,u)H(1,v) + G(1-v,0,-v,u)H(1,v) - 2G(1-v,1,0,u)H(1,v) - G(1-v,0,-v,1-v,u) \\
 & + G(1-v,-v,0,u)H(1,v) - 4G(1-v,-v,-v,u)H(1,v) + 2G(-v,0,1-v,u)H(1,v) + 2G(-v,1-v,0,u)H(1,v) - 4G(-v,1-v,-v,u)H(1,v) \\
 & - 4G(-v,-v,1-v,u)H(1,v) + 4G(-v,-v,-v,u)H(1,v) + G(0,0,u)H(0,0,v) + G(0,1-v,u)H(0,0,v) + G(1-v,0,u)H(0,0,v) + H(1,0,1,0,v) \\
 & - G(0,0,u)H(0,1,v) + G(0,-v,u)H(0,1,v) - G(1-v,0,u)H(0,1,v) + 2G(1-v,1-v,u)H(0,1,v) - 3G(1-v,-v,u)H(0,1,v) \\
 & - G(-v,0,u)H(0,1,v) - 2G(-v,1-v,u)H(0,1,v) + G(1,1,v) - G(0,0,u)H(1,0,v) + G(0,-v,u)H(1,0,v) - G(1,0,u)H(1,0,v) \\
 & + 2G(1-v,0,u)H(1,0,v) - 2G(1-v,1-v,u)H(1,0,v) + G(1,0,u)H(1,0,v) - G(-v,0,u)H(1,0,v) + 2G(-v,1-v,u)H(1,0,v) + G(0,0,u)H(1,1,v) \\
 & - 2G(0,-v,u)H(1,1,v) - 2G(-v,0,u)H(1,1,v) + 4G(-v,-v,u)H(1,1,v) + G(0,u)H(0,0,1,v) - 3G(1-v,u)H(0,0,1,v) + 4G(-v,u)H(0,0,1,v) \\
 & + G(0,u)H(0,1,0,v) + G(1-v,u)H(0,1,0,v) - G(0,u)H(0,1,1,v) + 2G(-v,u)H(0,1,1,v) + G(0,u)H(1,0,0,v) + G(1-v,u)H(1,0,0,v) + H(1,1,0,0,v) \\
 & - G(0,u)H(1,0,1,v) + 2G(-v,u)H(1,0,1,v) - G(0,u)H(1,1,0,v) + 4G(1-v,u)H(1,1,0,v) - 2G(-v,u)H(1,1,0,v) + H(0,0,1,1,v) + H(0,1,0,1,v) \\
 & + G(1-v,1-v,u)H(0,0,v) + 2G(1-v,1-v,-v,u)H(1,v) - G(1-v,-v,0,1-v,u) + H(0,1,1,0,v) + G(1-v,0,1-v,0,u) - G(0,1-v,1,0,u) \\
 & + 4G(-v,1-v,-v,1-v,u)
 \end{aligned}$$

[Gehrmann, Jaquier, Glover, Koukoutsakis 2011]

$$G(1-v, 1-v, v, 1-v; u)$$

QCD

Multiple polyLogarithm



$$\begin{aligned}
 & -2 \left[J_4 \left(-\frac{uv}{w} \right) + J_4 \left(-\frac{vw}{u} \right) + J_4 \left(-\frac{wu}{v} \right) \right] - 8 \sum_{i=1}^3 \left[\text{Li}_4 \left(1 - u_i^{-1} \right) + \frac{\log^4 u_i}{4!} \right] \\
 & - 2 \left[\sum_{i=1}^3 \text{Li}_2(1 - u_i) + \frac{\log^2 u_i}{2!} \right]^2 + \frac{1}{2} \left[\sum_{i=1}^3 \log^2 u_i \right]^2 - \frac{\log^4(uvw)}{4!} - \frac{23}{2} \zeta_4
 \end{aligned}$$

[Brandhuber, Travaglini, GY 2012]

N=4

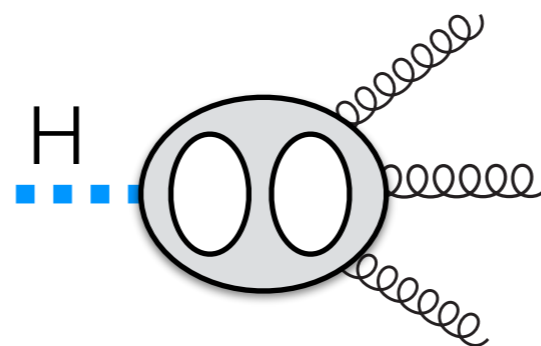


Maximal Transcendentality Principle



To which extend is this correspondence between correct?

Require more two-loop computations.



Two-loop QCD computation

Still a very challenging problem!

- 4-gluon amplitudes known many years ago

[Glover, Oleari, Tejeda-Yeomans 2001]
[Bern, De Freitas, Dixon 2002]

- 5-gluon results are still not fully known; numerical planar results last year

[Badger, Brønnum-Hansen, Hartanto, Peraro 2017]
[Abreu, Corderoa, Ita, Pagea, Zeng 2017]

.....

Our problem also involves high dimensional operators.

Strategy:

On-shell unitarity



Integration by parts

Content

- Motivations
- **Computation**
- Results
- Summary and outlook



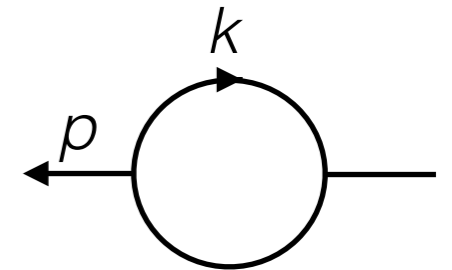
Integration by part reduction

[Chetyrkin, Tkachov 1981]

Integration by part (IBP): $\int d^D l_1 \dots d^D l_L \frac{\partial}{\partial l_i^\mu} (\text{integrand}) = 0.$

Solve a set of linear relations between different integrals.

Example: $J(a_1, a_2) := \int \frac{d^D k}{i\pi^{D/2}} \frac{1}{(-k^2 + m^2)^{a_1} (-(k+p)^2 + m^2)^{a_2}}$



→ $0 = \int \frac{d^D k}{i\pi^{D/2}} \frac{\partial}{\partial k^\mu} \left(k^\mu \frac{1}{(-k^2 + m^2)^{a_1} (-(k+p)^2 + m^2)^{a_2}} \right)$

→ $0 = (D - 2a_1 - a_2)J(a_1, a_2) - a_2 J(a_1 - 1, a_2 + 1) + 2m^2 a_1 J(a_1 + 1, a_2) + (2m^2 - p^2) a_2 J(a_1, a_2 + 1)$

$a_2 = 0$
→ $0 = (D - 2a_1)J(a_1, 0) + 2m^2 a_1 J(a_1 + 1, 0)$

$$J(1, 2) = \frac{(D - 2)}{2m^2(4m^2 - p^2)} J(1, 0) + \frac{(D - 3)}{4m^2 - p^2} J(1, 1)$$

2 Master Integral

Public packages:
Reduze 2, FIRE,
LiteRed, etc

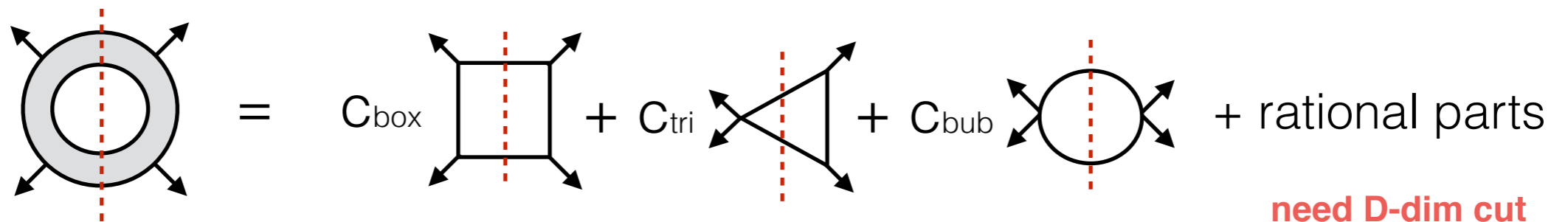
Unitarity method

As a replacement of Feynman diagram method:
construct the integrand from physical singularities,
i.e. poles or branch-cuts.

- At one-loop:

[Bern, Dixon, Dunbar, Kosower 1994]

[Britto, Cachazo, Feng 2004]



The diagrammatic equation shows the decomposition of a one-loop integrand. On the left is a shaded annulus with four external arrows and a vertical dashed red line. This is equal to the sum of three terms: C_{box} times a square loop with four external arrows and a vertical dashed red line; C_{tri} times a triangle loop with three external arrows and a vertical dashed red line; and C_{bub} times a bubble loop with two external arrows and a vertical dashed red line. The equation concludes with "+ rational parts".

need D-dim cut

Unitarity method

Challenges for higher loop QCD:

- need D-dimensional cuts (rational term issue)
- non-trivial to reconstruct full integrand (non-planar)
- need to further reduce the integrand, such as via IBP (sometimes IBP is the bottleneck)



Try new strategy

IBP for cut integrand

$$F^{(l)}|_{\text{cut}} = \sum_{\text{helicities}} F^{\text{tree}} \prod_j A_j^{\text{tree}} = \sum_i c_i M_i|_{\text{cut}}$$

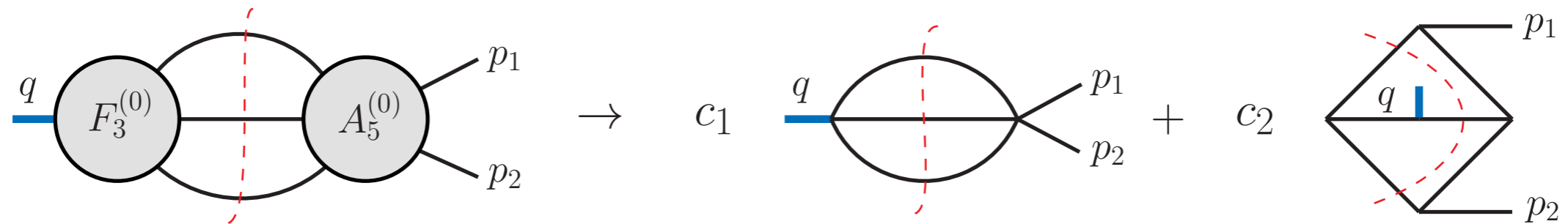
On-shell unitarity



Integration by parts

- D-dimensional cuts
- no need to reconstruct full integrand
- IBP is simplified

Example



- Tree by Feynman rules in D dimensions
- Helicity sum via contraction rule:

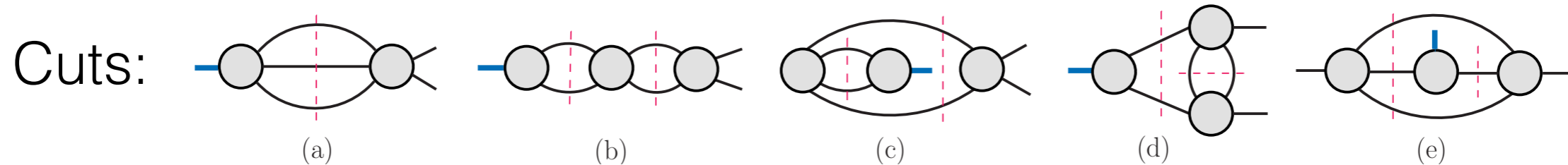
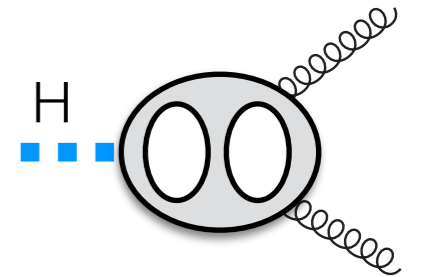
$$\sum_{\text{helicities}} \varepsilon_i^\mu \varepsilon_i^\nu = \eta^{\mu\nu} - \frac{q^\mu p_i^\nu + q^\nu p_i^\mu}{q \cdot p_i}$$

→ spinor helicity formalism for N=4 SYM

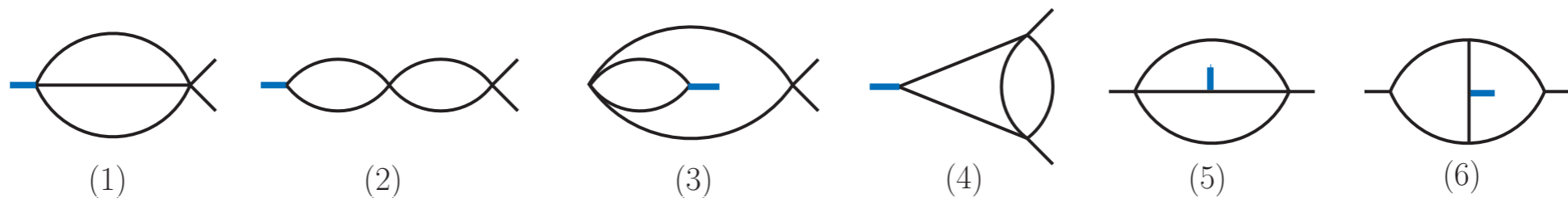
- Project to gauge invariant basis
- IBP reduction



2-loop 2-gluon



Master integrals:



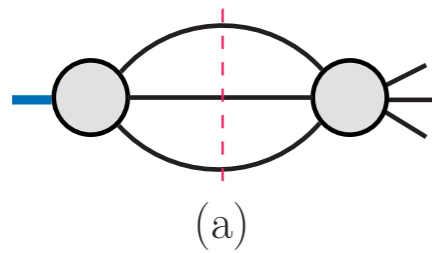
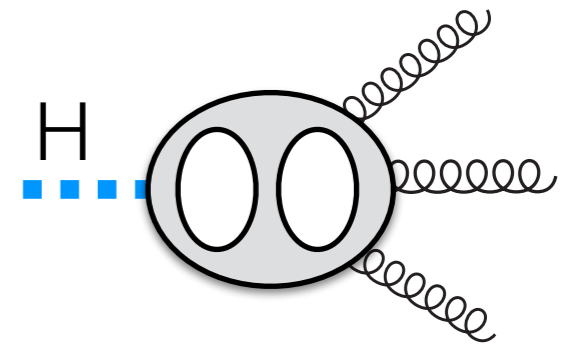
New feature (complication) of form factor:

'non-planar type' cuts appear for colour-planar part.

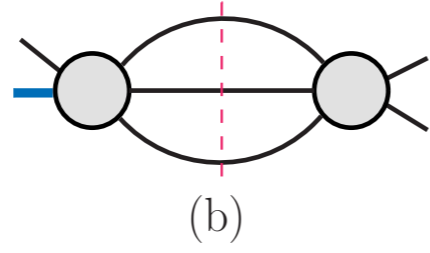
$$\text{Full result} = \left(\sum_{i=1}^4 c_i M_i + \frac{1}{2} \sum_{i=5,6} c_i M_i \right) + \text{perms}(p_1, p_2)$$

2-loop 3-gluon

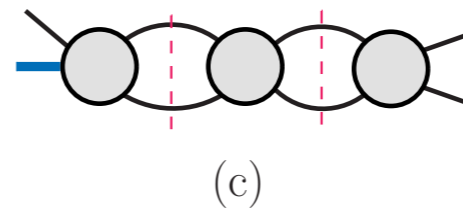
All cuts that are needed:



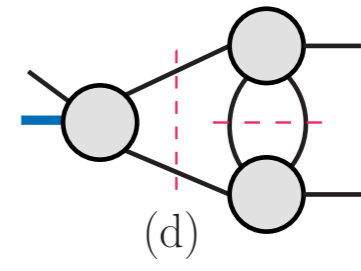
(a)



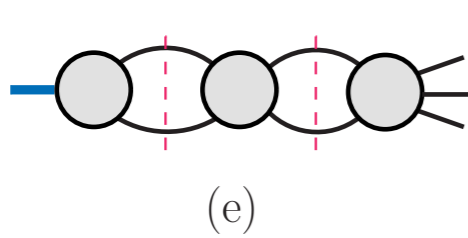
(b)



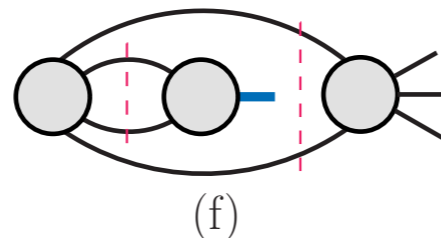
(c)



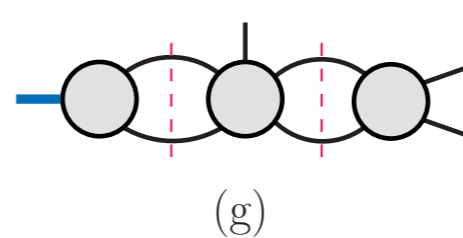
(d)



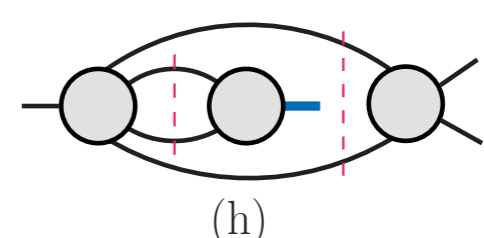
(e)



(f)

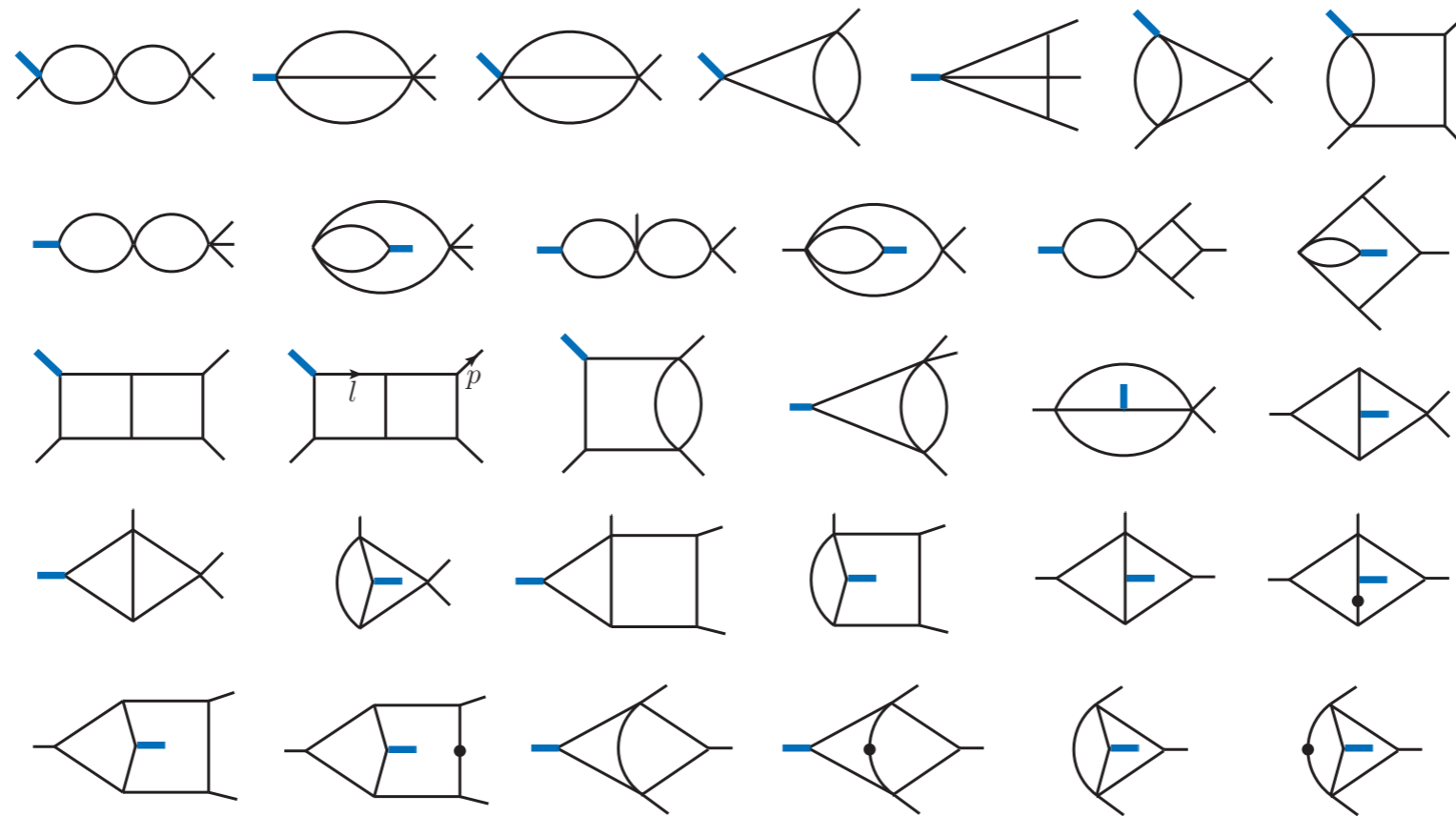


(g)



(h)

Master integrals



All analytic results are known, given in terms of
2d Harmonic polylogarithms.

[Gehrmann, Remiddi 2001]

Thus the results are given in terms of explicit functions.

Content

- Motivation
- Computation
- **Results**
- Summary and outlook



Divergence structure

UV divergences

both gauge coupling and local operator require renormalization

IR divergences

soft and collinear singularities from massless particle; universal!

Checks

- Consistent divergence structure: $\frac{1}{\epsilon^4}, \frac{1}{\epsilon^3}, \frac{1}{\epsilon^2}, \frac{1}{\epsilon}$

UV renormalisation (operator mixing) and Universal IR

$$Z_{\tilde{\mathcal{O}}_1}^{(2)} = -\frac{19}{24\epsilon^2} + \frac{25}{12\epsilon}, \quad \underline{\gamma_{\tilde{\mathcal{O}}_1}^{(2)}} = \frac{25}{3} \quad (\text{new result})$$

- Reproduce all known results, including the 2-loop Higgs to 3-gluon amplitudes in the heavy top limit $\mathcal{O}_0 = \text{tr}(G_{\mu\nu}G^{\mu\nu})$

- Results satisfy the linear relation:

$$F_{\mathcal{O}_2} = \frac{1}{2} q^2 F_{\mathcal{O}_0} - 4 g_{\text{YM}} F_{\mathcal{O}_1}$$

$$\mathcal{O}_0 = \text{tr}(G_{\mu\nu}G^{\mu\nu}).$$

$$\mathcal{O}_1 = \text{tr}(G_{\mu}^{\nu}G_{\nu}^{\rho}G_{\rho}^{\mu}),$$

$$\mathcal{O}_2 = \text{tr}(D_{\rho}G_{\mu\nu}D^{\rho}G^{\mu\nu}).$$

Degree 4:

$$\begin{aligned} & -2H[0, 1-x, 1, 0, y] + 2H[0, 1-x, 1, 0, z] + 2H[0, 1-y, 1, 0, x] - 2H[0, 1-y, 1, 0, z] - 2H[0, 1-z, 1, 0, x] + 2H[0, 1-z, 1, 0, y] - 2H[1-x, 0, 1, 0, y] + 2H[1-x, 0, 1, 0, z] - \\ & 2H[1-x, 1, 0, 0, y] + 2H[1-x, 1, 0, 0, z] - 3H[1-x, 1-x, 1, 0, y] + H[1-x, 1-x, 1, 0, z] + 2H[1-y, 0, 1, 0, x] - 2H[1-y, 0, 1, 0, z] + 2H[1-y, 1, 0, 0, x] - 2H[1-y, 1, 0, 0, z] + \\ & H[1-y, 1-y, 1, 0, x] - 3H[1-y, 1-y, 1, 0, z] - 2H[1-z, 0, 1, 0, x] + 2H[1-z, 0, 1, 0, y] - 2H[1-z, 1, 0, 0, x] + 2H[1-z, 1, 0, 0, y] - 3H[1-z, 1-z, 1, 0, x] + H[1-z, 1-z, 1, 0, y] + \\ & \frac{1}{6}\pi^2 \text{Log}[1-x]^2 - 2H[0, 1-x, 0, y] \text{Log}[x] + 2H[0, 1-x, 0, z] \text{Log}[x] - 2H[1-x, 0, 0, y] \text{Log}[x] + 2H[1-x, 0, 0, z] \text{Log}[x] - 3H[1-x, 1-x, 0, y] \text{Log}[x] + H[1-x, 1-x, 0, z] \text{Log}[x] - \\ & \text{Log}[1-x]^3 \text{Log}[x] - \frac{1}{6}\pi^2 \text{Log}[x]^2 + \frac{1}{2}\text{Log}[1-x]^2 \text{Log}[x]^2 - \frac{1}{6}\pi^2 \text{Log}[x] \text{Log}[1-y] + \frac{1}{6}\pi^2 \text{Log}[1-y]^2 + 2H[0, 1-y, 0, x] \text{Log}[y] - 2H[0, 1-y, 0, z] \text{Log}[y] + 2H[1-y, 0, 0, x] \text{Log}[y] - \\ & 2H[1-y, 0, 0, z] \text{Log}[y] + H[1-y, 1-y, 0, x] \text{Log}[y] - 3H[1-y, 1-y, 0, z] \text{Log}[y] - \frac{1}{6}\pi^2 \text{Log}[1-x] \text{Log}[y] + \frac{1}{3}\pi^2 \text{Log}[x] \text{Log}[y] + \text{Log}[1-x]^2 \text{Log}[x] \text{Log}[y] + \frac{1}{2}\text{Log}[1-x] \text{Log}[x]^2 \text{Log}[y] + \\ & \frac{1}{2}\text{Log}[x]^2 \text{Log}[1-y] \text{Log}[y] + \text{Log}[x] \text{Log}[1-y]^2 \text{Log}[y] - \text{Log}[1-y]^3 \text{Log}[y] - \frac{1}{6}\pi^2 \text{Log}[y]^2 + \frac{1}{2}\text{Log}[1-x] \text{Log}[x] \text{Log}[y]^2 - \frac{1}{4}\text{Log}[x]^2 \text{Log}[y]^2 - \frac{3}{2}\text{Log}[x] \text{Log}[1-y] \text{Log}[y]^2 + \frac{1}{2}\text{Log}[1-y]^2 \text{Log}[y]^2 - \\ & \frac{1}{6}\pi^2 \text{Log}[x] \text{Log}[1-z] - \frac{1}{6}\pi^2 \text{Log}[y] \text{Log}[1-z] + \frac{1}{6}\pi^2 \text{Log}[1-z]^2 - 2H[0, 1-z, 0, x] \text{Log}[z] + 2H[0, 1-z, 0, y] \text{Log}[z] - 2H[1-z, 0, 0, x] \text{Log}[z] + 2H[1-z, 0, 0, y] \text{Log}[z] - \\ & 3H[1-z, 1-z, 0, x] \text{Log}[z] + H[1-z, 1-z, 0, y] \text{Log}[z] - \frac{1}{6}\pi^2 \text{Log}[1-x] \text{Log}[z] + \frac{1}{3}\pi^2 \text{Log}[x] \text{Log}[z] + \text{Log}[1-x]^2 \text{Log}[x] \text{Log}[z] - \frac{3}{2}\text{Log}[1-x] \text{Log}[x]^2 \text{Log}[z] - \frac{1}{6}\pi^2 \text{Log}[1-y] \text{Log}[z] + \\ & \frac{1}{3}\pi^2 \text{Log}[y] \text{Log}[z] - 2\text{Log}[1-x] \text{Log}[x] \text{Log}[y] \text{Log}[z] + \text{Log}[x]^2 \text{Log}[y] \text{Log}[z] - 2\text{Log}[x] \text{Log}[1-y] \text{Log}[y] \text{Log}[z] + \text{Log}[1-y]^2 \text{Log}[y] \text{Log}[z] + \text{Log}[x] \text{Log}[y]^2 \text{Log}[z] + \frac{1}{2}\text{Log}[1-y] \text{Log}[y]^2 \text{Log}[z] + \\ & \frac{1}{2}\text{Log}[x]^2 \text{Log}[1-z] \text{Log}[z] - 2\text{Log}[x] \text{Log}[y] \text{Log}[1-z] \text{Log}[z] + \frac{1}{2}\text{Log}[y]^2 \text{Log}[1-z] \text{Log}[z] + \text{Log}[x] \text{Log}[1-z]^2 \text{Log}[z] + \text{Log}[y] \text{Log}[1-z]^2 \text{Log}[z] - \text{Log}[1-z]^3 \text{Log}[z] - \frac{1}{6}\pi^2 \text{Log}[z]^2 + \\ & \frac{1}{2}\text{Log}[1-x] \text{Log}[x] \text{Log}[z]^2 - \frac{1}{4}\text{Log}[x]^2 \text{Log}[z]^2 + \text{Log}[x] \text{Log}[y] \text{Log}[z]^2 + \frac{1}{2}\text{Log}[1-y] \text{Log}[y] \text{Log}[z]^2 - \frac{1}{4}\text{Log}[y]^2 \text{Log}[z]^2 + \frac{1}{2}\text{Log}[x] \text{Log}[1-z] \text{Log}[z]^2 - \frac{3}{2}\text{Log}[y] \text{Log}[1-z] \text{Log}[z]^2 + \\ & \frac{1}{2}\text{Log}[1-z]^2 \text{Log}[z]^2 - 2\text{Log}[1-x]^2 \text{PolyLog}[2, 1-x] - 3\text{Log}[1-x] \text{Log}[y] \text{PolyLog}[2, 1-x] + 5\text{Log}[1-x] \text{Log}[z] \text{PolyLog}[2, 1-x] - \text{Log}[1-x]^2 \text{PolyLog}[2, x] + \text{Log}[1-x] \text{Log}[y] \text{PolyLog}[2, x] + \\ & \frac{1}{2}\text{Log}[y]^2 \text{PolyLog}[2, x] + \text{Log}[1-x] \text{Log}[z] \text{PolyLog}[2, x] - 2\text{Log}[y] \text{Log}[z] \text{PolyLog}[2, x] + \frac{1}{2}\text{Log}[z]^2 \text{PolyLog}[2, x] + 5\text{Log}[x] \text{Log}[1-y] \text{PolyLog}[2, 1-y] - 2\text{Log}[1-y]^2 \text{PolyLog}[2, 1-y] - \\ & 3\text{Log}[1-y] \text{Log}[z] \text{PolyLog}[2, 1-y] + \frac{1}{6}\pi^2 \text{PolyLog}\left[2, -\frac{x}{-1+y}\right] - \text{Log}[1-y] \text{Log}[y] \text{PolyLog}\left[2, -\frac{x}{-1+y}\right] - \frac{1}{2}\text{Log}[y]^2 \text{PolyLog}\left[2, -\frac{x}{-1+y}\right] + \frac{1}{2}\text{Log}[x]^2 \text{PolyLog}[2, y] + \text{Log}[x] \text{Log}[1-y] \text{PolyLog}[2, y] \\ & \text{Log}[1-y]^2 \text{PolyLog}[2, y] - 2\text{Log}[x] \text{Log}[z] \text{PolyLog}[2, y] + \text{Log}[1-y] \text{Log}[z] \text{PolyLog}[2, y] + \frac{1}{2}\text{Log}[z]^2 \text{PolyLog}[2, y] - \text{PolyLog}\left[2, -\frac{x}{-1+y}\right] \text{PolyLog}[2, y] + \frac{1}{6}\pi^2 \text{PolyLog}\left[2, \frac{y}{1-x}\right] - \\ & \text{Log}[1-x] \text{Log}[x] \text{PolyLog}\left[2, \frac{y}{1-x}\right] + \frac{3}{2}\text{Log}[x]^2 \text{PolyLog}\left[2, \frac{y}{1-x}\right] - \text{PolyLog}[2, x] \text{PolyLog}\left[2, \frac{y}{1-x}\right] + \frac{1}{6}\pi^2 \text{PolyLog}\left[2, \frac{y}{1-x}\right] - \text{Log}[1-z] \text{Log}[z] \text{PolyLog}\left[2, \frac{y}{1-z}\right] - \frac{1}{2}\text{Log}[z]^2 \text{PolyLog}\left[2, \frac{y}{1-z}\right] - \\ & 3\text{Log}[x] \text{Log}[1-z] \text{PolyLog}[2, 1-z] + 5\text{Log}[y] \text{Log}[1-z] \text{PolyLog}[2, 1-z] - 2\text{Log}[1-z]^2 \text{PolyLog}[2, 1-z] + \frac{1}{6}\pi^2 \text{PolyLog}\left[2, -\frac{x}{-1+z}\right] - \text{Log}[1-z] \text{Log}[z] \text{PolyLog}\left[2, -\frac{x}{-1+z}\right] + \\ & \frac{3}{2}\text{Log}[z]^2 \text{PolyLog}\left[2, -\frac{x}{-1+z}\right] + \frac{1}{2}\text{Log}[x]^2 \text{PolyLog}[2, z] - 2\text{Log}[x] \text{Log}[y] \text{PolyLog}[2, z] + \frac{1}{2}\text{Log}[y]^2 \text{PolyLog}[2, z] + \text{Log}[x] \text{Log}[1-z] \text{PolyLog}[2, z] + \text{Log}[y] \text{Log}[1-z] \text{PolyLog}[2, z] - \\ & \text{Log}[1-z]^2 \text{PolyLog}[2, z] - \text{PolyLog}\left[2, \frac{y}{1-z}\right] \text{PolyLog}[2, z] - \text{PolyLog}\left[2, -\frac{x}{-1+z}\right] \text{PolyLog}[2, z] + \frac{1}{6}\pi^2 \text{PolyLog}\left[2, \frac{z}{1-x}\right] - \text{Log}[1-x] \text{Log}[x] \text{PolyLog}\left[2, \frac{z}{1-x}\right] - \frac{1}{2}\text{Log}[x]^2 \text{PolyLog}\left[2, \frac{z}{1-x}\right] - \\ & \text{PolyLog}[2, x] \text{PolyLog}\left[2, \frac{z}{1-x}\right] + \frac{1}{6}\pi^2 \text{PolyLog}\left[2, \frac{z}{1-y}\right] - \text{Log}[1-y] \text{Log}[y] \text{PolyLog}\left[2, \frac{z}{1-y}\right] + \frac{3}{2}\text{Log}[y]^2 \text{PolyLog}\left[2, \frac{z}{1-y}\right] - \text{PolyLog}[2, y] \text{PolyLog}\left[2, \frac{z}{1-y}\right] + 4\text{Log}[1-x] \text{PolyLog}[3, 1-x] + \\ & 3\text{Log}[y] \text{PolyLog}[3, 1-x] - 5\text{Log}[z] \text{PolyLog}[3, 1-x] + 4\text{Log}[y] \text{PolyLog}[3, x] - 4\text{Log}[z] \text{PolyLog}[3, x] - 5\text{Log}[x] \text{PolyLog}[3, 1-y] + 4\text{Log}[1-y] \text{PolyLog}[3, 1-y] + 3\text{Log}[z] \text{PolyLog}[3, 1-y] - \\ & 4\text{Log}[x] \text{PolyLog}[3, y] + 4\text{Log}[z] \text{PolyLog}[3, y] + 3\text{Log}[x] \text{PolyLog}[3, 1-z] - 5\text{Log}[y] \text{PolyLog}[3, 1-z] + 4\text{Log}[1-z] \text{PolyLog}[3, 1-z] + 4\text{Log}[x] \text{PolyLog}[3, z] - 4\text{Log}[y] \text{PolyLog}[3, z] - \\ & 4\text{PolyLog}[4, 1-x] - 4\text{PolyLog}[4, 1-y] - 4\text{PolyLog}[4, 1-z] - 2\text{PolyLog}[2, 2, x] - 2\text{PolyLog}[2, 2, y] - 2\text{PolyLog}[2, 2, z] + \text{Log}[x] \text{Zeta}[3] + \text{Log}[y] \text{Zeta}[3] + \text{Log}[z] \text{Zeta}[3] - \frac{\pi^4}{16} \end{aligned}$$

Finite remainder

Simplify via “**symbol**” for transcendental functions

$$\begin{aligned}\Omega_{\mathcal{O}_1;4}^{(2)} = & -\frac{3}{2}\text{Li}_4(u) + \frac{3}{4}\text{Li}_4\left(-\frac{uv}{w}\right) - \frac{3}{2}\log(w)\text{Li}_3\left(-\frac{u}{v}\right) + \frac{\zeta_2}{8} [5\log^2(u) - 2\log(v)\log(w)] \\ & + \frac{\log^2(u)}{32} [\log^2(u) + 2\log^2(v) - 4\log(v)\log(w)] - \frac{1}{4}\zeta_4 - \frac{1}{2}\zeta_3 \log(-q^2) + \text{perms}(u, v, w)\end{aligned}$$



$$\Omega_{\mathcal{O}_1;4}^{(2)} = \Omega_{\mathcal{O}_1;4}^{(2), \mathcal{N}=4}$$

for N=4 result see:
Brandhuber, Kostacinska, Penante,
Travaglini 2017]

It also appears as a universal function for length-3 operators.

[Brandhuber, Kostacinska, Penante, Travaglini, Wen, Young 2014, 2016]

[Loebbert, Nandan, Sieg, Wilhelm, GY 2015, 2016]



Finite remainder

Weight-3 part:

$$\Omega_{\mathcal{O}_1;3}^{(2)} = \left(1 + \frac{u}{w}\right) T_3 + \frac{143}{72} \zeta_3 - \frac{11}{24} \zeta_2 \log(-u q^2) + \text{perms}(u, v, w)$$

$$\Omega_{\mathcal{O}_1;3}^{(2), \mathcal{N}=4} = \left(1 + \frac{u}{w}\right) T_3 + \text{perms}(u, v, w)$$

$$T_3 := \left[-\text{Li}_3\left(-\frac{u}{w}\right) + \log(u) \text{Li}_2\left(\frac{v}{1-u}\right) - \frac{1}{2} \log(1-u) \log(u) \log\left(\frac{w^2}{1-u}\right) + \frac{1}{2} \text{Li}_3\left(-\frac{uv}{w}\right) + \frac{1}{12} \log^3(w) \right. \\ \left. + \frac{1}{2} \log(u) \log(v) \log(w) + (u \leftrightarrow v) \right] + \text{Li}_3(1-v) - \text{Li}_3(u) + \frac{1}{2} \log^2(v) \log\left(\frac{1-v}{u}\right) - \zeta_2 \log\left(\frac{uv}{w}\right).$$

T_3 function is also a building block appearing in many form factors in N=4 SYM: $T_3 = -\left(R_i^{(2)}\right)_{XXY}^{XYX} \Big|_3 - \zeta_2 \log(u)$

[Loebbert, Nandan, Sieg, Wilhelm, GY 2015]



Finite remainder

Degree 2 to 0:

$$\Omega_{\mathcal{O}_{1;2}}^{(2)} = \left\{ \left(\frac{u^2}{w^2} + \frac{v^2}{w^2} - 1 \right) \left[\text{Li}_2(1-u) + \frac{1}{2} \log(u) \log(v) - \frac{1}{2} \zeta_2 \right] - \frac{55}{48} \log^2(u) + \frac{73}{72} \log(u) \log(v) + \frac{23}{6} \zeta_2 + \text{perms}(u, v, w) \right\} - \frac{19}{36} \log(uvw) \log(-q^2) - \frac{19}{24} \log^2(-q^2).$$

$$\Omega_{\mathcal{O}_{1;1}}^{(2)} = \left(\frac{119}{18} + \frac{v}{w} + \frac{u^2}{2vw} \right) \log(u) + \left(\frac{119}{18} - \frac{1}{3uvw} \right) \log(-q^2) + \text{perms}(u, v, w)$$

$$\Omega_{\mathcal{O}_{2;0}}^{(2)} = \frac{487}{72} \frac{1}{uvw} - \frac{14075}{216}.$$

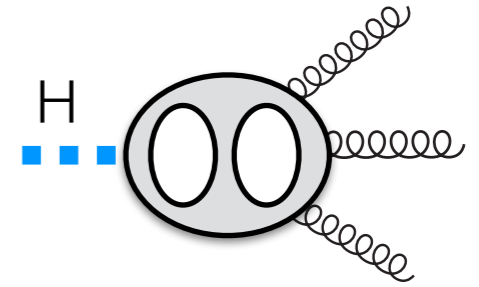
Lower transcendental terms with rational kinematic coefficients are also identical to the N=4 results.

Content

- Motivation
- Computation
- Results
- **Summary and outlook**



Summary



- Two-loop Higgs amplitudes with dim-7 operators
- Efficient method based on on-shell unitarity and IBP
- Simple analytic result which provide further evidence of transcendentality principle



Outlook

- Form factors with more general operators and more gluons
- Understand better the maximal transcendentality principle (more examples)
- Origin of the simplicity?
Is there a way to understand it directly? Bootstraps??



Thank you for your attention!



From function to “Symbol”

Recursion definition of “Symbol”:

$$df_k = \sum_i f_{k-1}^i d\text{Log}(R_i), \quad \text{Symbol}(f_k) = \sum_i \text{Symbol}(f_{k-1}^i) \otimes R_i$$

Some examples:

Function	Differential	symbol
R	d R	0
log(R)	d log(R)	R
log(R1)log(R2)	logR1 dlogR2+logR2 dlogR1	R1 ⊗ R2 + R2 ⊗ R1
Li ₂ (R)	Li ₁ (R) dlogR	-(1-R) ⊗ R

Symbol contains analytics properties of functions, e.g. branch cuts.

Symbol

Properties:

$$R_1 \otimes \dots \otimes (c R_i) \otimes \dots \otimes R_n = R_1 \otimes \dots \otimes R_i \otimes \dots \otimes R_n \quad c = \text{const}$$

$$R_1 \otimes \dots \otimes (R_i R_j) \otimes \dots \otimes R_n = R_1 \otimes \dots \otimes R_i \otimes \dots \otimes R_n + R_1 \otimes \dots \otimes R_j \otimes \dots \otimes R_n$$

Make it easy to prove non-trivial identities, e.g.:

$$\text{Li}_2\left(\frac{x}{1-y}\right) + \text{Li}_2\left(\frac{y}{1-x}\right) - \text{Li}_2(x) - \text{Li}_2(y) - \text{Li}_2\left(\frac{xy}{(1-x)(1-y)}\right) = \text{Log}(1-x)\text{Log}(1-y)$$

($x < 1$ and $y < 1$)

$$\text{Symbol (LHS)} = (1-x) \otimes (1-y) + (1-y) \otimes (1-x)$$

[back to main content](#)

应用

复杂结果 \longrightarrow symbol \longrightarrow 简单结果

更进一步：

从一些基本约束出发，
直接求解symbol

\longrightarrow 简单结果

(这类方法称为Bootstrap)

Gauge invariant basis projection

$$F_n(\varepsilon_i, p_i, l_a)|_{\text{cut}} = \sum_{\alpha} f_n^{\alpha}(p_i, l_a) B_{\alpha}$$

$$f_n^{\alpha}(p_i, l_a) = B^{\alpha} \circ F_n(\varepsilon_i, p_i, l_a)$$

$$\varepsilon_i^{\mu} \circ \varepsilon_i^{\nu} \equiv \sum_{\text{helicities}} \varepsilon_i^{\mu} \varepsilon_i^{\nu} = \eta^{\mu\nu} - \frac{q^{\mu} p_i^{\nu} + q^{\nu} p_i^{\mu}}{q \cdot p_i}$$

$$B^{\alpha} \circ B_{\beta} = \delta_{\beta}^{\alpha}, \quad B_{\alpha} = G_{\alpha\beta} B^{\beta}, \quad G_{\alpha\beta} = B_{\alpha} \circ B_{\beta}$$

Three gluon case:

$$B_1 = A_1 C_{23}, \quad B_2 = A_2 C_{31}, \quad B_3 = A_3 C_{12}, \quad B_4 = A_1 A_2 A_3$$

$$\text{where } A_i = \frac{\varepsilon_i \cdot p_j}{p_i \cdot p_j} - \frac{\varepsilon_i \cdot p_k}{p_i \cdot p_k}, \quad C_{ij} = \varepsilon_i \cdot \varepsilon_j - \frac{(p_i \cdot \varepsilon_j)(p_j \cdot \varepsilon_i)}{p_i \cdot p_j}$$

Two gluon case:

$$B_0 = C_{12}$$

UV renormalization

Coupling constant renormalisation:

$$\alpha_0 = \alpha_s S_\epsilon^{-1} \frac{\mu^{2\epsilon}}{\mu_0^{2\epsilon}} \left[1 - \frac{\beta_0 \alpha_s}{\epsilon 4\pi} + \left(\frac{\beta_0^2}{\epsilon^2} - \frac{\beta_1}{2\epsilon} \right) \left(\frac{\alpha_s}{4\pi} \right)^2 + \mathcal{O}(\alpha_s^3) \right]$$

Renormalisation constant Z for the operators:

$$\mathcal{O}_I^b \longrightarrow Z_{IJ} \mathcal{O}_J^b \quad Z = 1 + \sum_{l=1}^{\infty} \left(\frac{\alpha_s}{4\pi} \right)^l Z^{(l)}$$

anomalous dimension

$$\longrightarrow \gamma = \mu \frac{\partial}{\partial \mu} \log Z$$

Renormalized form factor: $F = g_s^x S_\epsilon^{-x/2} \sum_{l=0}^{\infty} \left(\frac{\alpha_s}{4\pi} \right)^l F^{(l)}$

$$F^{(0)} = F_b^{(0)},$$

$$F^{(1)} = S_\epsilon^{-1} F_b^{(1)} + \left(Z^{(1)} - \frac{x \beta_0}{2 \epsilon} \right) F_b^{(0)},$$

$$F^{(2)} = S_\epsilon^{-2} F_b^{(2)} + S_\epsilon^{-1} \left[Z^{(1)} - \left(1 + \frac{x}{2} \right) \frac{\beta_0}{\epsilon} \right] F_b^{(1)} \\ + \left[Z^{(2)} - \frac{x \beta_0}{2 \epsilon} Z^{(1)} + \frac{x^2 + 2x \beta_0^2}{8 \epsilon^2} - \frac{x \beta_1}{4 \epsilon} \right] F_b^{(0)}$$

IR subtraction

Universal IR structure:

[Catani 1998]

$$F^{(1)} = I^{(1)}(\epsilon)F^{(0)} + F^{(1),\text{fin}} + \mathcal{O}(\epsilon),$$

$$F^{(2)} = I^{(2)}(\epsilon)F^{(0)} + I^{(1)}(\epsilon)F^{(1)} + F^{(2),\text{fin}} + \mathcal{O}(\epsilon)$$

where

$$I^{(1)}(\epsilon) = -\frac{e^{\gamma_E \epsilon}}{\Gamma(1-\epsilon)} \left(\frac{N_c}{\epsilon^2} + \frac{\beta_0}{2\epsilon} \right) \sum_{i=1}^n (-s_{i,i+1})^{-\epsilon},$$

$$\begin{aligned} I^{(2)}(\epsilon) = & -\frac{1}{2} [I^{(1)}(\epsilon)]^2 - \frac{\beta_0}{\epsilon} I^{(1)}(\epsilon) \\ & + \frac{e^{-\gamma_E \epsilon} \Gamma(1-2\epsilon)}{\Gamma(1-\epsilon)} \left[\frac{\beta_0}{\epsilon} + \left(\frac{67}{9} - \frac{\pi^2}{3} \right) N_c \right] I^{(1)}(2\epsilon) \\ & + n \frac{e^{\gamma_E \epsilon}}{\epsilon \Gamma(1-\epsilon)} \left[\left(\frac{\zeta_3}{2} + \frac{5}{12} + \frac{11\pi^2}{144} \right) N_c^2 \right]. \end{aligned}$$



Anomalous dimension

Operator mixing:

$$\begin{aligned}
 F_{\mathcal{O}_1}^{(2)}(1^-, 2^-, 3^-)|_{Z^{(2)\text{-part}}} &= F_{\mathcal{O}_1}^{(0)}(1^-, 2^-, 3^-) \left(-\frac{19}{24\epsilon^2} + \frac{25}{12\epsilon} - \frac{1}{uvw} \frac{1}{\epsilon} \right) \\
 &= \left(-\frac{19}{24\epsilon^2} + \frac{25}{12\epsilon} \right) F_{\mathcal{O}_1}^{(0)}(1^-, 2^-, 3^-) - \frac{F_{\tilde{\mathcal{O}}_2}^{(0)}(1^-, 2^-, 3^-)}{\epsilon}
 \end{aligned}$$

$$F_{\mathcal{O}_1}^{(2)}(1^-, 2^-) = F_{\tilde{\mathcal{O}}_2}^{(0)}(1^-, 2^-) \left(-\frac{1}{\epsilon} + 2 \log s_{12} - \frac{487}{72} \right) + \mathcal{O}(\epsilon^1)$$

Eigen-operators:

$$\tilde{\mathcal{O}}_2 = -\frac{3}{2}(\mathcal{O}_2 + 8g_{\text{YM}} \mathcal{O}_1) = -\frac{3}{4} \partial^2 \mathcal{O}_0$$

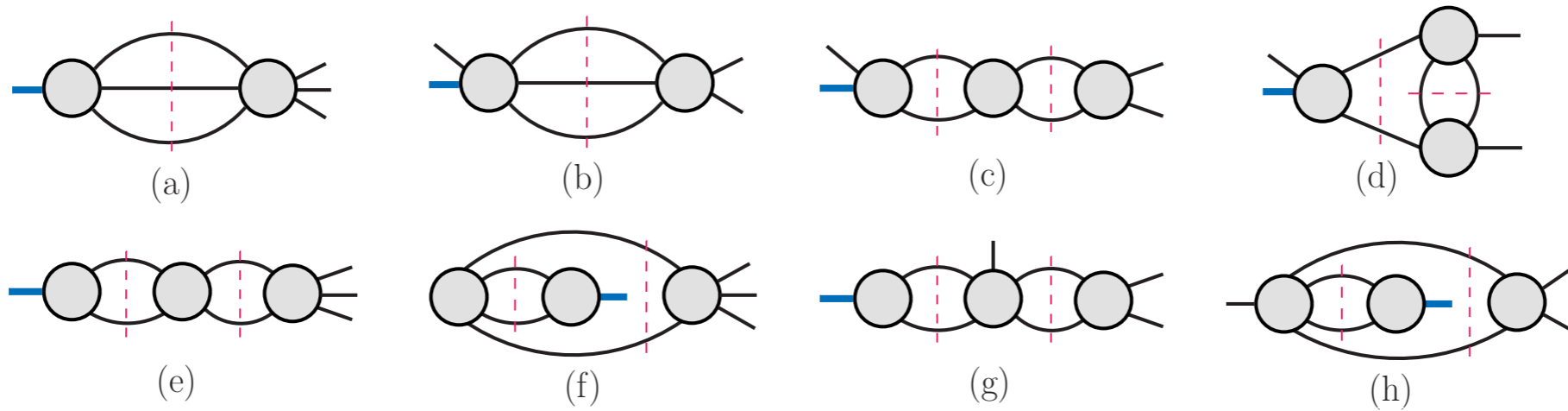
$$\tilde{\mathcal{O}}_1 = \mathcal{O}_1 + \frac{1}{\epsilon} \frac{1}{g_{\text{YM}}} \left(\frac{\alpha_s}{4\pi} \right)^2 \tilde{\mathcal{O}}_2$$

Anomalous dimension: $\gamma = \mu \frac{\partial}{\partial \mu} \log Z$

$$Z_{\tilde{\mathcal{O}}_1}^{(2)} = -\frac{19}{24\epsilon^2} + \frac{25}{12\epsilon}, \quad \underline{\gamma_{\tilde{\mathcal{O}}_1}^{(2)} = \frac{25}{3}}$$

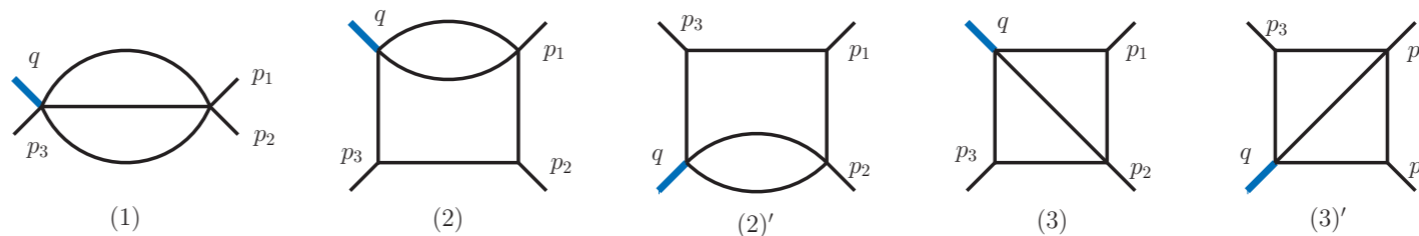


2-loop 3-gluon

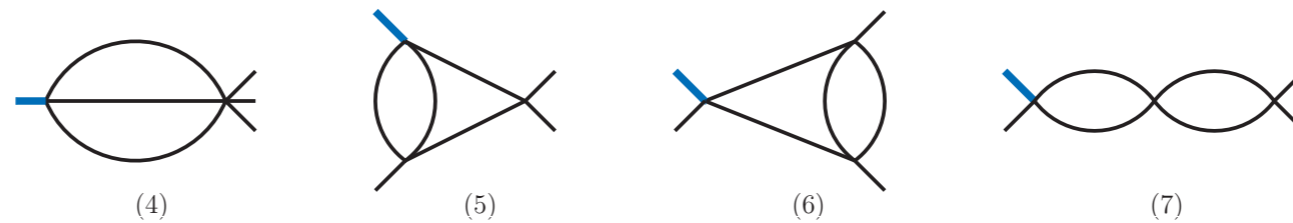


$\text{Tr}(F^3)$

cut (b):



Other MIs:



$$\text{Full result} = \frac{1}{2} \left(\sum_{i=1}^7 c_i M_i + \sum_{i=2,5} c_i M_i \right) + \text{perms}(p_1, p_2, p_3)$$