### Toward precision measurement of neutrinos

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# $\nu$ oscillation

### The solar neutrino problem:

 $\nu$  produced in Sun: pp  $\nu$ :  $p + p \rightarrow^2 \mathbf{H} + e^+ + \nu_e$ Boron  $\nu$ :  ${}^8\mathbf{B} \rightarrow^8 \mathbf{Be}^* + e^+ + \nu_e$ etc. ...

At low energy part of the spectrum about  $\frac{2}{3} \nu_e$  flux detected

At high energy part of the spectrum about  $\frac{1}{3} \nu_e$  flux detected

# Missing of atmospheric $\nu_{\mu}$ also discovered



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# u oscillation

When flavor eigenstates are not the mass eigenstates ( $\nu_1$ , $\nu_2$ ),

$$\begin{split} i\frac{d}{dx} \begin{pmatrix} \nu_e \\ \nu'_{\mu} \end{pmatrix} &= \frac{1}{4E} \begin{pmatrix} -\Delta m^2 \cos 2\theta & \Delta m^2 \sin 2\theta \\ \Delta m^2 \sin 2\theta & \Delta m^2 \cos 2\theta \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu'_{\mu} \end{pmatrix}; \\ \begin{pmatrix} \nu_e \\ \nu'_{\mu} \end{pmatrix} &= \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}. \end{split}$$

At distance x one obtains

$$\begin{aligned} |\nu_e(x)\rangle &= \cos\theta e^{-iE_1x} |\nu_1\rangle + \sin\theta e^{-iE_2x} |\nu_2\rangle; \\ |\langle\nu_e|\nu_e(x)\rangle|^2 &= 1 - \sin^2 2\theta \sin^2 \frac{\Delta m^2 x}{4E}. \end{aligned}$$

Averaging over phases or coherence lost(happened for  $\nu_{\odot}$ )

$$P_{ee} = 1 - \frac{1}{2}\sin^2 2\theta = \frac{1}{2}(1 + \cos^2 2\theta).$$

 $P_{ee} \ge 1/2$  in vacuum oscillation.

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### L. Wolfenstein 1978, Mikheyev and Smirnov 1986.

Coherent forward scattering by medium modifies dispersion relation:

$$E=E_k+V,$$

 $E_k$ , the kinetic energy; V, the potential energy. Examples include

- the optics : the case of electromagnetism;
- neutron optics: the case of strong interaction;
- MSW: the case of weak interaction,  $V_e = \sqrt{2}G_F N_e$ :

$$i\frac{d}{dx}\begin{pmatrix}\nu_{e}\\\nu'_{\mu}\end{pmatrix} = \begin{pmatrix}-\frac{\Delta m^{2}}{4E}\cos 2\theta + V_{e} & \frac{\Delta m^{2}}{4E}\sin 2\theta\\\frac{\Delta m^{2}}{4E}\sin 2\theta & \frac{\Delta m^{2}}{4E}\cos 2\theta\end{pmatrix}\begin{pmatrix}\nu_{e}\\\nu'_{\mu}\end{pmatrix};$$
$$\begin{pmatrix}\nu_{e}\\\nu'_{\mu}\end{pmatrix} = \begin{pmatrix}\cos \theta_{m} & \sin \theta_{m}\\-\sin \theta_{m} & \cos \theta_{m}\end{pmatrix}\begin{pmatrix}\nu_{m1}\\\nu_{m2}\end{pmatrix}.$$

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### LMA MSW solution:

- $$\begin{split} \Delta m^2 &= (3.8-10)\times 10^{-5} {\rm eV}^2, \\ \tan^2\theta &= 0.32-0.47 \end{split}$$
- ► *L*<sub>osc</sub> ~ 200 km.
- ► For  $E \gtrsim 7$  MeV,  $P_{ee} = \frac{1}{2}(1 - \cos 2\theta) \approx 0.3$ ;
- ► for  $E \lesssim 1$  MeV,  $P_{ee} = \frac{1}{2}(1 + \cos^2 2\theta) \approx 0.6$ .
- ►  $2EV_e/\Delta m_{21}^2 \lesssim 0.08$  in Earth; matter effect small in Earth.

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The present status Oscillation of  $\nu$  confirmed, LMA MSW established

Solar ν experiment(Homestake, Super-K, SNO, etc.), plus long baseline ν experiment(Kamland):

$$\Delta m^2_{21}pprox 0.8 imes 10^{-4}~{
m eV}^2,~{
m tan}^2 heta_{12}pprox 0.4$$

• Atmospheric  $\nu$  experiments(Super-K, SNO, etc.):

$$|\Delta m^2_{32}| pprox 3 imes 10^{-3}~{
m eV}^2,~{
m tan}^2 heta_{23} pprox 1.0$$

Reactor v experiment(Chooz, etc.):

$$\sin^2 2 heta_{13} \lesssim 0.1$$

### In the future we will measure

- $\theta_{13}$  (reactor  $\nu$ , long baseline exp)
- mass hierarchy(long baseline exp)
- CP violation in neutrinos(long baseline exp)
- absolute mass scale
- the Earth matter effect(long baseline, solar ν exp)
- nature of neutrino mass
- magnetic moment of neutrino, etc.

Earth matter effect is important to most measurements

# Earth matter effect

PReliminary Earth Model(PREM)



Matter profile very complicated

Earth matter density has many layers and changes

- sharply between two layers
- slowly in a layer: density height

$$h = \left(\frac{dV}{dx}\right)^{-1} \sim R_{Earth}$$

Earth matter effect is a challenge to future precision measurement

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# Quest to understand the effect of the complex Earth matter in neutrino oscillation

### Earth matter effect

Oscillation length tells how fast  $\nu$  oscillates

$$L_{21} = rac{4\pi E}{\Delta m_{21}^2}, \ \ L_{32} = rac{4\pi E}{|\Delta m_{32}^2|}$$

Note that available measurements tell us

• for  $E \sim 10$  MeV (solar, supernova neutrinos)

 $L_{21} \sim \textbf{hundreds } \textbf{km}, \ \ L_{32} \sim \textbf{few } \textbf{km},$ 

 $L_{21}, L_{32} \ll h$ • for  $E \gtrsim 500$  MeV

$$L_{21}\gtrsim h$$

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### An important lesson:

oscillation well approximated by 1-2 plus 2-3 oscillation i.e., oscillation in matter and in vacuum if  $\theta_{13}$  small

$$H = \frac{1}{2E} U^* \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^{\dagger} + \begin{pmatrix} V_e & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Rotation in 2-3 sector does not change the potential term

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Crucial to compare  $L_{21}$  and h to understand matter effect;  $L_{31}$  not crucial when  $\theta_{13}$  small

For  $E \lesssim 30$  MeV,  $L_{21} \ll h$  (solar and supernova neutrinos)

 $\nu_{e}$  survival well reduced to 1-2 oscillation

$$i\frac{d}{dx}\begin{pmatrix}\psi_{1m}\\\psi_{2m}\end{pmatrix} = \begin{pmatrix}-\frac{\Delta(x)}{4E} & -i\dot{\theta}_m(x)\\i\dot{\theta}_m(x) & \frac{\Delta(x)}{4E}\end{pmatrix}\begin{pmatrix}\psi_{1m}\\\psi_{2m}\end{pmatrix},$$

 $(\psi_{m1}, \psi_{m2})^T$ , neutrino mass state in matter.

### The survival probability of $\nu_{\odot}$ on the Earth

$$P_{ee} = \frac{1}{2}(1 + \cos 2\theta_m(x_0)\cos 2\theta_{12}) - \cos 2\theta_m(x_0)f_{reg}.$$

 $f_{reg}=P(
u_2
ightarrow
u_e)-\sin^2 heta_{12}$ , the regeneration by the Earth

### Adiabatic perturbation theory (de Holanda, Liao, Smirnov, 2004)

Search for the solution of the following form

$$\begin{pmatrix} \psi_{1m}(x) \\ \psi_{2m}(x) \end{pmatrix} = \begin{pmatrix} e^{i\Phi(x)} & c(x)e^{-i\Phi(x)} \\ -c^*(x)e^{i\Phi(x)} & e^{-i\Phi(x)} \end{pmatrix} \begin{pmatrix} \psi_{1m}(x_0) \\ \psi_{2m}(x_0) \end{pmatrix}$$

$$\Phi(x) = \frac{1}{4E} \int_{x_0}^x dx' \Delta(x').$$

where  $|c(x)| \ll 1$  is supposed to hold(adiabatic perturbation) We get

$$c(x) = -\int_{x_0}^{x} dx' \frac{d\theta_m(x')}{dx'} \exp\left[-i \int_{x}^{x'} dx'' \frac{\Delta(x'')}{2E}\right]$$

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#### $\nu$ regeneration in the Earth



$$f_{reg} = \frac{2E \sin^2 2\theta}{\Delta m^2} \sin \Phi_0 \sum \Delta V_i \sin \Phi_i,$$
  
$$\Phi_i = \int_{-L_i/2}^{L_i/2} dx \frac{\Delta(x)}{4E}.$$

- Leading contributions are from potential jumps between layers
- Analytic and numerical computations perfectly agree
- Oscillatory pattern is well understood using the adiabatic perturbation theory

#### $\nu$ regeneration in the Earth



#### **Properties:**

- Averaging over energy tremendously simplifies the oscillation pattern
- Neutrinos of horizontal direction still give complicated oscillation pattern in Earth regeneration



Extension to asymmetric matter profile

$$f_{reg} = -\frac{E\sin^2 2\theta_{12}}{\Delta m_{21}^2} \sum_{i=0}^k \delta V_i \cos 2\phi_i, \ \phi_i = \int_{x_i}^{x_k} dx \frac{\Delta(x)}{4E}, \ i = 0, \cdots, k$$

Main uncertainties in the Earth matter effects are from contributions of small structures close to the detectors.

For  $E \gtrsim 0.5$  GeV,  $L_{21} \gtrsim h$ ,  $L_{31} > \mathbf{or} < h$ 

 $\boldsymbol{\nu}$  oscillation seems very complicated.

It turns out that

1) we can have a perturbation theory which perfectly describes the oscillation pattern

2) neutrino oscillation can be substantially simplified in the interested energy range

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### A nice example:

For E > 10 GeV ( $L_{21}, L_{31} > h$ ), neutrinos

- can not see the structure of the Earth very well
- see baseline dependent average potential in Earth

Earth matter effect is well described by a formulation using the baseline dependent average potential

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Liao, 2008; Liao, 2008

### A baseline dependent perturbation theory

$$H = \bar{H} + \delta H, \quad \bar{H} = H_0 + \bar{V}$$
$$\delta H = \delta V = V(x) - \bar{V}, \quad \bar{V} = \frac{1}{L} \int_0^L dx \ V(x)$$

### $\bar{V}$ depends on baseline

The transition matrix is found

$$M(L) = \bar{U}_m e^{-i\frac{\Delta}{2E}L} (1 - iC) \bar{U}_m^{\dagger}$$
$$C = \int_0^L dx \ e^{i\frac{\Delta}{2E}L} \bar{U}_m^{\dagger} \delta V(x) \bar{U}_m e^{-i\frac{\Delta}{2E}L}$$

This is a perturbation expanded using  $\delta V$  around the baseline dependent average potential  $\bar{V}$ 

So

$$C_{jj} = \int_0^L dx (\bar{U}_m^{\dagger} \ \delta V(x) \ \bar{U}_m)_{jj} = 0,$$
  

$$C_{jk} = \int_0^L dx \ e^{i\frac{\Delta_j - \Delta_k}{2E}} (\bar{U}_m^{\dagger} \ \delta V(x) \ \bar{U}_m)_{jk}, \ j \neq k$$

 $|C_{jk}| \ll 1$  needed as a good perturbation approximation

C<sub>jk</sub> suppressed by 1)  $\delta V_e / \bar{V}_e \lesssim 0.3$ 2) small  $\Delta m_{21}^2 / (4E\bar{V}_e)$  and  $\sin \theta_{13}$ 

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 $\nu_{\mu} \rightarrow \nu_{e}$  conversion vs. E



For  $L \lesssim 6000$  km, the Earth matter effect is very well described by the baseline dependent average potential, the only parameter for a fixed baseline

Plus 1st order correction, the theory perfectly describes the oscillation of high energy  $\nu$  in the Earth

#### Time reversal asymmetry vs. E



# The CP violating effect in Earth matter is very well described by this formulation of $\nu$ oscillation

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 $\nu_{\mu}-\nu_{\tau}$  is mainly vacuum oscillation, as can be seen  $\nu_e-\nu_{\tau}$  oscillation is very well described by the perturbation theory

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# Conclusions

### Conclusions

- ▶ We found two nice theories which perfectly describe neutrino oscillation of interested energy range in Earth matter.
- ► Earth matter effect of interested energy range is well understood in these theories.
- ► For E ≤ 30 MeV, the adiabatic perturbation theory says that main contributions are from potential jumps between layers of the Earth matter and averaging over energy can tremendously simplify the oscillation pattern.

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- ► For E ≥ 500 MeV, expanding the potential around the baseline dependent average potential gives a perturbation theory
- This perturbation theory perfectly describes the probability and CP violating effect of oscillation of high energy neutrinos
- Ambiguities of Earth matter effect in measuring neutrino parameters in future experiments, in particular in measuring  $\delta_{CP}$ , are properly addressed in the framework of this perturbation theory.

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